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SOME ESTIMATES OF GREEN'S FUNCTIONS IN THE SHADOW

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0. Introduction

The purpose of this work is to investigate the asymptotic behaviour of Green's functions in the so-called shadow for Laplace operator in an exterior domain. As a consequence a field scattered by a non-trapping obstacle will be examined at high frequencies.

These asymptotics have been studied by many authors since Keller's article [6] appeared. It was shown that for some convex obstacles the scattered field in the shadow should be as small as the exponent $\exp(-A|k|^{1/3})$, A>0, is when the frequency k tends to infinity. Such an estimate is believed to take place for a large class of domains but it has not been proved yet even for strictly convex obstacles except for some special cases. In [12], Ludwig constructed an asymptotic solution u_N for Helmholtz equation in the deep shadow which behaved like $\exp(-A|k|^{1/3})$, A>0, as $k\to\infty$, but he did not show that the difference between u_N and the exact solution could be estimated by the same exponent.

The asymptotics of Green's functions in the shadow were investigated in [1], [2], [3], [14]. Recently, an asymptotic solution of Green's functions in the deep shadow was obtained by Zayaev and Philippov [4] for planar strictly convex obstacles. Probably, the technique developed in [8], [9], [11] may be used to obtain the asymptotic expansions of Green's functions at high frequences for any strictly convex obstacle in \mathbb{R}^n , $n \ge 2$.

Let K be a compact in \mathbb{R}^n , $n \ge 2$, with a real analytic boundary Γ and let $\Omega = \mathbb{R}^n \setminus K$. The obstacle K is called non-trapping if for any R > 0 with $K \subset B_R = \{x \in \mathbb{R}^n; |x| \le R\}$ there exists $T_R > 0$ such that there are no generalized geodesics, (for definition see [13]), with length T_R within $\overline{\Omega} \cap B_R$. Denote by Δ_0 , respectively by Δ_D , the self-adjoint extension of the Laplace operator in \mathbb{R}^n , respectively in Ω with Dirichlet boundary conditions. Let

$$R_j^{\pm}(k) = (-\Delta_j - k^2)^{-1}$$

be the resolvent of the operator $-\Delta_j$, j=0, D in $\pm \text{Im}k > 0$. Consider the

cut-off resolvents

$$(0.1) \quad \{k \in C; \pm \mathrm{Im}k > 0\} \ni k \to R^{\pm}_{D,x}(k) = \chi R^{\pm}_{D}(k) \ \chi \in \mathcal{L}(L^{2}(\Omega), L^{2}(\Omega))$$

where $\chi \in C_{(0)}^{\infty}(\overline{\Omega}) = \{ \varphi \in C^{\infty}(\overline{\Omega}) ; \text{ supp} \varphi \text{ is compact} \}$ and $\chi(x) = 1$ in a neighbourhood of Γ . Hereafter $\mathcal{L}(H_1, H_2)$ stands for the Banach space of bounded linear operators mapping from the Banach space H_1 into the Banach space H_2 and equipped with the usual norm. Obviously the functions (0.1) are analytic with respect to k in + Imk > 0.

Our first result is

Theorem 1. Suppose K non-trapping. Then the function (0.1) admits an analytic continuation in the region

$$U_{\boldsymbol{\alpha},\boldsymbol{\beta}}^{\pm} = \{\boldsymbol{k} \in \boldsymbol{C}; \ \mp \operatorname{Im} \boldsymbol{k} \leq \boldsymbol{\alpha} | \boldsymbol{k} |^{1/3} - \boldsymbol{\beta} \}$$

for some positive constants α and β .

This theorem was proved for strictly convex obstacles with C^{*} boundaries and for n=3 by Babich and Grigorieva [2]. Recently, in [8], [9], Bardos, Lebeau and Rauch showed that the region $U_{\overline{a},\beta}^-$ is free of poles of the scattering matrix for any non-trapping obstacle with an analytic boundary, provided $n\geq 3$ odd. They investigated the generator B of the semi-group Z(t) introduced by Lax and Phillips in [7]. Using the propagation of the Gevrey singularities of the solutions of the mixed problem for the wave equation they proved the estimate $||B^jZ(t_0)|| \leq AC^j(3j)!$ for some t_0 and for any $j \in \mathbb{Z}^+$. Then the region $U_{\overline{a},\beta}^-$ does not contain poles of the scattering matrix according to the results in [7], §3. This result can be obtained also from Theorem 1 since the poles of the scattering matrix coincide with the poles of the meromorphic continuation of $R_{\overline{p},x}(k)$.

A result close to Theorem 1 was proved by Vainberg [18] and Rauch [16] when K is non-trapping and Γ is smooth. In this case the functions (0.1) have analytic continuations in $\{k \in C; \mp \text{Im } k \leq \alpha \text{ Log } |k| - \beta\}$. It is an open problem if Theorem 1 can be extended to hold for any smooth, non-trapping obstacle.

Let us now consider the distribution kernel $G^+(k, x, y)$ $(G^-(k, x, y))$ of the resolvent $R_{\overline{D}}^{\pm}(k)$ in $\pm \text{Im } k \ge 0$ which is usually called outgoing (incoming) Green's function. For any k > 0 the distribution $G^{\pm}(k, x, y)$ solves the problem

(0.2)
$$\begin{bmatrix} (\Delta+k^{2})G^{\pm}(k, x, y) = -\delta(x-y), & (x, y) \in \Omega \times \Omega \\ B G^{\pm} = 0 \\ G^{\pm}(k, x, y) = 0(r^{(1-n)/2}), & \frac{dG^{\pm}}{dr} \mp ik \ G^{\pm} = o(r^{(1-n)/2}) \\ as \ r = |x-y| \to \infty \quad and \quad k \in \mathbf{R}_{+}^{1} = (0, \infty) \end{bmatrix}$$

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where $B u = u_{/\Gamma}$.

The point $x_0 \in \overline{\Omega}$ belongs to the shadow $Sh(y_0)$ of K with respect to a given point $y_0 \in \overline{\Omega}$ if there are no generalized geodesics starting at y_0 and passing through x_0 . Denote by d(x, y) the distance function in $\overline{\Omega}$, i.e.

 $d(x, y) = \inf \{ \text{length of } \gamma; \gamma \text{ is a path in } \overline{\Omega} \text{ connecting } x \text{ and } y \}$.

Denote $D_x^p = D_1^{p_1} \cdots D_n^{p_n}$, where $D_j = i^{-1} \partial \partial x_j$ and $p = (p_1, \cdots, p_n) \in \mathbb{Z}_+^n, \mathbb{Z}_+ = \{0, 1, \cdots\}$.

Theorem 2. Suppose K non-trapping and $x_0 \in Sh(y_0)$. Then there exists a neighbourhood \mathcal{O} of (x_0, y_0) in $\overline{\Omega} \times \overline{\Omega}$ such that

$$(0.3) |D_k^m D_x^b D_y^q G^{\pm}(k, x, y)| \le C \exp(-A|k|^{1/3} \mp d(x, y) \operatorname{Im} k)$$

in $U_{\alpha,\beta}^{\pm} \times \mathcal{O}$ for any $(m, p, q) \in \mathbb{Z}_{+}^{2n+1}$ and for some positive constants α, β, A , and C = C(m, p, q).

Now consider the scattering of plane waves by the obstacle K. Let $\omega \in S^{n-1} = \{\theta \in \mathbb{R}^n; |\theta| = 1\}$ and denote $L_s = \{x \in \mathbb{R}^n; \langle x, \omega \rangle = s\}$ where $\langle x, \omega \rangle = \sum_{i=1}^n x_i \omega_i$. Consider the solution $u_s(k, x)$ of the problem

$$\begin{split} & (\Delta+k^2) \, u_S(k, \, x) = 0 \\ & u_{S/x \in \Gamma} = -e^{ik\langle x, \omega \rangle} / x \in \Gamma \\ & \cdot \, u_S = O(r^{(1-n)/2}), \, \frac{d}{dr} \, u_S - iku_S = o(r^{(1-n)/2}) \quad \text{as} \quad r = |x| \to \infty \; . \end{split}$$

The point x_0 belongs to the shadow $Sh(K, \omega)$ of K with respect to a given direction ω if non of the generalized geodesics $\gamma(t)$, t>0, starting at L_s for some $s < \min_{y \in \Gamma} \langle y, \omega \rangle$ and having ω as an initial direction passes through the point x_0 (t is the natural parameter on γ).

Theorem 3. Suppose K non-trapping and $x_0 \in Sh(K, \omega)$. Then there exists a neighbourhood \mathcal{O} of x_0 in $\overline{\Omega}$ such that

(0.4) $|D_k^m D_x^p(u_s(k, x) + e^{ik\langle x, w \rangle})| \le C \exp(-A |k|^{1/3})$

in $[k_0, \infty) \times \mathcal{O}$ for some A > 0 and any $k_0 > 0, m \ge 0, p \in \mathbb{Z}_+^n$.

An immediate consequence of (0.4) is the Kirchoff approximation of $\frac{\partial}{\partial \nu} u_{S/\Gamma}$ in the shadow, where ν is the outward normal to Γ .

An estimate close to (0.3) was obtained for strictly convex obstacles in [2]. Moreover, some asymptotic expansions in the shadow for x_0 and y_0 sufficiently close to Γ and n=2 were recently obtained by Zayaev and Philippov in [4]. Provided $x_0 \in Sh(y_0)$ and Γ smooth the Green's functions G^{\pm} were

estimated in [14] as follows

$$|G^{\pm}(k, x, y)| \leq C_N k^{-N}$$

for any N > 0 in $k \ge k_0 > 0$ and (x, y) in a neighbourhood of (x_0, y_0) .

The estimate (0.4) was predicted by Keller's geometrical theory of diffraction [5], [6], see also [12].

The method we use is close to that developed by Vainberg [18] (see also [16]) in order to prove uniform decay of the local energy for hyperbolic equations. The propagation of Gevrey singularities for the mixed problem studied in [10], [11] and the non-trapping condition allow us to compare the solutions of the mixed problem with suitably chosen solutions of the Cauchy problem for the wave equation. This is used in Proposition 1 to prove that the kernels of the cut-off resolvents $R_{D,x}^{\pm}(k)$ coincide with the Fourier transforms of some compactly supported distributions modulo exponentially decreasing functions, holomorphic in $U_{\sigma,\beta}^{\pm}$. The theorems follow from Proposition 1 by using once more the results on the propagation of Gevrey singularities for the mixed problem.

1. Estimates of Green's functions

In this section we prove theorems 1 and 2. Let us denote by $U_0(t)$ and U(t) the propagators of the Cauchy problem and the mixed problem respectively, i.e.

(1.1)
$$\begin{cases} (\partial_t^2 - \Delta) \ U_0(t) f(x) = 0 & \text{in } (t, x) \in \mathbf{R}^1 \times \mathbf{R}^n \\ U_0(0) f(x) = 0, \quad \partial_t U_0(0) f(x) = f(x), \quad f \in C_0(\mathbf{R}^n), \end{cases}$$

(1.2)
$$\begin{cases} (0_t - \Delta) \ O(t) f(x) \equiv 0 & \text{in } (t, x) \in \mathbf{K}^* \times \Omega \\ B \ U(t) f(x) = 0 \\ U(0) f(x) = 0 , \quad \partial_t U(0) f(x) = f(x) , \quad f \in C_0^{\infty}(\Omega) . \end{cases}$$

Using standard energy estimates one can extend the operators $U_0(t)$ and U(t) by continuity in $L^2(\mathbb{R}^n)$ and in $L^2(\Omega)$ respectively. Recall that a function f(z) defined in a domain $M \subset \mathbb{R}^p$ belongs to the Gevrey class $G^s(M)$, $s \ge 1$, if for any compact $M_1 \subset M$ there exist some constants $A = A(M_1, f)$, $B = B(M_1, f)$ such that

$$\sup_{z \in \mathcal{M}_1} |D^{\alpha} f(z)| \leq A B^{|\alpha|} (\alpha!)^s$$

for any α , $|\alpha| = \alpha_1 + \cdots + \alpha_n$, $\alpha! = (\alpha_1!) \cdots (\alpha_n!)$.

Let $\chi \in G^3(\mathbb{R}^n)$, $\chi(x)=1$ in a neighbourhood of $B_R = \{x; x \leq R\}$ and $\chi(x) = 0$ for $x \notin B_{R_1}$ for some $R_1 > R$. In view of the non-trapping condition there exists $T > R_1$ such that any generalized geodesic starting at B_{R_1} leaves it by

the time T. Then from the theorem about the propagation of Gevrey G^3 singularities proved by G. Lebeau [10] follows that the distribution kernel U(t, x, y) of U(t) is a G^3 function in

$$Q_{0} = [\mathbf{R}^{1} \setminus (-T, T)] \times (B_{R_{1}} \cap \overline{\Omega}) \times (B_{R_{1}} \cap \overline{\Omega}).$$

Therefore the estimate

(1.3)
$$|D_t^j D_x^{\alpha} D_y^{\beta} U(t, x, y)| \leq A_Q C_Q^{j+|\alpha|+|\beta|} ((j+|\alpha|+|\beta|)!)^3$$

holds in $(t, x, y) \in Q$ for any compact $Q \subset Q_0$ and any j, α, β . Moreover the constants A_Q and C_Q do not depend on $(t, x, y) \in Q$ and on j, α, β .

Let $\zeta \in G^3(\mathbb{R}^{n+1})$, $\zeta = 1$ in a neighbourhood of the set $\{(t, x) \in \mathbb{R}^{n+1}; ||x| -t| < T\}$ and $\zeta(t, x) = 0$ if ||x| - t| > T + 1. Consider the operators

$$U_{\mathtt{x}}(t) = {\mathtt{X}} U(t) {\mathtt{X}}, \quad U_{\mathtt{0},\mathtt{x}}(t) = {\mathtt{X}} U_{\mathtt{0}}(t) {\mathtt{X}}, \quad E(t) = {\mathtt{\zeta}} U(t) {\mathtt{X}},$$

Next we write the modified resolvent $R_{D,x}^+(k)$ in the form

(1.4)
$$R_{D,x}^+(k) = \chi \hat{E}(k) + Z_x(k)$$

where

$$\chi \hat{E}(k) = \int_0^\infty e^{ikt} \chi E(t) dt$$
, Im $k > 0$,

denotes the Fourier-Laplace transform of $\chi E \in L^1(\mathbb{R}^1, \mathcal{L}(L^2(\Omega), L^2(\Omega)))$. Note that the operator-valued function $\chi E(t)$ has a compact support with respect to t since $\chi(x) \zeta(t, x)$ has. Therefore $\chi \hat{E}(k)$ is an analytic function with values in the space $\mathcal{L}(L^2(\Omega), L^2(\Omega))$, while $Z_{\chi}(k)$ is analytic in $\{k \in \mathbb{C}; \text{ Im } k > 0\}$. Let $H^s(\Omega), s \ge 0, s \in \mathbb{Z}$, be the closure of $C^{\infty}_{(0)}(\Omega)$ with respect to the Sobolev norm $||u||_s^2 = \sum_{|x| \le s} ||D^{\alpha}u||_{L^2(\Omega)}^2$ and let $H^{-s}(\Omega)$ be the dual space of $H^s(\Omega)$. We shall use also the domain D^s of the operator $(-\Delta_D)^{s/2}, s \ge 0, s \in \mathbb{Z}$, equipped with the graph topology, where the operator $(-\Delta_D)^{s/2}$ is given by the functional calculas. Denote by D^{-s} the dual space of D^s . Theorems 1 and 2 will follow from

Proposition 1. The function $Z_{x}(k)$ can be extended as an analytic function

$$\{k \in C; \operatorname{Im} k > 0\} \ni k \mapsto Z_{x}(k) \in \mathcal{L}(H^{-s}(\Omega), H^{s}(\Omega))$$

for any $s \ge 0$, $s \in \mathbb{Z}$. Moreover, there exist some positive constants α and β such that $Z_x(k)$ has an analytic continuation in $U^+_{\alpha,\beta}$ and

(1.5)
$$||D_k^m Z_{\chi}(k)||_{\mathcal{L}(H^{-s},H^s)} \leq C \exp(-A|k|^{1/3} - T \operatorname{Im} K), \quad m \geq 0,$$

in $k \in U_{\alpha,\beta}^{+}$ for some positive constants A and C = C(m, s).

Proof. Let us denote $F(t) = [\partial_t^2 - \Delta, \zeta] U(t) \chi$, where $[F_1, F_2] = F_1 F_2 - F_2 F_1$

is the commutator of the operators F_1 and F_2 and ζ stands for the operator of multiplication by the function $\zeta(t, x)$. Then E(t) is the propagator of the problem

(1.6)
$$\begin{bmatrix} (\partial_t^2 - \Delta) E(t) f(x) = F(t) f(x) \\ BE(t) f = 0 \\ E(0) f(x) = 0, \quad \partial_t E(0) f(x) = \chi(x) f(x), \quad f \in L^2(\Omega). \end{bmatrix}$$

The distribution kernel F(t, x, y) of the operator F(t) belongs to the Gevrey class $G^3(\mathbf{R}^1 \times \overline{\Omega} \times \overline{\Omega})$ in view of the propagation of Gevrey singularities of U(t, x, y) and the definition of the functions $\zeta(t, x)$ and $\chi(x)$. Moreover

(1.7)
$$\sup F \subset \{(t, x, y) \in \mathbb{R}^{1} \times \overline{\Omega} \times \overline{\Omega}; |t| > T, T \le ||x| - t| \le T + 1, |y| \le \mathbb{R}_{1}\}$$

in view of the finite propagation speed for the wave equation.

Let $\tilde{F}(t, x, y)$ be a G^3 continuation of the function F(t, x, y) such that (1.7) continues to hold. Denote by $\tilde{F}(t)$ the operator with a distribution kernel $\tilde{F}(t, x, y)$ and consider the problem

(1.8)
$$\begin{cases} (\partial_t^2 - \Delta) \ W(t) \ f(x) = \widetilde{F}(t) \ f(x) \\ W(0) \ f(x) = \partial_t W(0) \ f(x) = 0, \quad f \in C_0^{\infty}(\mathbf{R}^n) \end{cases}$$

The distribution kernel W(t, x, y) of W(t) is a G^3 function since the function $\widetilde{F}(t, x, y)$ is such, $\widetilde{F}(t)=0$ in |t| < T and since

$$W(t) = \int_0^t U_0(s) \, \widetilde{F}(t-s) ds$$

Let $\psi \in C^{\infty}(\mathbb{R}^n)$, $\psi(x)=0$ in a neighbourhood of B_R and $\chi(x)=1$ on supp $(1-\psi)$. Denote

$$Q(t) f(x) = (\partial_t^2 - \Delta) (E(t) f(x) - \psi W(t) f(x))$$

= (1-\psi) F(t) f(x)+[\Delta, \psi] W(t) f(x)

in $x \in \overline{\Omega}$ for $f \in C^{\infty}_{(0)}(\overline{\Omega})$. In view of (1.6), (1.7) and Duhamel's formula we obtain

$$E(t)f - \psi W(t)f = U(t) \chi f + \int_0^t U(t-s) \chi Q(s) f ds, \quad f \in L^2(\Omega).$$

Multiplying the last equality by χ and performing Fourier-Laplace transform with respect to t we obtain

(1.9)
$$\chi \hat{E}(k) f = R^{+}_{D,x}(k) f + R^{+}_{D,x}(k) \hat{Q}(k) f + \psi \chi \hat{W}(k) f$$

for Im k>0. We are going to prove that the functions $\psi \chi \hat{W}(k)$ and $\hat{Q}(k)$ can

be continued analytically for $\text{Im } k \leq 0$.

Let $\mathcal{H} \in C^{\infty}(\mathbf{R}^n)$, $\mathcal{H}(x) = 0$ for $x \in B_T$, $\mathcal{H}(x) = 1$ outside B_{T+1} and set

$$G(t) f(x) = (\partial_t^2 - \Delta) (W(t) f(x) - \mathcal{H}(x) E(t) f(x))$$

= $(1 - \mathcal{H}) \widetilde{F}(t) f(x) - [\Delta, \mathcal{H}] E(t) f(x).$

The function $\mathbf{R}^1 \ni t \mapsto E(t) \in \mathcal{L}(D^{-s}, D^{-s+1})$ is bounded for any $s \in \mathbf{Z}$, $[\Delta, \mathcal{H}] \in \mathcal{L}(D^{-s+1}, H^{-s}(\mathbf{R}^n))$, and $H^{-s}(\Omega) \subset D^{-s}$ for any $s \ge 0$, $s \in \mathbf{Z}$. Then $\mathbf{R}^1 \ni t \mapsto [\Delta, \mathcal{H}] E(t)$ is a bounded function with values in $\mathcal{L}(H^{-s}(\Omega), H^{-s}(\mathbf{R}^n))$, $s \ge 0$, $s \in \mathbf{Z}$, and

$$(1.10) ||G(t)||_{\mathcal{L}(H^{-s}(\Omega), H^{-s}(\mathbf{R}^n))} \leq C$$

for any $t \in \mathbf{R}^1$.

In view of (1.6), (1.7), (1.10) and Duhamel's formula we write

$$W(t) f(x) = \mathcal{H}(x) E(t) f(x) + \int_0^t U_0(t-s) G(s) f(x) ds, f \in H^{-s}(\Omega).$$

Note that the support of the distribution kernel of G(t) is contained in $\{(t, x, y); |t| \le 2T+2, |x| \le T+1, |y| \le T+1\}$. Therefore

(1.11)
$$\chi_2 W(t) f = \chi_2 \int_0^{T_1} U_0(t-s) \chi_1 G(s) f ds, \quad f \in H^{-s}(\Omega)$$

for any $T_1 > 2T + 2$, where $\chi_1 \in C_0^{\infty}(\mathbb{R}^n)$, $\chi_1(x) = 1$ in B_{T+1} and $\chi_2 \in C_0^{\infty}(B_T)$.

Lemma 1. Let $\chi_2 \in C_0^{\infty}(B_T \setminus B_R)$. Then $\chi_2 U_0(t) \chi_1 \in \mathcal{L}(H^{-s}(\mathbb{R}^n), H^s(\Omega))$ for any $s \in \mathbb{R}^1$ and any $t \in [2T+3, \infty)$. Moreover the function

$$[2T+3, \infty) \ni t \mapsto \chi_2 U_0(t) \chi_1 \in \mathcal{L}(H^{-s}(\mathbf{R}^n), H^s(\Omega))$$

can be continued analytically in $\{t \in C; |t| > 2T+3\}$ and

(1.12)
$$||D_{t}^{j} \chi_{2} U_{0}(t) \chi_{1}||_{\mathcal{L}(H^{-s}(\mathbf{R}^{n}), H^{s}(\Omega))} \leq A(j!) |t|^{-2}$$

for any $t \in C$, |t| > 2T+3, for $j \ge \max(0, 3-n)$, and for some A which does not depend on j.

Proof. The conclusion is obvious when n is odd because of Huyghens principle. Suppose $n \ge 2$ is even, $j \ge 1$, and set $\mathcal{O}_T = \{(t, x, y) \in \mathbb{C}^{2n+1}; |t| > 2T + 3, |x| \le T, |y| \le T+1\}$. Then $U_0(t, x, y) = C_n(t^2 - |x-y|^2)^{-(n-1)/2}$ for any $(t, x, y) \in \mathcal{O}_T$ and for some constant C_n . Using Cauchy integral formula we obtain for any $j \ge 1$, α , β the estimate

$$\begin{aligned} |D_t^{j} D_x^{\omega} D_y^{\beta} U_0(t, x, y)| &\leq (2\pi)^{-2n-1} (j-1)! (\alpha + \beta)! 2^{|\omega + \beta|} \max\{|D_t U_0(z, \tilde{x}, \tilde{y})|; \\ |z - t| &= 1, \ |x - \tilde{x}| + |y - \tilde{y}| = 1/2\} \leq A_{\omega, \beta}(j!) |t|^{-2} \end{aligned}$$

in \mathcal{O}_T which yields (1.12).

According to (1.10), (1.11) and lemma 1 the function

$$[T_2, \infty) \ni t \to \chi_2 W(t) \in \mathcal{L}(H^{-s}(\Omega), H^s(\Omega)), \quad T_2 = 2T_1 + 2,$$

can be continued as an analytical one in $\{t \in C; |t| > T_2\}$ for any $t \ge 0$ and any $\chi_2 \in C_0^{\infty}(B_T \setminus B_R)$. Moreover the estimate

(1.13)
$$||D_{t}^{j} \chi_{2} W(t)||_{\mathcal{L}(H^{-s}, H^{s})} \leq A(3j)! |t|^{-2}$$

is valid in $|t| > T_2$ for any $j \ge \max(0, 3-n)$ and any $s \ge 0$, $s \in \mathbb{Z}$ where the constant A does not depend on j.

Now we can estimate the norm of the Fourier-Laplace transform of $\chi_2 W(t)$ in $\mathcal{L}(H^{-s}, H^s)$. Let Re $k \ge k_0 > 0$ for some $k_0 > 0$. Since W(t) = 0 in |t| < Twe can write

$$\chi_2 \hat{W}(k) = k^{-1} \int_0^{T_2} e^{ikt} D_t \chi_2 W(t) dt + k^{-1} \int_{T_2}^{\infty} e^{ikt} D_t \chi_2 W(t) dt$$

Using (1.13) we can change the contour of integration in the second integral to obtain

$$\begin{split} \exp(C \,|\,k\,|^{1/3}) \,\chi_2 \hat{W}(k) &= \sum_{j=0}^{\infty} \,C^j \,|\,k\,|^{j/3-1} (j!)^{-1} \,[\int_0^{T_2} e^{ikt} \,D_t \chi_2 W(t) \,dt \\ &+ e^{ikT_2} \int_0^{\infty} e^{-kt} \,\chi_2(D_t W) \,(T_2 + it) \,dt] \,. \end{split}$$

Integrating [j/3] times by parts in any member of the last sum we have

$$\exp(C |k|^{1/3}) \chi_2 \hat{W}(k) = \sum_{j=0}^{\infty} C^j k^{j/3 - [j/3] - 1} (j!)^{-1} \\ \left[\int_0^{T_2} e^{ikt} \chi_2 D_t^{[j/3] + 1} W(t) dt + e^{ikT_2} \int_0^{\infty} e^{-kt} \chi_2 (D_t^{[j/3] + 1} W) (T_2 + it) dt \right]$$

where [m] denotes the integer part of $m \in \mathbb{R}^1_+$. Since $W \in G^3$ and in view of (1.13) any member of the last sum can be estimated by

$$A_1 C^j B_1^{j/3} e^{-B \operatorname{Im} k}$$
, $B = \begin{cases} T & \text{when Im } k \ge 0 \\ T_2 & \text{when Im } k < 0 \end{cases}$

in $\{k \in C; \text{ Re } k \ge k_0 > 0\}$, where the constants A_1 and B_1 do not depend on $j \in \mathbb{Z}$. Provided that $C < B_1^{-1/3}$ we obtain

(1.14)
$$||\chi_2 \ddot{W}(k)||_{\mathcal{L}(H^{-s}(\Omega), H^s(\Omega))} \leq C_s \exp(-C |k|^{1/3} - B \operatorname{Im} k)$$

for Re $k \ge k_0 > 0$, where $C_s = A_1(1 - CB_1^{1/3})^{-1}$. Proceeding in the same way when Re $k \le -k_0 < 0$ we can continue $\chi_2 \hat{W}(k)$ analytically in $C \setminus \{k; \text{ Im } k \le 0, | \text{ Re } k| \le k_0\}$ so that (1.14) holds in this region for any $k_0 > 0$. Then the Fourier-Laplace transform $\hat{Q}(k)$ of $Q(t) = (1 - \psi)F(t) + [\Delta, \psi] W(t)$ can be continued analytically in $C \setminus [i0, -i\infty)$ and

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(1.15)
$$||\hat{Q}(k)||_{\mathcal{L}(H^{-s}(\Omega),H^{s}(\Omega))} \leq C_{s} \exp(-C |k|^{1/3} - B \operatorname{Im} k)$$

is fulfilled in $C \setminus \{k; \operatorname{Im} k \leq 0, |\operatorname{Re} k| \leq k_0\}$ for any $k_0 > 0$.

Lemma 2. The function $C \ni k \mapsto \chi \hat{E}(k) \in \mathcal{L}(H^{s}(\Omega), H^{s}(\Omega))$ is analytic and

(1.16)
$$||\chi \hat{E}(k)||_{\mathcal{L}(H^{s}(\Omega), H^{s}(\Omega))} \leq C(1+|k|)^{2s} e^{(2T+1)\max(0, -\operatorname{Im} k)}$$

for any $s \ge 0$, $s \in \mathbb{Z}$.

Proof. The assertion is obvious for t=0 since U(t) is a bounded function in \mathbb{R}^1 with values in $\mathcal{L}(L^2(\Omega), L^2(\Omega))$ and $\chi \zeta(t, x)=0$ for any t>2T+1. Suppose $s\geq 1$ and consider

(1.17)
$$\begin{bmatrix} (\Delta - 1) \chi \hat{E}(k) f = L(k) f - \chi f + ([\Delta, \chi] - (k^2 + 1) \chi) \hat{E}(k) f \\ \chi \hat{E}(k) f_{/\Gamma} = 0 \end{bmatrix}$$

for $f \in H^{s}(\Omega)$. Here

$$L(k)f = -\int_0^\infty e^{ikt} \, \chi F(t) \, f dt \in H^s(\Omega)$$

and L(k) satisfies the estimate (1.16) for any $s \ge 0$ since the distribution kernel of the operator $\chi F(t)$ is smooth and $\operatorname{supp}(\chi F) \subset \{|x| \le R_1, |y| \le R_1, |t| < 2T + 1\}$ in view of (1.7). Then

$$||\chi \hat{E}(k)f||_{s} \leq C((1+|k|^{2})||\chi_{1}\hat{E}(k)f||_{s-1}+e^{(2T+1)\max(0,-\operatorname{Im} k)}||f||_{s})$$

 $f \in H^{s}(\Omega)$, for some $\chi_{1} \in C_{(0)}^{\infty}(\overline{\Omega})$, $\chi_{1}=1$ in a neighbourhood of supp(χ) which proves (1.16) by induction. Differentiating (1.17) with respect to k and using (1.16) it is easy to prove that $\frac{d}{dk} \chi \hat{E}(k) \in \mathcal{L}(H^{s}(\Omega), H^{s}(\Omega))$ for any $s \geq 0$, $s \in \mathbb{Z}$. Thus $\chi \hat{E}(k)$ is an analytic function.

According to (1.15) the operator $I + \hat{Q}(k)$: $H^{s}(\Omega) \mapsto H^{s}(\Omega)$ is invertable for any $k \in U^{+}_{\alpha,\beta}$ and for some α, β . Then $R^{+}_{D,x}(k)$ is an analytic function in $U^{+}_{\alpha,\beta}$ with values in $\mathcal{L}(H^{s}(\Omega), H^{s}(\Omega))$ and satisfies (1.16) in view of (1.9) and Lemma 2. Now, (1.5) follows for m=0 from (1.9), (1.14) and (1.15), choosing α and β small enough. Using Cauchy integral formula we obtain (1.5) for any $m \in \mathbb{Z}_{+}$.

To prove theorem 2 we choose some neighbourhoods \mathcal{O}_1 and \mathcal{O}_2 of x_0 , respectively y_0 , $\mathcal{O}_j \subset \overline{\Omega}$, so that none of the generalized geodesics starting at \mathcal{O}_2 passes through \mathcal{O}_1 . Set $\mathcal{O}=\mathcal{O}_1 \times \mathcal{O}_2$ and suppose that $\mathcal{O}_j \subset B_R$ and T>sup $\{d(x, y); (x, y) \in \mathcal{O}\}$. According to proposition 1 we have

$$G^{+}(k, x, y) = \int_{0}^{\infty} e^{ikt} \zeta(t, x) U(t, x, y) dt + Z_{x}(k, x, y)$$

where

$$|D_{k}^{m}D_{x}^{p}D_{y}^{q}Z_{x}(k, x, y)| = |\langle D_{x}^{p} \delta_{x}, D_{k}^{m} Z_{x}(k) D_{y}^{q} \delta_{y} \rangle| \leq ||D_{k}^{m}Z_{x}(k)||_{\mathcal{L}(H^{-s}, H^{s})} ||\delta_{x}||_{p+q-s}^{2}$$

$$\leq C \exp(-A|k|^{1/3} - T \operatorname{Im} k) \leq C \exp(-A_{0}|k|^{1/3} - d(x, y) \operatorname{Im} k)$$

in $U^+_{\alpha,\beta} \times \mathcal{O}$ for some $\alpha > 0$ and $A_0 > 0$. Here $\langle \delta_x, \varphi \rangle = \varphi(x)$ for any $\varphi \in C^{\infty}_{(0)}(\overline{\Omega})$ and s > n+p+q. On the other hand $\zeta(t, x) U(t, x, y)$ is a G^3 function in $\mathbb{R}^1 \times \mathcal{O}$ with a compact support with respect to t. Moreover, U(t, x, y) = 0 for |t| < d(x, y) since the propagation speed for the solutions of the mixed problem for the wave equation equals one (see [17]). Now the arguments used in the proof of (1.14) yield (0.3).

Denote by $e(\lambda, x, y)$ the spectral function of the operator $-\Delta_D$ given as the distribution kernel of the spectral projector E_{λ} of $-\Delta_D$. Since $E_{\lambda} \rightarrow I$ in $L^2(\Omega)$ as $\lambda \rightarrow \infty$ and

$$\frac{de}{d\lambda}(\lambda^2, x, y) = (2\pi i)^{-1} \{G^+(\lambda, x, y) - G^-(\lambda, x, y)\} \quad \text{for} \quad x \neq y, \, \lambda > 0 \,,$$

it is easy to obtain from theorem 2 the following

Corollary 1. Suppose K non-trapping and $x_0 \in Sh(y_0)$. Then

 $|D_{\lambda}^{m} D_{x}^{p} D_{y}^{q} e(\lambda, x, y)| \leq C \exp(-A \lambda^{1/6}), \quad A > 0,$

in $[\lambda_0, \infty) \times \mathcal{O}$ for $(m, p, q) \in \mathbb{Z}_+^{2n+1}, \lambda_0 > 0$.

2. Asymptotics of the scattered waves

In this section we prove theorem 3. Translating the origin to a given point $z_0 \in \mathbb{R}^n$ the function $u_S(k, x)$ is multiplied by $\exp(ik \langle z_0, \omega \rangle)$. Thus we can suppose that $K \subset B_R(x_0) = \{x \in \mathbb{R}^n; |x-x_0| \leq R\}$ and $\langle x, \omega \rangle > 0$ for any $x \in B_{R+1}(x_0)$. Consider the function

$$v(k, x) = u_{S}(k, x) + \varphi(x) e^{ik\langle x, \omega \rangle}$$

where $\varphi \in G^3(B_{R+1}(x_0))$ and $\varphi(x)=1$ on $B_R(x_0)$, supp $\varphi \subset B_{R+1}(x_0)$. Then

$$\begin{bmatrix} (\Delta+k^2) v(k, x) = [\Delta, \varphi] e^{ik\langle x, w \rangle} \\ v(k, x)_{/\Gamma} = 0 \end{bmatrix}$$

and v(k, x) satisfies the outgoing Sommerfeld's condition at infinity. Therefore

$$egin{aligned} v(k,\,x) &= R^+_{D,\,\chi}(k)\,([\Delta,\,arphi]\,e^{ik\langle x,\,oldsymbol{arphi}
angle}) \ &= Z_\chi(k)([\Delta,\,arphi]\,e^{ik\langle x,\,oldsymbol{arphi}
angle}) + \chi \hat{E}(k)\,([\Delta,\,arphi]\,e^{ik\langle x,\,oldsymbol{arphi}
angle)} \end{aligned}$$

for $x \in B_R(x_0)$ where $\chi \in G^3(\mathbb{R}^n)$, $\chi = 1$ on $B_{R+1}(x_0)$, $\operatorname{supp}(\chi) \subset B_{R+2}(x_0)$.

The first term of the last equality is estimated by proposition 1. The second one is equal to the Fourier-Laplace transform of the distribution

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$$v_1(t, x) = \chi(x) \int_{-\infty}^t \zeta(t-s, x) \ U(t-s) \left[\Delta, \varphi\right] \delta(s - \langle x, \omega \rangle) \ ds$$

since $v_2(s, y) = [\Delta, \varphi] \delta(s - \langle y, \omega \rangle)$ vanishes for s < 0. The distribution v_1 is well-defined since v_2 has a compact support, $v_2 \in D^{-m}$ for m > 3 and $\zeta(t-s)$ U(t-s) is a continuous function with valued in $\mathcal{L}(D^{-m}, D^{-m})$.

We are going to prove that there exists a neighbourhood \mathcal{O} of x_0 such that v_1 is a G^3 function in $\mathbb{R}^1 \times \mathcal{O}$.

Let us write $v_1 = Q(v_2)$ where the operator Q has a distribution kernel $Q(t, s, x, y) = \chi(x) \zeta(t-s) H(t-s) U(t, x, y) \chi(y)$ and H(s) = 0 for $s \le 0$, H(s) = 1 for s > 0. We shall evaluate the Gevrey G^3 wave front $SS^3(v_1)$ of v_1 using the relation $SS^3(v_1) \subset SS^3(Q)^1 \circ SS^3(v_2)$. We have

$$SS^{3}(v_{2}) \subset \{(s, y; \tau, \eta); s = \langle y, \omega \rangle > 0, y \notin B_{R}(x_{0}), \eta = -\tau\omega, \tau \neq 0\}$$

Moreover, theorem 1.4 in [10] yields

$$SS^{3}(Q)^{1} \subset \{(\phi^{t-s}(s, y, \tau, \eta); s, y, \tau, \eta); s \leq t, \tau \neq 0\} \cup \{(0, y, \tau, \xi; 0, y, \tau, \eta)\}$$

where $\phi^t(s, y, \tau, \eta) = (t+s, x^t(s, y, \tau, \eta), \tau, \xi^t(s, y, \tau, \eta))$ is the generalized bicharacteristic starting at (s, y, τ, η) and t is the natural parameter on it. Thus we have

$$SS^{3}(v_{1}) \subset \{(t, x^{t-s}(s, y, \tau, -\tau\omega), \tau, \xi); \tau \neq 0, 0 < s = \langle y, \omega \rangle \leq t, y \in B_{\mathbb{R}}(x_{0})\}.$$

Note that the initial codirection of the generalized geodesic $\gamma(t) = x^t(s, y, \tau, \eta)$ is $\frac{d\gamma}{dt}(0) = -\eta/\tau$ for any $y \in \Omega$. Then

$$SS^{3}(v_{1}) \subset \{(t, \gamma(t-s), \tau, \xi); \gamma \text{ is a generalized geodesic with} \\ \gamma(0) \oplus B_{R}(x_{0}), \frac{d\gamma}{dt}(0) = \omega, 0 < s = \langle \gamma(0), \omega \rangle \leq t \}.$$

Moreover $\gamma(t) \notin B_R(x_0)$ for any $t \ge 0$ when $\gamma(0) \notin B_R(x_0)$ and $\langle \gamma(0), \omega \rangle \ge \langle x_0, \omega \rangle$ while $\gamma(t-s)=\gamma_1(t)$, $\gamma_1(t)$ is the generalized geodesic with initial data $\gamma_1(0)=$ $\gamma(0)-s \omega \in L_0$, $\frac{d\gamma_1}{dt}(0)=\omega$, when $\gamma(0) \notin B_R(x_0)$ and $\langle \gamma(0), \omega \rangle \le \langle x_0, \omega \rangle$. Therefore

(sing supp
$$_{G^3}(v_1)) \cap B_R(x_0) \subset \{x = \gamma(t); t > 0 \text{ and } \gamma \text{ is a generalized}$$

geodesic with $\gamma(0) \in L_0, \frac{d\gamma}{dt}(0) = \omega\}$.

Since $x_0 \in Sh(K, \omega)$ we can choose a neighbourhood \mathcal{O} of x_0 such that (sing $\operatorname{supp}_{G^3}(v_1)) \cap \mathcal{O} = \phi$ which proves theorem 3 since $\operatorname{supp}(v_1)$ is compact.

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