



Title	Scattering Theory for First Order Systems with Long-range Perturbations : the Existence of the Modified Wave Operators
Author(s)	菊地, 光嗣
Citation	大阪大学, 1988, 博士論文
Version Type	VoR
URL	https://hdl.handle.net/11094/2459
rights	
Note	

The University of Osaka Institutional Knowledge Archive : OUKA

<https://ir.library.osaka-u.ac.jp/>

The University of Osaka

SCATTERING THEORY FOR FIRST ORDER SYSTEMS
WITH LONG-RANGE PERTURBATIONS:
THE EXISTENCE OF THE MODIFIED WAVE OPERATORS

Koji KIKUCHI

1 Introduction

The equation of acoustic, electromagnetic and elastic waves can be considered as first order symmetric hyperbolic systems of partial differential equations for C^m -valued function u of the form

$$(1.1) \quad D_t u = \Lambda u,$$

where

$$\Lambda = E(x)^{-1} \sum_{j=1}^n A_j D_j.$$

Here $D_t = (1/i)(\partial/\partial t)$ and $D_j = (1/i)(\partial/\partial x_j)$, A_j 's are constant $m \times m$ hermitian matrices, and $E(x)$ is a positive definite hermitian matrix. It is measurable and satisfies

$$0 < c_1 I \leq E(x) \leq c_2 I \quad \text{for some } c_1 \text{ and } c_2.$$

We shall consider the case that there exists a constant matrix E_0 and a positive number δ such that

$$|E(x) - E_0| \leq C \langle x \rangle^{-\delta} \quad (\langle x \rangle = (1 + |x|^2)^{1/2}).$$

Then the equation (1.1) is regarded as the perturbation of the

equation

$$D_t u = \Lambda^0 u,$$

where

$$\Lambda^0 = E_0^{-1} \sum_{j=1}^n A_j D_j.$$

When $\delta > 1$, the perturbation is said to be short-range, and when $\delta \leq 1$, it is said to be long-range. In this paper we shall consider a class of first order systems with long-range perturbations.

Now we put $J = E^{-1/2} E_0^{-1/2}$. The operator W_{\pm} is called the wave operator if the limit

$$(1.2) \quad W_{\pm} u = \lim_{t \rightarrow \pm\infty} e^{it\Lambda} J e^{-it\Lambda^0} P_{ac}(\Lambda^0) u$$

exists. From the point of view of the existence and the completeness of wave operators we can say that the general theory of systems with short-range perturbations is at the satisfactory stage (for example [1], [8]). When the perturbation is long-range, as is known for the case of Schrödinger operators, the limit (1.2) does not exist generally, that is, Λ^0 cannot be regarded as the free operator. Thus our problem is to construct another operator W_{\pm}^D (so-called

modified wave operator) which satisfies the intertwining property

$$e^{is\Lambda} w_{\pm}^D = w_{\pm}^D e^{is\Lambda^0}.$$

The fundamental problems of the theory of long-range perturbations are the existence and the completeness of w_{\pm}^D . In connection with these problems H.Tamura first proved the limiting absorption principle for uniformly propagative systems with long-range perturbations ([5]). This result has been extended to wide classes ([6], [7], [8] and [9]). For the case of 2-body Schrödinger operators the limiting absorption principle alone enables us to prove the existence and completeness of (modified) wave operators even in the case of long-range perturbations. This is also the case for first order hyperbolic systems when all characteristic roots of the unperturbed operator are simple. However, if we consider the case that characteristic roots are non smooth, for example the Maxwell equation in a crystal optics, we encounter serious difficulties in developing the scattering theory by using the above mentioned stationary results.

In this paper we prove the existence of the modified wave operator for some classes of systems by using the time dependent method. In section 2 we shall state the exact conditions for Λ^0 , which includes the Maxwell equations in biaxial crystals. Their characteristic roots are non smooth,

and our theory admits including such equations.

Now we want to define the modified wave operator as the limit

$$(1.3) \quad W_{\pm}^D u = \lim_{t \rightarrow \pm\infty} e^{it\Lambda} J X(t) u$$

for some partial isometric operator $X(t)$ instead of $e^{-it\Lambda^0}$. In the construction of $X(t)$ we shall use the solution of the Hamilton-Jacobi equations

$$(1.4) \quad \frac{\partial w_k}{\partial t} = \lambda_k(\nabla_{\xi} w_k, \xi),$$

where $\lambda_k(x, \xi)$'s are the eigenvalues of the principal symbol of Λ (an $m \times m$ matrix).

In [2] Lars Hörmander has proved the existence of the modified wave operators for Schrödinger equations with long-range potentials. Here he solved the Hamilton-Jacobi equation and used the stationary phase method on \mathbb{R}^n . But in our case, since $\lambda_k(x, \xi)$ is positively homogeneous of degree 1 with respect to ξ , its Hessian vanishes everywhere. Hence we cannot apply the Hörmander calculus, and we shall use the stationary phase method on a hypersurface. Then our calculus is rather different from that of [2].

This paper is organized as follows. In section 2 we shall give the exact assumptions for Λ^0 and some preliminary results. In section 3 we shall define the operator $X(t)$ and

show some properties. In section 4 we shall prove the existence of the limit (1.3) by the use of the stationary phase method on a hypersurface.

The author would like to express his sincere gratitude to Prof.Mitsuru Ikawa and Dr.Tomio Umeda for their encouragements and useful suggestions.

2 Assumptions and some preliminary results

We are now treating the case that the perturbation is long-range. More precisely we assume

Assumption (E) $E(x) \in C^\infty(\mathbb{R}^n)$ and

$$|\partial_x^\alpha(E(x) - E_0)| \leq C|x|^{-\delta-|\alpha|}$$

for $\delta > 0$ and $|\alpha| \geq 0$.

On Λ^0 we also require some assumptions. We put

$$\Lambda^0(\xi) = E_0^{-1} \sum_{j=1}^n A_j \xi_j \quad (\text{the symbol of } \Lambda^0)$$

and

$$\rho_0 = \max_{\xi \in \mathbb{R}^n} \#\{\text{positive eigenvalues of } \Lambda^0(\xi)\}.$$

Then

Assumption (F) 1) Λ^0 is strongly propagative, that is,
for some d

$$\text{rank } \Lambda^0(\xi) = m - d \quad (\text{for } \xi \neq 0).$$

$$2) \quad \rho_0 = (m - d)/2.$$

Here d is called the deficit. The condition 2) of (F) is equivalent to that the multiplicities of non-zero eigenvalues are all simple outside a conic null set. As an example we consider the Maxwell equations in crystals:

$$\nabla \times H - \epsilon \frac{\partial E}{\partial t} = 0, \quad \nabla \times E + \mu_0 \frac{\partial H}{\partial t} = 0.$$

Here $\epsilon = (\epsilon_{ij})$ is a tensor dielectric constant and μ_0 is a scalar magnetic permeability. Let ϵ_1, ϵ_2 and ϵ_3 be eigenvalues of ϵ . We may assume that

$$\epsilon_1 \geq \epsilon_2 \geq \epsilon_3 > 0.$$

There are three classes which are defined by the condition (1) $\epsilon_1 > \epsilon_2 > \epsilon_3$, (2) $\epsilon_1 > \epsilon_2 = \epsilon_3$ or $\epsilon_1 = \epsilon_2 > \epsilon_3$ and (3) $\epsilon_1 = \epsilon_2 = \epsilon_3$ (isotropic). The cases of (1) and (2) are covered in our class, and the case of (3) is not covered (refer to C.H. Wilcox [11]).

Similar to the case of Λ^0 we put

$$\Lambda(x, \xi) = E(x)^{-1} \sum_{j=1}^n A_j \xi_j \quad (\text{the symbol of } \Lambda)$$

and

$$\rho(x) = \max_{\xi \in \mathbb{R}^n} \# \{ \text{positive eigenvalues of } \Lambda(x, \xi) \}.$$

The eigenvalues of $\Lambda^0(\xi)$ and $\Lambda(x, \xi)$ can be enumerated as follows:

$$(2.1) \quad \begin{aligned} \lambda_{\rho_0}^0(\xi) &\geq \dots \geq \lambda_1^0(\xi) > \lambda_0^0(\xi) \equiv 0 \\ &> \lambda_{-1}^0(\xi) \geq \dots \geq \lambda_{-\rho_0}^0(\xi) \end{aligned}$$

and

$$(2.2) \quad \begin{aligned} \lambda_{\rho(x)}(x, \xi) &\geq \dots \geq \lambda_1(x, \xi) > \lambda_0(x, \xi) \equiv 0 \\ &> \lambda_{-1}(x, \xi) \geq \dots \geq \lambda_{-\rho(x)}(x, \xi), \end{aligned}$$

respectively. If $\lambda_j^0(\xi) \equiv \lambda_k^0(\xi)$ ($\lambda_j(x, \xi) \equiv \lambda_k(x, \xi)$ for fixed x), then $j = k$.

Here we note that 2) of (F) guarantees that $\lambda_k(x, \xi)$ is smooth for large $|x|$ when $\lambda_k^0(\xi)$ is smooth. It is not the case if 2) of (F) is not satisfied, and (1.4) does not have classical solutions in general. Hence it seems to be difficult to remove the condition 2) of (F) as long as our method is used.

To study Λ^0 and Λ in one Hilbert space $\mathcal{H} = L^2(\mathbb{R}^n, \mathbb{C}^m)$ with the usual inner product

$$(u, v) = \int_{\mathbb{R}^n} u(x)^* v(x) dx$$

(note that Λ^0 and Λ are not self-adjoint in \mathcal{H}), we set

$$\tilde{\chi}^0 = \sum_{j=1}^n E_0^{-1/2} A_j E_0^{-1/2} D_j$$

and

$$\tilde{\chi} = \tilde{E}(x)^{-1/2} \tilde{\chi}^0 \tilde{E}(x)^{-1/2} \quad (\tilde{E}(x) = E_0^{-1/2} E(x) E_0^{-1/2}),$$

which are self-adjoint operators in \mathcal{H} with domains

$$\mathcal{D}(\tilde{\chi}^0) = \mathcal{D}(\tilde{\chi}) = \{u \in \mathcal{H}; \tilde{\chi}^0 u \in \mathcal{H}\}.$$

Now we can reduce our problem to studying $\tilde{\chi}^0$ and $\tilde{\chi}$ in \mathcal{H} (see for example [9]), and then our purpose is now to construct a partial isometric operator $\tilde{\chi}(t)$ in \mathcal{H} and to show the existence of the limits

$$\tilde{w}_\pm^D = \lim_{t \rightarrow \pm\infty} e^{it\tilde{\chi}} \tilde{\chi}(t).$$

Note that the eigenvalues of $\tilde{\chi}^0(\xi)$ and $\tilde{\chi}(x, \xi)$ (symbols of $\tilde{\chi}^0$ and $\tilde{\chi}$) coincide with those in (2.1) and (2.2). Clearly \tilde{E} satisfies the Assumption (E) with $E_0 = I$, and $\tilde{\chi}^0$ satisfies the Assumption (F). Then we can replace E_0 with I without loss of generality, and hereafter we shall omit the sign " \sim " for simplicity.

Here we give some fundamental concepts which are important to develop the theory of first order symmetric hyperbolic

systems. (About these concepts refer to C.H. Wilcox [10] and K. Kikuchi [3]). Let $P_k^0(\xi)$ and $P_k(x, \xi)$ be orthogonal projections on the eigenspaces corresponding to $\lambda_k^0(\xi)$ and $\lambda_k(x, \xi)$, respectively. The properties of $\lambda_k^0(\xi)$ and $P_k^0(\xi)$ are given in [3, (2.5) ~ (2.13)]. $\lambda_k(x, \xi)$ and $P_k(x, \xi)$ also satisfy the same properties. Especially $\lambda_k(x, \xi)$'s are positively homogeneous of degree 1 with respect to ξ . The slowness surface for Λ^0 is defined by

$$S = \{ \xi \in \mathbb{R}^n; \det(I - \Lambda^0(\xi)) = 0 \}.$$

The condition 1) of the Assumption (F) is equivalent to the fact that S is bounded. We put

$$S_k = \{ \xi \in \mathbb{R}^n; \lambda_k^0(\xi) = 1 \}, \quad j = 1, 2, \dots, \rho,$$

and then $S = \bigcup_{k=1}^{\rho} S_k$. Let $z_S^{(1)}$ be a set of all algebraic singularities of S . $\lambda_k^0(\xi)$ may not be smooth on $z_S^{(1)}$. $z_S^{(2)}$ denotes a set defined in [3, page 579], which contains all points at which the Gaussian curvature vanishes. (For the definition of the Gaussian curvature see [4]). And we put

$$z_S = z_S^{(1)} \cup z_S^{(2)}.$$

In this paper we shall denote for $M \subset \mathbb{R}^n$

$$\bar{M} = \{ \xi = rs; s \in M \text{ and } r \in R \}.$$

For example $\bar{Z}_S = \{ \xi = rs; s \in Z_S \text{ and } r \in R \}$. (Note that \bar{Z}_S is a closed null set). The similar concepts are defined for Λ , for example S_x, Z_{S_x} and \bar{Z}_{S_x} .

About the eigenvalues and projections we have the following proposition.

Proposition 2.1 Under the Assumptions (E) and (F) the following facts hold:

i) There exists an $R > 0$ such that

$$\rho(x) = \rho_0 \quad \text{when } |x| > R.$$

ii) Let $K \subset R^n \setminus \bar{Z}_S$ be compact. Take R large enough. Then $\lambda_k(x, \xi)$ and $P_k(x, \xi)$ are smooth for $(x, \xi) \in \{|x| > R\} \times K$ and satisfies

$$|\partial_x^\alpha \partial_\xi^\beta (\lambda_k(x, \xi) - \lambda_k^0(\xi))| \leq C|x|^{-\delta - |\beta|}$$

for any $\alpha, \beta, |x| > R$ and $\xi \in K$.

$$\text{iii) } |\partial_x^\alpha \partial_\xi^\beta (P_k(x, \xi) - P_k^0(\xi))| \leq C|x|^{-\delta - |\beta|}$$

for any $\alpha, \beta, |x| > R$ and $\xi \in K$.

Here C is uniform for $(x, \xi) \in \{|x| > R\} \times K$.

Proof Since $\rho_0 = (m-d)/2$, the algebraic equation with respect to λ

$$\det(\lambda I - \Lambda^0(\xi)) = 0$$

does not have a double root other than $\lambda = 0$ when $\xi \notin \bar{z}_S^{(1)}$.

Then the continuity of the roots gives i) and

$$|\lambda_k(x, \xi) - \lambda_k^0(\xi)| \rightarrow 0 \quad \text{as} \quad |x| \rightarrow \infty$$

uniformly for ξ in a compact set $K \subset \mathbb{R}^n \setminus \bar{z}_S^{(1)}$. The rest of the assertions are easily proved by a straightforward calculation.

Q.E.D.

In the rest of this section we state a theorem of the stationary phase method on a hypersurface. It will play an important role in the argument of section 4.

Let $\{s_t\}_{t \geq T}$ be a family of smooth hypersurfaces. For each t a compact support smooth function $\mu_t(s)$ defined on s_t is given. Suppose that they satisfy

(i). For sufficiently small domain U of \mathbb{R}^n there exist smooth functions $\{\varphi_t\}$ such that

$$s_t \cap U = \{ \varphi_t(\xi) = 0 \}.$$

ii) There exists another smooth function φ such that

$$|\partial^\alpha \varphi_t - \partial^\alpha \varphi| \leq c_\alpha t^{-\delta} \quad \text{for } |\alpha| \geq 0.$$

iii) The Gaussian curvature $K_t(s)$ of S_t satisfies

$$0 < c_1 \leq |K_t(s)| \leq c_2 \quad \text{for } s \in \text{supp } \mu_t,$$

where c_1 and c_2 are constants independent of t .

iv) The derivatives of $\mu_t(s)$ are uniformly bounded with respect to t .

We consider here the integral

$$I_t(x) = \int_{S_t} e^{ixs} \mu_t(s) dS_t, \quad x \in \mathbb{R}^n$$
$$(dS_t = (2\pi)^{-(n-1)/2} ds_t).$$

Let $s_t^\gamma(\theta)$'s be the points where the exterior unit normal of S_t is θ , and let $\rho_t(\theta)$ be the number of $\{s_t^\gamma(\theta)\}$. Put

$$(2.3) \quad \psi_t^\pm(s) = \exp\{\pm(\pi i/4)(p_t^+(s) - p_t^-(s))\},$$

where $p_t^\pm(s)$ denotes the number of positive (negative) principal curvatures of S_t at s . Now the following fact holds.

Theorem 2.2 If $\{s_t\}_{t \geq T}$ and $\{\mu_t\}_{t \geq T}$ are given as above, then $I_t(x)$ satisfies the following expansion formula:

$$I_t(x) = \sum_{\gamma=1}^{\rho_t(\vartheta)} e^{ixs} \psi_t^+(s) \mu_t(s) |K_t(s)|^{-1/2} \Big|_{s=s_t^{\gamma(\vartheta)}} |x|^{-(n-1)/2}$$

$$+ \sum_{\gamma=1}^{\rho_t(-\vartheta)} e^{ixs} \psi_t^-(s) \mu_t(s) |K_t(s)|^{-1/2} \Big|_{s=s_t^{\gamma(-\vartheta)}} |x|^{-(n-1)/2}$$

$$+ \sum_{\gamma=1}^{\rho_t(\vartheta)} e^{ixs} C^+(t, s) |K_t(s)|^{-1/2} \psi_t^+(s) \Big|_{s=s_t^{\gamma(\vartheta)}} |x|^{-n/2}$$

$$+ \sum_{\gamma=1}^{\rho_t(-\vartheta)} e^{ixs} C^-(t, s) |K_t(s)|^{-1/2} \psi_t^-(s) \Big|_{s=s_t^{\gamma(-\vartheta)}} |x|^{-n/2}$$

$$+ q(t, x),$$

where $\vartheta = x/|x|$,

$$|q(t, x)| \leq C \sum_{|\alpha| \leq n+2} |\partial_\sigma^\alpha \mu_t(s(\sigma))| |x|^{-(n+2)/2}$$

and

$$|C^\pm(t, s)| \leq C \sum_{|\alpha|=0, 1} |\partial_\sigma^\alpha \mu_t(s(\sigma))|_{s(\sigma)=s}|$$

(σ is a local coordinate)

for some constant C independent of t .

The proof of Theorem 2.2 can be carried out in the same line of the proofs written in many literatures (see, for example, M. Matsumura [4, § 4 and §5]).

3 Definition and properties of $X(t)$

We denote the set $\{u \in C^\infty(\mathbb{R}^n); \hat{u} \in C_0^\infty(\mathbb{R}^n \setminus \bar{z}_S)\}$ by \mathcal{A}_S .

The unitary group $e^{-it\Lambda^0}$ is given for $u \in \mathcal{A}_S$ by

$$(3.1) \quad e^{-it\Lambda^0} u = P_0^0(D_x) u + \sum_{|k|=1}^{p_0} \int_{\mathbb{R}^n} e^{ix\xi - it\lambda_k^0(\xi)} P_k^0(\xi) \hat{u}(\xi) d\xi.$$

We must replace $e^{-it\Lambda^0}$ in the definition of the wave operator by another operator $X(t)$. So we shall seek $X(t)$ in the form

$$(3.2) \quad X(t)u = \sum_{|k|=1}^{p_0} \int_{\Omega_t} e^{ix\xi - iW_k(t, \xi)} P_k^0(\xi) \hat{u}(\xi) d\xi,$$

where $W_k(t, \xi)$ satisfies the Hamilton-Jacobi equation (1.4)

and Ω_t satisfies $\bigcup_t \Omega_t = \mathbb{R}^n \setminus \bar{z}_S$. The solutions of the Hamilton-Jacobi equations are obtained in a standard way. (Here we state only the case of $t \rightarrow +\infty$).

Lemma 3.1 Under the Assumptions (E) and (F), let $\{t_\ell\}_{\ell=1}^\infty$ be a given sequence with $t_1 < t_2 < \dots$ and $\lim_{\ell \rightarrow \infty} t_\ell = \infty$. Then

there exist a sequence of conic open sets $\{\Omega_\ell\}_{\ell=0}^\infty$ with $\Omega_0 \subset$

$\Omega_1 \subset \dots$ and $\bigcup_{\ell=0}^\infty \Omega_\ell = \mathbb{R}^n \setminus \bar{z}_S^{(1)}$, and C^∞ -functions $w_k(x, \xi)$

($|k| = 1, \dots, p_0$) defined on $\bigcup_{\ell=1}^\infty [t_\ell, \infty) \times \Omega_\ell$ having the

following properties:

i) $w_k(x, \xi)$ satisfies

$$(3.3) \quad \frac{\partial w_k}{\partial t} = \lambda_k(\nabla_\xi w_k, \xi) \quad \text{on} \quad \bigcup_{\ell=1}^{\infty} [t_\ell, \infty) \times \Omega_\ell.$$

For any compact set $K \subset \mathbb{R}^n \setminus \bar{Z}_S^{(1)}$

ii) $|\partial_\xi^\alpha w_k(t, \xi)| \leq C_\alpha t$ for $|\alpha| \geq 1$ and $\xi \in K$

iii) $|\partial_\xi^\alpha (t^{-1} w_k(t, \xi) - \lambda_k^0(\xi))| + |\partial_\xi^\alpha (\lambda_k(\nabla w_k, \xi) - \lambda_k^0(\xi))| \leq C_\alpha t^{-\delta}$ for $|\alpha| \geq 0$ and $\xi \in K$

iv) For $r > 0$ $w_k(t, r\xi) = r w_k(t, \xi)$

v) $w_k(t, -\xi) = -w_{-k}(t, \xi)$.

The proof of this lemma is similar to that of Lemma 3.8 of [2]. Assertions iv) and v) follow from the fact that $\lambda_k(x, \xi)$'s have the same properties.

From Lemma 3.1 ii) and iii) for $|\alpha| = 1$ we have

$$(3.4) \quad C't \leq |\nabla w_k(t, \xi)| \leq Ct.$$

Now we define $X(t)$ by (3.2) for $u \in \mathcal{X}_S$ with functions w_k constructed in Lemma 3.1 and domains Ω_t given by

$$(3.5) \quad \Omega_t = \Omega_\ell \quad \text{for} \quad t \in [t_\ell, t_{\ell+1}).$$

Let $I_{W_k}(t)$ be operators defined by

$$I_{W_k}(t)u = \int_{\Omega_t} e^{ix\xi - iW_k(t, \xi)} \hat{u}(\xi) d\xi \quad \text{for } u \in \mathcal{J}_S,$$

and we put $X_k(t) = I_{W_k}(t)P_k^0(D_x)$. Then $X(t) = \sum_k X_k(t)$, $Y_k(t)$ and

$Y(t)$ are other operators given by $Y_k(t) = I_{W_k}(t)P_k(\nabla W_k(D_x), D_x)$

and $Y(t) = \sum_k Y_k(t)$. Here note that from Proposition 2.1 and

Lemma 3.1 the matrices $P_k(\nabla W_k, \xi)$'s are smooth on $\text{supp } \hat{u}$ when t is large.

Let $\rho(r) \in C^\infty(\mathbb{R}^1)$ be a function satisfying

$$\rho(r) = \begin{cases} 1 & \text{if } r \geq 1 \\ 0 & \text{if } r \leq 1/2, \end{cases}$$

and we put

$$(3.6) \quad \chi(t, x) = \rho(2C_W|x|/t),$$

where C_W is a constant which satisfies

$$(3.7) \quad C_W^{-1} \leq |\nabla \lambda_k^0(\xi)| \leq C_W \quad \text{for } |k| = 1, 2, \dots, \rho_0.$$

The existence of such constant easily follows from 1) of

Assumption (F).

Now we have

Lemma 3.2 i) Let $u \in \mathcal{Q}_S$ be fixed. Take sufficiently large t . Then for any $s > 0$

$$\| (1 - \chi(t, x)) I_{W_k}(t) u \|_{L^2} \leq C_{s,u} t^{-s}.$$

ii) u and t are the same as above. Then for any $s > 0$

$$\| \langle x \rangle^{-s} I_{W_k}(t) u \|_{L^2} \leq C_{s,u} t^{-s}.$$

iii) $x_k(t) - y_k(t) \rightarrow 0$ strongly as $t \rightarrow \infty$.

Proof i) Lemma 3.1 iii) gives

$$|t^{-1} \nabla w_k(t, \xi)| \geq |\nabla \lambda_k^0(\xi)| - c_1 t^{-\delta},$$

and this implies

$$|\nabla_\xi(x\xi - w_k(t, \xi))| \geq (C_W^{-1}/4)t$$

when t is large and (t, x) is in the support of $1 - \chi(t, x)$. Then the assertion easily follows from the equality

$$e^{ix\xi - iW_k(t, \xi)} = -i \frac{\nabla_\xi(x\xi - w_k)}{|\nabla_\xi(x\xi - w_k)|^2} \nabla_\xi e^{ix\xi - iW_k(t, \xi)}$$

and integration by parts.

ii) The assertion follows from the equality

$$e^{-iW_k(t, \xi)} = -i|\nabla W_k|^{-2} \nabla W_k \cdot \nabla e^{-iW_k(t, \xi)},$$

the inequality (3.4) and integration by parts.

iii) The assertion follows from (3.4) and Proposition 2.1
iii).

Q.E.D.

It is easy to see that $X(t)$ is a partial isometry as an operator from \mathcal{H} ($= L^2(\mathbb{R}^n; \mathbb{C}^m)$) to \mathcal{H} with the domain \mathcal{A}_S .

Since \mathcal{A}_S is dense in \mathcal{H} , $X(t)$ can be extend to an operator whose domain is the whole space \mathcal{H} . We denote this operator also by $X(t)$.

4 The existence of the modified wave operators

The purpose of this section is to prove the main theorem:

Theorem 4.1 The modified wave operator

$$(4.1) \quad W_{\pm}^D = \text{s-lim}_{t \rightarrow \pm\infty} e^{it\Lambda} X(t)$$

exists, and it is a partially isometric operator with the intertwining property

$$e^{is\Lambda} W_{\pm}^D u = W_{\pm}^D e^{is\Lambda^0} u, \quad s \in \mathbb{R} \quad \text{and} \quad u \in \mathcal{H}_c(\Lambda^0).$$

Proof If the limit

$$W_{\pm}^D u = \text{s-lim}_{t \rightarrow \pm\infty} e^{it\Lambda} X(t) u$$

exists for a dense set of $u \in \mathcal{H}$, it exists for every $u \in \mathcal{H}$. Hence we may assume $u \in \mathcal{H}_S$. If t is so large that $\text{supp } \hat{u} \subset \Omega_t$, we may consider $\Omega_t = \mathbb{R}^n$ in (3.2). Then we shall take such t for any fixed u .

Here we consider only the case of W_+^D . (W_-^D can be treated in the same way).

In the same way as in the proof of intertwining property for the case of the Schrodinger operators (see [2]) we have

$$\lim_{t \rightarrow \infty} X(t+s)^* X(t) e^{-is\Lambda^0} u = (I - P_0^0(D_x)) u,$$

and the intertwining property has been proved. (Note that $P_c(\Lambda^0) = P_{ac}(\Lambda^0) = I - P_0^0(D_x)$).

Now we prove the existence of the limit (4.1). From Lemma 3.2 iii) the existence of $W_+^D u$ is equivalent to the existence of

$$(4.2) \quad s\text{-}\lim_{t \rightarrow \infty} e^{it\Lambda} y_k(t) u \quad (|k| = 1, 2, \dots, p_0).$$

When we have for some T

$$\int_T^\infty \left\| \frac{d}{dt} (e^{it\Lambda} y_k(t) u) \right\|_{L^2} dt < \infty,$$

we know that the limit (4.2) exists. Here note the equality

$$\begin{aligned} i^{-1} e^{-it\Lambda} \frac{d}{dt} (e^{it\Lambda} y_k(t) u) &= \Lambda y_k(t) u - i \frac{dy_k(t)}{dt} u \\ &= \int_{\mathbb{R}^n} e^{ix\xi - iW_k(t, \xi)} \left[(\Lambda(x, \xi) - \frac{\partial W_k}{\partial t}(t, \xi)) P_k(\nabla W_k, \xi) \hat{u}(\xi) \right. \\ &\quad \left. + R(t, x, \xi) \hat{u}(\xi) \right] d\xi, \end{aligned}$$

where

$$|R(t, x, \xi)| \leq C(\langle x \rangle^{-1-\delta} + t^{-1-\delta}).$$

From Lemma 3.2 ii) we easily have

$$\left\| \int_{\mathbb{R}^n} e^{ix\xi - iW_k(t, \xi)} R(t, x, \xi) \hat{u}(\xi) d\xi \right\|_{L^2} \leq C t^{-1-\delta}.$$

Let $\chi(t, x)$ be of (3.6). From Lemma 3.2 i) we have

$$\begin{aligned} & \left\| \int_{\mathbb{R}^n} e^{ix\xi - iW_k(t, \xi)} (1 - \chi(t, \xi)) \right. \\ & \quad \cdot \left. (\Lambda(x, \xi) - \frac{\partial W_k}{\partial t}(t, \xi)) P_k(\nabla W_k, x) \hat{u}(\xi) d\xi \right\| \leq C t^{-1-\delta}. \end{aligned}$$

Hence we have only to consider

$$\begin{aligned} (4.3) \quad I(t, x) = & \int_{\mathbb{R}^n} e^{ix\xi - iW_k(t, \xi)} \chi(t, x) \\ & \cdot (\Lambda(x, \xi) - \frac{\partial W_k}{\partial t}(t, \xi)) P_k(\nabla W_k, \xi) \hat{u}(\xi) d\xi. \end{aligned}$$

Our purpose is to prove that the L^2 -norm of $I(t, x)$ is integrable with respect to t .

Here we introduce a new surface which is like the slowness surface:

$$S_k(t) = \{\xi \in \Omega_t; t^{-1}W_k(t, \xi) = 1\} \quad (\Omega_t \text{ is of (3.5)}).$$

Put $K = \text{supp } \hat{u}$ and $\bar{K} = \{\xi = r\xi'; r \in R \text{ and } \xi' \in K\}$. It

follows from (3.4) that $S_k(t) \cap \bar{K}$ is a smooth hypersurface when we take sufficiently large t (depending on \bar{K}). We put $\Sigma_t = S_k(t) \cap \bar{K}$.

In (4.3) we make the change of variables $\xi \mapsto (r, s)$ ($r > 0, s \in \Sigma_t$) by $\xi = rs$. From Lemma 3.1 iv) and v) this transformation is surely one to one. We put

$$(4.4) \quad F_k(t, x, \xi) = \chi(t, x)(\Lambda(x, \xi) - \frac{\partial w_k}{\partial t}(t, \xi))\hat{p}_k(\nabla w_k, \xi).$$

In the same way as [10, §4] we have

$$\begin{aligned} I(t, x) &= \int_{-\infty}^{\infty} \left(\int_{\Sigma_t} e^{i(rxs - w_k(t, rs))} F_k(t, x, rs) \hat{u}(rs) |T_t(s)|^{-1} \right. \\ &\quad \left. \cdot dS(t) \right) |r|^{n-1} dr \\ &= \int_{-\infty}^{\infty} e^{-irt} r |r|^{n-1} \left(\int_{\Sigma_t} e^{irxs} F_k(t, x, s) \hat{u}(rs) |T_t(s)|^{-1} \right. \\ &\quad \left. \cdot dS(t) \right) dr, \end{aligned}$$

where $dS(t)$ is the surface element of Σ_t and

$$(4.5) \quad T_t(s) = t^{-1} \nabla w_k(t, s) \quad \text{for } s \in S_k(t).$$

Now we put

$$v_k(t, x, y, r) = \int_{\Sigma_t} e^{ixs} F_k(t, y, s) \hat{u}(rs) |T_t(s)|^{-1} dS(t).$$

Then

$$(4.6) \quad I(t, x) = \int_{-\infty}^{\infty} e^{-irt} v_k(t, rx, x, r) r |r|^{n-1} dr.$$

We shall use the stationary phase method (Theorem 2.2) for the integral on Σ_t . For any parametrization $s = s(\sigma)$ the gradient of the phase $x \cdot s(\sigma)$ with respect to σ satisfies $\vec{\nabla}_\sigma(x \cdot s) = \dot{x} \cdot \vec{\nabla}_\sigma s$. Here note that the column vector of $\vec{\nabla}_\sigma s(\sigma)$ construct the tangent plane of $S_k(t)$ at $s(\sigma)$. Hence the stationary points are the points s where x is normal to $S_k(t)$, that is, $N_t(s) = \pm \theta$ ($\theta = x/|x|$). We denote the Gaussian curvature of $S_k(t)$ by $K_t(s)$. Since $K \subset \mathbb{R}^n \setminus \bar{Z}_S$, it easily follows that $K_t(s)$ does not vanish for sufficiently large t in Σ_t . This fact implies that the Gauss map $N_t(s)$ on Σ_t is a C^∞ -injective map into S^{n-1} . We put $\omega_t = N_t(\Sigma_t)$. Then N_t is bijective from Σ_t to ω_t . We denote the inverse map of this map by s_t .

It is easy to see that the integral of v_k satisfies all assumptions of Theorem 2.2. Hence we have

$$(4.7) \quad v_k(t, x, y, r) = e^{ix \cdot s} F_k(t, y, s) \hat{u}(rs) |T_t(s)|^{-1} |K_t(s)|^{-1/2} \cdot \psi_t^+(s) \Big|_{s=s_t(\theta)} |x|^{-(n-1)/2} + e^{ix \cdot s} F_k(t, y, s) \hat{u}(rs) \cdot |T_t(s)|^{-1} |K_t(s)|^{-1/2} \psi_t^-(s) \Big|_{s=s_t(-\theta)} |x|^{-(n-1)/2}$$

$$\begin{aligned}
& + e^{ix \cdot s} C^+(t, y, r, s) |K_t(s)|^{-1/2} \psi_t^+(s) \Big|_{s=s_t(\vartheta)} |x|^{-n/2} \\
& + e^{ix \cdot s} C^-(t, y, r, s) |K_t(s)|^{-1/2} \psi_t^-(s) \Big|_{s=s_t(-\vartheta)} |x|^{-n/2} \\
& + q(t, x, y, r) \\
= & v_{0,+} + v_{0,-} + v_{1,+} + v_{1,-} + q,
\end{aligned}$$

where

$$|q(t, x, y, r)| \leq C(\langle y \rangle^{-\delta} + t^{-\delta}) |x|^{-(n+2)/2}$$

and

$$(4.8) \quad |C^\pm(t, y, r, s)| \leq C \sum_{|\alpha|=0,1} |\{\partial_\xi^\alpha (F_k \hat{u})\}_{\xi=s}|$$

with constants independent of t . It is easy to see that the supports of C^\pm and q with respect to r are compact in $\mathbb{R}^1 \setminus \{0\}$.

From (4.6) and (4.7)

$$I(t, x) = I_{0,+} + I_{0,-} + I_{1,+} + I_{1,-} + Q,$$

where

$$I_{j,\pm}(t, x) = \int_{-\infty}^{\infty} e^{-irt} v_{j,\pm}(t, rx, x, r) r |r|^{n-1} dr \quad (j = 0, 1)$$

and $Q(t, x)$ is defined similarly. Q satisfies the estimate

$$(4.9) \quad |Q(t, x)| \leq C|x|^{-(n+2)/2}(\langle x \rangle^{-\delta} + t^{-\delta}).$$

Now we put

$$(4.10) \quad \phi_{\pm}^0(\tau, s; x, t) = F_k(t, x, s) |T_t(s)|^{-1} |K_t(s)|^{-1/2} \cdot \int_{-\infty}^{\infty} e^{ir\tau} r |r|^{(n-1)/2} \hat{u}(rs) \psi_t^{\text{sign}(\pm r)}(s) dr$$

and

$$(4.11) \quad \phi_{\pm}^1(\tau, s; x, t) = |K_t(s)|^{-1/2} \int_{-\infty}^{\infty} e^{ir\tau} \cdot C^{\pm}(t, x, r, rs) |r|^{n/2-1} \psi_t^{\text{sign}(\pm r)}(s) dr.$$

(Recall that ψ_t^{\pm} is defined in (2.3)). Then $I_{j,\pm}$ are written as

$$I_{j,\pm}(t, x) = \begin{cases} |x|^{-(n-1+j)/2} \phi_{\pm}^j(x \cdot s_t(\pm \theta) - t, s_t(\pm \theta); x, t) & (\pm \theta \in \omega_t) \\ 0 & (\text{otherwise}). \end{cases}$$

It easily follows from Lemma 3.1 iv) and (4.5) that $T_t(s) = (s \cdot N_t(s))^{-1} N_t(s)$. This equality implies

$$x \cdot s_t(-\delta) - t = -|x|(-\delta) \cdot s_t(-\delta) - t = -|x| |T_t(s_t(-\delta))|^{-1} - t.$$

Hence we have

$$(4.12) \quad |x \cdot s_t(-\delta) - t| > c|x| + t.$$

On the other hand integration by parts gives

$$(4.13) \quad \Phi_-^j(\tau, s; x, t) = \{\dots\}(-1/i\tau)^2 \int_{-\infty}^{\infty} e^{i\tau r} (\partial/\partial r)^2 \{\dots\} dr.$$

Then from (4.12) and (4.13) $I_{j,-}(t, x)$ ($j = 0, 1$) satisfy the same estimate of (4.9). Now we write

$$(4.14) \quad I(t, x) = I_{0,+} + I_{1,+} + \tilde{Q},$$

where $\tilde{Q} = I_{0,-} + I_{1,-} + \tilde{Q}$. Since \tilde{Q} satisfies the same estimate of (4.9) and $|x| > (1/4C_W)t$ on the support of $\chi(t, x)$, we have

$$(4.15) \quad \|\tilde{Q}(t, \cdot)\|_{L^2} \leq ct^{-1-\delta}.$$

Next we calculate the L^2 -norm of $I_{j,+}$:

$$\begin{aligned} \|I_{j,+}(t, \cdot)\|_{L^2}^2 &= \int_0^{\infty} \int_{S^{n-1}} I_{j,+}(t, x)^* I_{j,+}(t, x) |x|^{n-1} \\ &\quad \cdot ds^{n-1} d|x| \end{aligned}$$

$$= \int_0^\infty \int_{\omega_t} \phi_+^j(x \cdot s_t(\vartheta) - t, s_t(\vartheta); x, t)^* \\ \cdot \phi_+^j(x \cdot s_t(\vartheta) - t, s_t(\vartheta); x, t) ds_{\vartheta}^{n-1} d|x|.$$

Now we make a change of variables $s = s_t(\vartheta)$ ($\vartheta \in \omega_t$). It is well-known that

$$ds^{n-1} = |K_t(s)| ds(t).$$

It is easy to see that $\vartheta \cdot s_t(\vartheta) = |T_t(s_t(\vartheta))|^{-1}$ and $x = |x|\vartheta = |x|N_t(s_t(\vartheta))$. Thus

$$||I_{j,+}(t, \cdot)||_{L^2}^2 = \int_0^\infty \int_{\omega_t} \phi_+^j(|x||T_t(s)|^{-1} - t, s; \\ |x|N_t(s), t)^* \phi_+^j(|x||T_t(s)|^{-1} - t, s; |x|N_t(s), t) \\ \cdot |K_t(s)| ds(t) d|x|.$$

Here we make another change of variable $|x| \rightarrow \tau$ by

$$|x||T_t(s)|^{-1} - t = \tau.$$

Then $d|x| = |T_t(s)|^{-1} d\tau$, and we have

$$(4.16) \quad ||I_{j,+}(t, \cdot)||_{L^2}^2 = \int_{-t}^\infty \int_{\Sigma_t} \phi_+^j(\tau, s; (\tau + t)T_t(s), t)^* \\ \cdot \phi_+^j(\tau, s; (\tau + t)T_t(s), t) d\tau ds.$$

$$\cdot \psi_+^j(\tau, s; (\tau + t)T_t(s), t) |T_t(s)| |K_t(s)| ds(t) d\tau.$$

First we consider the case of $j = 0$. Recall the definition (4.10) of ϕ_+^0 . We begin with F_k . From (3.3) and (4.4)

$$\begin{aligned} F_k(t, x, s) &= \chi(t, x)(\Lambda(x, s) - \lambda_k(\nabla w_k, s)) P_k(\nabla w_k, s) \\ &= \chi(t, x)(\Lambda(x, s) - \Lambda(\nabla w_k, s)) P_k(\nabla w_k, s) \\ &= \chi(t, x) \int_0^1 \sum_{\mu=1}^n (x - \nabla w_k)_\mu (\partial_{x_\mu} \Lambda)((x - \nabla w_k)_\theta + \\ &\quad \nabla w_k, s) d\theta \cdot P_k(\nabla w_k, s). \end{aligned}$$

From (4.5) we have

$$\begin{aligned} F_k(t, (\tau + t)T_t(s), s) &= \chi(t, (\tau + t)T_t(s)) \\ &\cdot \int_0^1 \sum_{\mu=1}^n \tau(T_t(s))_\mu (\partial_{x_\mu} \Lambda)(\tau T_t(s)_\theta + t T_t(s), s) d\theta \cdot P(t T_t(s), s). \end{aligned}$$

This gives

$$\begin{aligned} (4.17) \quad |F_k(t, (\tau + t)T_t(s), s)| & \\ &\leq C |\tau| \int_0^1 <|\tau\theta + t| |T_t(s)|>^{-1-\delta} d\theta \leq C' |\tau| <t>^{-1-\delta} (1 + \tau^2/2). \end{aligned}$$

Integration by parts gives for any non-negative integer ℓ

$$(4.18) \quad \tau^{\ell} (1 + \tau^2/2) \int_{-\infty}^{\infty} e^{ir\tau} r |r|^{(n-1)/2} \hat{u}(rs) \psi_t^{\text{sign}(\pm r)}(s) dr$$

$$= \int_{-\infty}^{\infty} e^{ir\tau} (-D_r)^{\ell} (1 + D_r^2/2) (r |r|^{(n-1)/2} \hat{u}(rs) \psi_t^{\text{sign}(\pm r)}(s)) dr.$$

Clearly the right hand side of (4.18) is bounded. Then we easily have for any $s > 0$

$$(4.19) \quad |\tau(1 + \tau^2/2) \int_{-\infty}^{\infty} e^{ir\tau} r |r|^{(n-1)/2} \cdot \hat{u}(rs) \psi_t^{\text{sign}(\pm r)}(s) dr| \leq C_s |\tau|^{-s}.$$

If we estimate the right hand side of (4.10) by using (4.17) and (4.19) with $s = (1 - \delta)/2$, we obtain

$$|\phi_+^0(\tau, s; (\tau + t)T_t(s), t)| \leq C t^{-1-\delta} |\tau|^{-(1-\delta)/2}.$$

Thus from (4.16) we have

$$(4.20) \quad \|I_{0,+}(t, \cdot)\|_{L^2}^2 \leq C t^{-2(1+\delta/2)}.$$

For the case of $j = 1$ note the inequality (4.8) and the equality

$$\partial_{\xi_\mu} F_k(t, x, \xi)$$

$$= x(t, x) [\{(\partial_{\xi_\mu} \Lambda)(x, \xi) - (\partial_{\xi_\mu} \Lambda)(\nabla w_k, \xi)$$

$$- 32 -$$

$$\begin{aligned}
& - \sum_{v=1}^n \partial_{\xi_\mu} \partial_{\xi_v} w_k \cdot (\partial_{x_v} \Lambda) (\nabla w_k, \xi) \} p_k (\nabla w_k, \xi) \\
& + (\Lambda(x, \xi) - \Lambda(\nabla w_k, \xi)) \partial_{\xi_\mu} (p_k (\nabla w_k, \xi)),
\end{aligned}$$

and in the same way as in the proof of the case of $j = 0$ we have

$$(4.21) \quad \left\| I_{1,+}(t, \cdot) \right\|_{L^2}^2 \leq C' t^{-2(1 + \delta/2)}.$$

Thus the proof is complete.

Q.E.D.

Reference

- [1] V. Georgiev and P. Stefanov, Existence of the wave operators for dissipative hyperbolic systems with variable multiplicities, preprint.
- [2] L. Hörmander, The existence of wave operators in scattering theory, Math. Z., 146, (1976), 69 - 91.
- [3] K. Kikuchi, Asymptotic behavior at infinity of the Green function of a class of systems including wave propagation in crystals, Osaka J. Math., 22 (1985), 575 - 635.
- [4] M. Matsumura, Asymptotic behavior at infinity for Green's functions of first order systems with characteristics of nonuniform multiplicity, Publ. RIMS, Kyoto Univ., 12 (1976), 317 - 377.
- [5] H. Tamura, The principle of limiting absorption for uniformly propagative systems with perturbations of long-range class, Nagoya Math. J., 82 (1981), 141 - 174.
- [6] H. Tamura, The principle of limiting absorption for propagative systems in crystal optics with perturbations of long-range class, Nagoya Math. J., 84 (1981), 169 - 193.
- [7] H. Tamura, The principle of limiting absorption and decay of local energy for the linearized equation of magnetofgasdynamics, Nagoya Math. J., 89 (1983), 13 - 45.
- [8] H. Tamura, Spectral and scattering theory for symmetric

systems of non constant deficit, J. Funct. Anal. 67
(1986), 73 - 104.

[9] R. Weder, Spectral analysis of strongly propagative systems, J. Reine Angewandte Math., 354 (1984), 95 - 122.

[10] C.H. Wilcox, Asymptotic wave functions and energy distributions in strongly propagative anisotropic media, J. Math. pures et appl., 57 (1978), 275 - 321.

[11] C.H. Wilcox, Electromagnetic signal propagation in crystals, Appl. Anal., 8 (1978), 83 - 94.

Department of Mathematics

Faculty of Sience

Osaka University

560 Osaka Toyonaka, Japan