



Title	Asymptotic behavior of solutions of $x'' = e^{\alpha \lambda t} x^{1+\alpha}$ where $-1 < \alpha < 0$
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# ASYMPTOTIC BEHAVIOR OF SOLUTIONS OF $x'' = e^{\alpha \lambda t} x^{1+\alpha}$ WHERE $-1 < \alpha < 0$

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## 1. Introduction

Let us consider second order nonlinear differential equations

$$(1.1)_{\pm} \quad x'' = \pm t^{\beta} x^{1+\alpha}$$

$$(1.2)_{\pm} \quad x'' = \pm e^{\sigma t} x^{1+\alpha}$$

where  $' = d/dt$ , the double signs correspond in the same order in every equation and  $\alpha, \beta, \sigma$  are parameters. Using Chapter 7 of [1], we can state value of solving these. First these can be derived from an important second order nonlinear differential equation

$$(1.3) \quad \frac{d}{dt} \left( t^{\rho} \frac{du}{dt} \right) \pm t^{\sigma} u^n = 0,$$

$\rho, \sigma, n$  being parameters, which contains the Emden equation of astrophysics and the Fermi-Thomas equation of atomic physics and so has several interesting physical applications. Second (1.3) is mathematically interesting, because (1.3) is nontrivial, nonlinear and has a large class of solutions whose behavior can be ascertained with astonishing accuracy nevertheless these cannot be generally obtained explicitly. In addition (1.1) $_{\pm}$ , (1.2) $_{\pm}$  are examples of differential equations positive radial solutions of a nonlinear elliptic partial differential equation satisfy (cf. [17]).

Actually many authors have considered (1.1) $_{\pm}$ , (1.2) $_{\pm}$  and (1.3) in more general form in [2], [5] through [9], [13], [20] and so on. In these papers they mainly discussed asymptotic behavior of the solution continuable to  $\infty$ . On the other hand, initial value problems of (1.1) $_{+}$ , (1.2) $_{+}$ , (1.2) $_{-}$  and (1.1) $_{-}$  were considered in [10], [11], in [14], [16], in [15], [16] and in [17] respectively in case of  $\alpha > 0$  and asymptotic behavior of all the solutions was studied.

In the case  $\alpha < 0$ , the initial value problems of (1.1) $_{\pm}$ , (1.2) $_{\pm}$  are not considered yet, while in [5], [8], [20] etc. this case was already considered for differential equations with more general form than (1.1) $_{\pm}$ , (1.2) $_{\pm}$  and for the solutions continuable to  $\infty$ . So in this paper, we shall consider (1.2) $_{+}$  where  $-1 < \alpha < 0$  as a first step. Since

it is convenient to put  $\sigma = \alpha\lambda$ , the equation to be considered has the form

$$(E) \quad x'' = e^{\alpha\lambda t} x^{1+\alpha}$$

where  $-1 < \alpha < 0$ ,  $\lambda < 0$ . It is noteworthy that the case  $\lambda > 0$  can be reduced to our case if we replace  $t$  with  $-t$ . A domain where (E) will be considered is given as

$$(1.4) \quad -\infty < t < \infty, \quad 0 < x < \infty.$$

Notice that if  $p$  is a positive number and  $r$  is a real number, then throughout this paper  $p^r$  always takes its positive branch.

The initial condition given to (E) is

$$(I) \quad x(t_0) = a, \quad x'(t_0) = b$$

where

$$-\infty < t_0 < \infty, \quad a > 0, \quad -\infty < b < \infty.$$

$t_0$  will be fixed arbitrary and  $a$  suitably. For every  $b$ , we shall study asymptotic behavior of all solutions of an initial value problem (E), (I).

For this, we shall use the method which follows the arguments originally done in [10], [11] and applied in [14] through [19]. In this method, we adopt a transformation

$$(T) \quad y = \psi(t)^{-\alpha} \phi(t)^\alpha, \quad z = y'$$

where  $\psi(t) = \lambda^{2/\alpha} e^{-\lambda t}$  ( $\lambda^{2/\alpha} = (\lambda^2)^{1/\alpha}$ ) is a particular solution of (E) and  $\phi(t)$  is a solution of (E). This transforms (E) into a first order rational differential equation

$$(R) \quad \frac{dz}{dy} = \frac{(\alpha - 1)z^2 + 2\alpha\lambda yz - \alpha^2\lambda^2(y^2 - y^3)}{\alpha yz}.$$

Using a parameter  $s$ , we rewrite this as a 2-dimensional dynamical system

$$(D) \quad \begin{aligned} \frac{dy}{ds} &= \alpha yz \\ \frac{dz}{ds} &= (\alpha - 1)z^2 + 2\alpha\lambda yz - \alpha^2\lambda^2(y^2 - y^3). \end{aligned}$$

Graphs of solutions of (R) have the same shape as orbits of solutions of (D) except on the  $y$  and  $z$  axes. Since from (1.4) only a positive solution of (E) is considered,  $y$  is always positive. Finally we note that

$$(1.5) \quad z = \alpha y \left( \lambda + \frac{\phi'(t)}{\phi(t)} \right)$$

got from (T) will be often used.

## 2. On solutions of (E) obtained from orbits of (D) connecting its two singularities

The singularities of (D) are points  $(0, 0)$ ,  $(1, 0)$ . From orbits of (D) connecting these points, we get the following through (T):

**Theorem I.** *Let  $\phi(t)$  be a solution of the initial value problem (E), (I) and suppose*

$$0 < a < \psi(t_0).$$

*Then there exist  $b_1$ ,  $b_2$  ( $b_1 < b_2$ ) such that*

(i) *if  $b = b_1$ ,  $\phi(t)$  is defined for  $-\infty < t < \infty$  so that in the neighborhood of  $t = \infty$ ,  $\phi(t)$  is represented as*

$$(2.1) \quad \phi(t) = \lambda^{2/\alpha} e^{-\lambda t} \left[ 1 + Ce^{(\mu_1/\alpha)t} + \sum_{n=2}^{\infty} a_n \left\{ Ce^{(\mu_1/\alpha)t} \right\}^n \right]$$

*where  $C$ ,  $a_n$  are constants and*

$$\mu_1 = (1 + \sqrt{1 + \alpha})\alpha\lambda,$$

*and as  $t \rightarrow -\infty$ ,*

$$(2.2) \quad \phi(t) = ct + d + \frac{(ct)^{1+\alpha}}{\alpha^2 \lambda^2} e^{\alpha \lambda t} (1 + o(1))$$

*where  $c (< 0)$ ,  $d$  are constants,*

(ii) *if  $b = b_2$ ,  $\phi(t)$  is defined for  $-\infty < t < \infty$  so that in the neighborhood of  $t = \infty$ ,  $\phi(t)$  is represented as*

$$(2.3) \quad \phi(t) = \lambda^{2/\alpha} e^{-\lambda t} \left[ 1 + (A + gB^N t) e^{(\mu_1/\alpha)t} + Be^{(\mu_2/\alpha)t} + \sum_{m+n \geq 2} a_{mn} \left\{ (A + gB^N t) e^{(\mu_1/\alpha)t} \right\}^m \left\{ Be^{(\mu_2/\alpha)t} \right\}^n \right]$$

*where  $A$ ,  $B$ ,  $g$ ,  $a_{mn}$  are constants,  $B \neq 0$ ,  $N = \mu_1/\mu_2$ ,*

$$\mu_2 = (1 - \sqrt{1 + \alpha})\alpha\lambda$$

*and  $g \neq 0$  only if  $\mu_1/\mu_2$  is a positive integer, and in the neighborhood of  $t = -\infty$ ,  $\phi(t)$  is represented as*

$$(2.4) \quad \phi(t) = \lambda^{2/\alpha} C^{1/\alpha} \left\{ 1 + \sum_{n=1}^{\infty} a_n (Ce^{\alpha \lambda t})^n \right\}$$

where  $C, a_n$  are constants,

(iii) if  $b_1 < b < b_2$ ,  $\phi(t)$  is defined for  $-\infty < t < \infty$  and represented as (2.3) in the neighborhood of  $t = \infty$  and (2.2) in the neighborhood of  $t = -\infty$ .

Let us start the proof. First we consider  $(1, 0)$ . Putting

$$y = 1 + \eta, \quad z = \zeta,$$

we get from (D)

$$(2.5) \quad \begin{aligned} \frac{d\eta}{ds} &= \alpha\zeta + \dots \\ \frac{d\zeta}{ds} &= \alpha^2\lambda^2\eta + 2\alpha\lambda\zeta + \dots \end{aligned}$$

where  $\dots$  denotes terms whose degrees are greater than the previous terms. The coefficient matrix of the linear terms of (2.5) has eigenvalues  $\mu_1, \mu_2$ . Since  $-1 < \alpha < 0$ , we get

$$(2.6) \quad \mu_1 > \mu_2 > 0.$$

A linear transformation

$$\begin{pmatrix} \eta \\ \zeta \end{pmatrix} = \begin{pmatrix} \alpha & \alpha \\ \mu_1 & \mu_2 \end{pmatrix} \begin{pmatrix} \tilde{\eta} \\ \tilde{\zeta} \end{pmatrix}$$

transforms (2.5) into

$$\frac{d\tilde{\eta}}{ds} = \mu_1\tilde{\eta} + \dots, \quad \frac{d\tilde{\zeta}}{ds} = \mu_2\tilde{\zeta} + \dots.$$

Owing to Theorem A of [3] and its proof, a transformation

$$\tilde{\eta} = w_1 + \dots, \quad \tilde{\zeta} = w_2 + \dots$$

holomorphic in the neighborhood of  $(w_1, w_2) = (0, 0)$  transforms this into

$$(2.7) \quad \frac{dw_1}{ds} = \mu_1 w_1 + gw_2^N, \quad \frac{dw_2}{ds} = \mu_2 w_2$$

where  $g$  is a constant such that  $g \neq 0$  only if  $\mu_1/\mu_2$  is a positive integer and  $N = \mu_1/\mu_2$ . That is

$$(2.8) \quad w_1 = (A + gB^N s)e^{\mu_1 s}, \quad w_2 = Be^{\mu_2 s}$$

where  $A, B$  are arbitrary constants. Therefore an arbitrary solution of (D) converging to  $(1, 0)$  is given as

$$(2.9) \quad y = 1 + \alpha w_1 + \alpha w_2 + \sum_{m+n \geq 2} a_{mn} w_1^m w_2^n$$

$$(2.10) \quad z = \mu_1 w_1 + \mu_2 w_2 + \sum_{m+n \geq 2} b_{mn} w_1^m w_2^n$$

in the neighborhood of  $(w_1, w_2) = (0, 0)$ , namely of  $s = -\infty$ . From (2.7), (2.9) we get

$$z = y' = \{\alpha(\mu_1 w_1 + g w_2^N) + \alpha \mu_2 w_2 + \dots\} s'.$$

If we fix  $s, s'$  and vary  $A, B$  and compare this with (2.10), then since

$$\frac{\partial(w_1, w_2)}{\partial(A, B)} = e^{(\mu_1 + \mu_2)s} \neq 0$$

and hence  $w_1, w_2$  can attain arbitrary values, we conclude

$$s = \frac{t}{\alpha} + C$$

where  $C$  is an arbitrary constant. So if we replace  $(A + gB^N C)e^{\mu_1 C}$ ,  $g/\alpha$  and  $Be^{\mu_2 C}$  with  $A, g$  and  $B$  respectively, then since

$$B^N e^{\mu_1 C} = (Be^{\mu_2 C})^N$$

in case of  $g \neq 0$ , we get

$$y = 1 + \alpha(A + gB^N t)e^{(\mu_1/\alpha)t} + \alpha B e^{(\mu_2/\alpha)t} + \sum_{m+n \geq 2} a_{mn}((A + gB^N t)e^{(\mu_1/\alpha)t})^m (Be^{(\mu_2/\alpha)t})^n.$$

Because  $t \rightarrow \infty$  as  $s \rightarrow -\infty$ , this is valid in the neighborhood of  $t = \infty$ . Consequently we get the following through (T):

**Lemma 2.1.** *From (2.9), (2.10) we obtain a solution  $\phi(t)$  of (E) with the form (2.3) valid in the neighborhood of  $t = \infty$ .*

Note that  $B$  of (2.3) is not always nonzero in Lemma 2.1. If  $B \neq 0$ , then from (2.9), (2.10) we have

$$(2.11) \quad \lim_{s \rightarrow -\infty} \frac{z}{y - 1} = \frac{\mu_2}{\alpha}.$$

However if  $B = 0$ , then we get

$$(2.12) \quad \lim_{s \rightarrow -\infty} \frac{z}{y - 1} = \frac{\mu_1}{\alpha}.$$

Therefore there exists uniquely a solution of (R) satisfying (2.12). Indeed from (2.9), (2.10) and  $B = 0$  this is represented as

$$(2.13) \quad z = \frac{\mu_1}{\alpha}(y - 1) + \sum_{n=2}^{\infty} c_n(y - 1)^n$$

in the neighborhood of  $y = 1$ . Conversely we get (2.10) from substituting (2.9) into (2.13). In this way we get (2.9), (2.10) from (2.13). Hence if we take  $A = C$ ,  $B = 0$  in Lemma 2.1, we obtain

**Lemma 2.2.** *From (2.13) we get a solution  $\phi(t)$  of (E) with the form (2.1) in the neighborhood of  $t = \infty$ .*

For further discussions, we examine the sign of  $dz/ds$  in (D). So we put  $dz/ds = 0$ . Then we get  $z = Z_{\pm}(y)$  where

$$Z_{\pm}(y) = \frac{\alpha\lambda}{1-\alpha}y \left\{ 1 \pm \sqrt{\alpha - (\alpha - 1)y} \right\}.$$

$Z_{\pm}(y)$  are defined for  $y \geq -\alpha/(1-\alpha)$  and

$$Z_{\pm} \left( -\frac{\alpha}{1-\alpha} \right) = -\frac{\alpha^2\lambda}{(1-\alpha)^2}, \quad Z_+(y) > Z_-(y) \quad \text{for } y > -\frac{\alpha}{1-\alpha}.$$

$Z_+(y)$  is monotone increasing. If  $\alpha \leq -1/3$ , then  $Z_-(y)$  is monotone decreasing. If  $-1/3 < \alpha < 0$ , then  $Z_-(y)$  has a minimal value and a maximal value and is monotone increasing in the interval between these extrema and monotone decreasing outside this interval. Moreover we get

$$\lim_{y \rightarrow \infty} Z_{\pm}(y) = \pm\infty.$$

**Lemma 2.3.** *If  $y < -\alpha/(1-\alpha)$  or  $z > Z_+(y)$  or  $z < Z_-(y)$ , then*

$$\frac{dz}{ds} < 0,$$

*if  $z = Z_{\pm}(y)$ , then*

$$\frac{dz}{ds} = 0,$$

*and if  $Z_-(y) < z < Z_+(y)$ , then*

$$\frac{dz}{ds} > 0.$$

Let us consider how the solution (2.9), (2.10) of (D) behaves asymptotically as  $s \rightarrow \infty$ , if  $B = 0$ . For this we introduce a curve

$$z = f(y) = \alpha \lambda (y - y^2).$$

Owing to the proof of Proposition 4 of [14], we get

$$(2.14) \quad \frac{d}{ds}(z - f(y)) = (\alpha + 1)\alpha^2 \lambda^2 y^3(1 - y) > 0$$

when a solution  $(y, z)$  of (D) passes the curve  $z = f(y)$  where  $0 < y < 1$ . Moreover we have

$$\frac{\mu_1}{\alpha} = -\frac{\alpha \lambda}{1 - \sqrt{1 + \alpha}} < -\alpha \lambda = f'(1).$$

It follows from (2.12) and  $Z'_-(1) = -\alpha \lambda / 2$  that (2.13) lies between  $z = f(y)$  and  $z = Z_+(y)$  in the  $yz$  plane. Hence from Lemma 2.3 and Poincaré-Bendixon's theorem, (2.9), (2.10) where  $B = 0$  tend to  $(0, 0)$  as  $s \rightarrow \infty$ . Since (2.13) is obtained uniquely from (2.9), (2.10), we get

**Lemma 2.4.** *There exists the unique solution  $z = z_1(y)$  of (R) defined for  $0 \leq y \leq 1$  such that (2.13) holds. Moreover we obtain*

$$\lim_{y \rightarrow +0} z_1(y) = \lim_{y \rightarrow 1-0} z_1(y) = 0.$$

Next let us consider the singularity  $(0, 0)$  of (D). As preparation of this, we show

**Lemma 2.5.** *If there exists a solution  $y = y(x)$  of a Briot-Bouquet differential equation*

$$(2.15) \quad x \frac{dy}{dx} = f(x, y)$$

where  $f(x, y)$  is a holomorphic function in the neighborhood of  $(x, y) = (0, 0)$  with the form

$$(2.16) \quad f(x, y) = \lambda y + ax + \sum_{j+k \geq 2} a_{jk} x^j y^k, \quad \lambda < 0$$

and if the accumulation points of  $y(x)$  contain 0 as  $x$  tends to 0 with bounded  $\arg x$ , then  $y(x)$  is the unique holomorphic solution.

Proof. It is known that if  $\lambda$  is not a positive integer, then there exists the unique holomorphic solution  $y = h(x)$  of (2.15) such that  $h(0) = 0$ . So we put  $z = y - h(x)$ .

Then we get

$$x \frac{dz}{dx} = zg(x, z)$$

where  $g(x, z)$  is a holomorphic function in the neighborhood of  $(x, z) = (0, 0)$  and  $g(0, 0) = \lambda$ . Hence it suffices to show that the solution  $y(x)$  of (2.15) satisfying the assumption of this lemma is identically zero, when we get

$$(2.17) \quad f(x, y) = \lambda y \left( 1 + \sum_{j+k \geq 1} a_{jk} x^j y^k \right)$$

instead of (2.16).

Suppose the contrary. Then  $y(x)$  is not identically zero. Now if  $c(\neq 0)$  is an accumulation point of  $y(x)$  as  $x$  tends to 0 with bounded  $\arg x$ , then there exists a compact neighborhood  $U$  of 0 such that  $c \notin U$ . Since  $y(x)$  intersects the boundary of  $U$  infinitely many times and the boundary of  $U$  is compact,  $y(x)$  has its accumulation point on the boundary of  $U$  as  $x$  tends to 0 with bounded  $\arg x$ . Since  $U$  can be taken sufficiently small, we may suppose that  $c$  is so small that  $f(x, c) \neq 0$ . Hence we get from (2.15)

$$\frac{dx}{dy} = \frac{x}{f(x, y)}$$

which implies the contradiction  $x \equiv 0$  from Painlevé's theorem (cf. Theorem 3.2.1 of [4]) and the uniqueness theorem. Consequently  $y(x)$  converges to 0 as  $x$  tends to 0 with bounded  $\arg x$ .

Substituting  $y = y(x)$  into (2.15), we have from (2.17)

$$x \frac{dy}{dx} = \lambda y(1 + o(1))$$

as  $x$  tends to 0 with bounded  $\arg x$ . Therefore for some  $x_0$  in the neighborhood of 0 and  $y_0 = y(x_0)$  we get

$$\int_{y_0}^y \frac{dy}{y} = \lambda \int_{x_0}^x \frac{1 + o(1)}{x} dx$$

where  $|x| < |x_0|$  and  $y = y(x)$ . From Cauchy's theorem we take  $\Gamma$  as a path of integration of the righthand side so that if  $x \in \Gamma$ , then  $|x|$  and  $\arg x$  vary monotonously along  $\Gamma$ . Then we get

$$\int_{x_0}^x \frac{1 + o(1)}{x} dx = \int_{x_0}^x \frac{dx}{x} (1 + o(1))$$

and thus

$$\log \frac{y}{y_0} = \lambda \left( \log \frac{x}{x_0} \right) (1 + o(1))$$

whose real parts deduce a contradiction as  $x$  tends to 0 with bounded  $\arg x$ . Hence  $y(x)$  is identically zero.  $\square$

**Lemma 2.6.** *Let  $z = z(y)$  be a solution of (R) such that*

$$\lim_{y \rightarrow 0} z(y) = 0.$$

*Then we have*

$$(2.18) \quad \lim_{y \rightarrow 0} y^{-1} z = \alpha \lambda.$$

Proof. Since Lemma 5 of [14] is valid also in our case, we get

$$\lim_{y \rightarrow 0} y^{-1} z = \alpha \lambda, \quad \pm\infty.$$

If  $\lim_{y \rightarrow 0} y^{-1} z = \pm\infty$ , we put  $w = yz^{-1}$ . Then as in [14] we obtain

$$y \frac{dw}{dy} = \frac{w}{\alpha} - 2\lambda w^2 + \alpha \lambda^2 (1 - y) w^3.$$

Since  $1/\alpha < 0$ , it follows from Lemma 2.5 that this has the unique holomorphic solution  $w \equiv 0$ . This is a contradiction and (2.18) is valid.

Concerning a solution  $z = z(y)$  of (R) converging to 0 as  $y \rightarrow 0$  and satisfying (2.18), we shall get some lemmas.  $\square$

**Lemma 2.7.** *There exists uniquely a solution  $z = z_2(y)$  of (R) such that (2.18) holds and*

$$(2.19) \quad \lim_{y \rightarrow 0} y^{-1} v = \lambda$$

where

$$v = y^{-1} z - \lambda.$$

Furthermore in the neighborhood of  $y = 0$  we get

$$(2.20) \quad z_2(y) = \alpha \lambda y + \lambda y^2 + \dots$$

Proof. Putting  $w = y^{-1} v - \lambda$ , we get

$$w \rightarrow 0 \quad \text{as } y \rightarrow 0$$

from (2.19) and

$$\frac{dw}{dy} = \frac{-(\alpha+1)y(w+\lambda)^2 - \alpha^2\lambda w}{\alpha y(yw + \lambda y + \alpha\lambda)}.$$

This is (27) of [14]. Hence if we follow discussion of [14] after (27), then the proof is completed.

Using (2.20) and (T), we have □

**Lemma 2.8.** *From  $z_2(y)$  we obtain a solution  $\phi(t)$  of (E) with the form (2.4) in the neighborhood of  $t = -\infty$ .*

Here in the same way as in [16] we conclude

**Lemma 2.9.** *A solution of (R) satisfying (2.18) and not satisfying (2.19) is given as*

$$(2.21) \quad z = \alpha\lambda y + \frac{\alpha^2\lambda y}{\log|y|} \left( 1 + O\left( \frac{\log|\log|y||}{\log|y|} \right) \right) \quad \text{as } y \rightarrow 0.$$

In the neighborhood of  $y = 0$ , we get

$$(2.21) > \alpha\lambda y > f(y) > z_2(y)$$

since  $-1 < \alpha < 0$ . However a solution of (D) satisfies

$$\frac{dy}{ds} = 0, \quad \frac{dz}{ds} = -\alpha^2\lambda^2(y^2 - y^3) < 0$$

on the segment  $0 < y < 1$ ,  $z = 0$ . Therefore from (2.14) and Poincaré-Bendixon's theorem we have

$$(y, z_2(y)) \rightarrow (1, 0) \quad \text{as } s \rightarrow -\infty.$$

Namely  $z_2(y)$  is defined for  $0 \leq y \leq 1$  and

$$\lim_{y \rightarrow 1-0} z_2(y) = 0.$$

Because only  $z_1(y)$  satisfies (2.12) and only  $z_2(y)$  satisfies (2.18) and (2.19), we conclude

**Lemma 2.10.** *On  $0 < y < 1$ , we get*

$$z_1(y) > z_2(y) > 0$$

and  $z_1(y)$  is represented as (2.21) in the neighborhood of  $y = 0$  and  $z_2(y)$  as (2.9), (2.10) where  $B \neq 0$  in the neighborhood of  $y = 1$ .

Now we continue the argument of [16] used for obtaining Lemma 2.9 and have

$$\phi(t) \sim ct \quad \text{as } t \rightarrow -\infty.$$

Moreover integrating both sides of (E) twice, we conclude

**Lemma 2.11.** *From (2.21) we get a solution  $\phi(t)$  of (E) with the form (2.2) as  $t \rightarrow -\infty$ .*

Finally, note that from (T) and (1.5) the initial condition (I) gives an initial condition

$$(2.22) \quad z(y_0) = z_0$$

to (R). Here

$$(2.23) \quad y_0 = \psi(t_0)^{-\alpha} a^\alpha, \quad z_0 = \alpha y_0 \left( \lambda + \frac{b}{a} \right).$$

Since  $z_1(y)$ ,  $z_2(y)$  are defined for  $0 < y < 1$ , we take

$$0 < y_0 < 1.$$

This is equivalent to

$$(2.24) \quad 0 < a < \psi(t_0).$$

Fix  $t_0$  arbitrarily and  $a$  so as to satisfy (2.24). Then varying  $b$ ,  $y_0$  is fixed and  $z_0$  varies. Suppose that from (R) and (2.22) we get  $z_1(y)$  if  $b = b_1$  and  $z_2(y)$  if  $b = b_2$ . Then from Lemma 2.10, we obtain  $b_1 < b_2$  since  $z_0$  is monotone decreasing in  $b$ . If  $b_1 < b < b_2$ , then we have a solution  $z(y)$  of (R) and (2.22) such that

$$z_1(y) > z(y) > z_2(y).$$

Thus the unique existence of  $z_1(y)$  and  $z_2(y)$  implies

**Lemma 2.12.** *Let  $z(y)$  be a solution of (R), (2.22) with  $b_1 < b < b_2$ . Then  $z(y)$  is defined for  $0 \leq y \leq 1$  and represented as (2.21) in the neighborhood of  $y = 0$  and as (2.9), (2.10) where  $B \neq 0$  in the neighborhood of  $y = 1$ .*

Consequently in Theorem I, (i) follows from Lemmas 2.2, 2.4, 2.10, 2.11 and (ii) from Lemmas 2.1, 2.8, 2.10 and (iii) from Lemmas 2.1, 2.11, 2.12.

### 3. Preliminaries for the further discussions

Before considering the cases not treated yet, we need the following discussions.

If we put  $y = 1/\eta$  in (R), we get

$$(3.1) \quad \frac{dz}{d\eta} = -\frac{(\alpha - 1)\eta^3 z^2 + 2\alpha\lambda\eta^2 z - \alpha^2\lambda^2(\eta - 1)}{\alpha\eta^4 z}.$$

If we put  $z = 1/\zeta$  in (R), then

$$(3.2) \quad \frac{d\zeta}{dy} = -\frac{(\alpha - 1)\zeta + 2\alpha\lambda y\zeta^2 - \alpha^2\lambda^2(y^2 - y^3)\zeta^3}{\alpha y}.$$

Furthermore if we put  $y = 1/\eta$ ,  $z = 1/\zeta$ , then we have

$$(3.3) \quad \frac{d\zeta}{d\eta} = \frac{(\alpha - 1)\eta^3\zeta + 2\alpha\lambda\eta^2\zeta^2 - \alpha^2\lambda^2(\eta - 1)\zeta^3}{\alpha\eta^4}.$$

Moreover if we put

$$w = \eta^{-3/2}\zeta, \quad \xi = \eta^{1/2},$$

then we obtain a Briot-Bouquet differential equation

$$(3.4) \quad \xi \frac{dw}{d\xi} = -\frac{\alpha + 2}{\alpha}w + 4\lambda\xi w^2 - \alpha\lambda^2(\xi^2 - 1)w^3.$$

Let  $z = z(y)$  be a solution of (R). Then through (T) we obtain a solution  $x = \phi(t)$  of (E). If  $(\omega_-, \omega_+)$  denotes a domain of  $\phi(t)$  and if  $y$  is a function obtained from  $\phi(t)$  through (T), then we get

**Lemma 3.1.**  $y \rightarrow \infty$  as  $t \rightarrow \omega_{\pm}$  imply that  $\omega_{\pm}$  are finite respectively.

Proof. If a solution  $z$  of (R) is bounded, then (3.1) implies a contradiction  $\eta \equiv 0$ . Hence  $z$  is unbounded.

So we consider (3.4). If  $\xi = 0$ , then the righthand side of (3.4) vanishes if and only if  $w = 0, \pm\rho$  where

$$\rho = \frac{1}{\alpha\lambda} \sqrt{\frac{\alpha + 2}{2}}.$$

Here let  $c$  be an accumulation point of a solution  $w$  of (3.4) as  $\xi \rightarrow 0$ , namely  $y \rightarrow \infty$ . Suppose that  $c \neq 0, \pm\rho, \pm\infty$ . Then from (3.4) we get a contradiction  $\xi \equiv 0$ . Hence  $c = 0, \pm\rho, \pm\infty$ .

If  $c = 0$ , then we get from (3.4)

$$w = C\xi^{-(1+2/\alpha)} \left[ 1 + \sum_{m+n \geq 1} a_{mn} \xi^m \{C\xi^{-(1+2/\alpha)}\}^n \right]$$

where  $C$  is an arbitrary constant and the power series converges in the neighborhood of  $\xi = 0$ , since  $-(\alpha + 2)/\alpha > 0$  and the righthand side of (3.4) is divisible by  $w$ . Returning the variables, we have

$$(3.5) \quad y^{1/\alpha} \left\{ 1 + \sum_{m+n \geq 1} b_{mn} y^{-m/2 + ((\alpha+2)/2\alpha)n} \right\} = \frac{t - \omega_-}{\alpha C} \quad \text{or} \quad \frac{t - \omega_+}{\alpha C}$$

where  $-\infty < \omega_- < \omega_+ < \infty$ .

If  $c = \pm\rho$ , we put  $\theta = w - c$ . Then we have

$$\xi \frac{d\theta}{d\xi} = \frac{2(\alpha+2)}{\alpha^2 \lambda} \xi + \frac{2(\alpha+2)}{\alpha} \theta + \dots$$

Since  $2(\alpha+2)/\alpha < 0$  and  $\theta$  is real so that  $\arg \theta$  is bounded, Lemma 2.5 implies that  $\theta$  is holomorphic and represented as

$$\theta = \sum_{n=1}^{\infty} a_n \xi^n.$$

Here, return the variables. Then we get

$$(3.6) \quad -2cy^{-1/2} - \sum_{n=1}^{\infty} \frac{2a_n}{n+1} y^{-(n+1)/2} = t - \omega_- \quad \text{or} \quad t - \omega_+$$

where  $-\infty < \omega_- < \omega_+ < \infty$ .

Now we suppose  $c = \pm\infty$ . Putting  $w = 1/\theta$ , we have

$$\frac{d\xi}{d\theta} = \frac{\alpha \xi \theta}{(\alpha+2)\theta^2 - 4\alpha\lambda\xi\theta + 2\alpha^2\lambda^2(\xi^2 - 1)}.$$

These imply a contradiction  $\xi \equiv 0$ . Consequently  $c \neq \pm\infty$ .  $\square$

**Corollary 3.2.** *If  $c = 0$ , we get*

$$(3.7) \quad \phi(t) = \frac{\lambda^{2/\alpha} e^{-\lambda\omega_-}}{\alpha C} (t - \omega_-) \times \left\{ 1 + \sum_{l+m+n \geq 1} d_{lmn} (t - \omega_-)^l (t - \omega_-)^{-\alpha m/2} (t - \omega_-)^{(\alpha+2)n/2} \right\}$$

in the neighborhood of  $t = \omega_-$  and

$$(3.8) \quad \phi(t) = \frac{\lambda^{2/\alpha} e^{-\lambda\omega_+}}{\alpha C} (\omega_+ - t) \\ \times \left\{ 1 + \sum_{l+m+n \geq 1} d_{lmn} (\omega_+ - t)^l (\omega_+ - t)^{-\alpha m/2} (\omega_+ - t)^{(\alpha+2)n/2} \right\}$$

in the neighborhood of  $t = \omega_+$  where  $C$  is an arbitrary constant and  $d_{lmn}$  are constants. Moreover if  $c = \pm\rho$ , we get

$$(3.9) \quad \phi(t) = \left\{ \frac{2(\alpha+2)}{\alpha^2} \right\}^{1/\alpha} e^{-\lambda\omega_-} (t - \omega_-)^{-2/\alpha} \left\{ 1 + \sum_{n=1}^{\infty} c_n (t - \omega_-)^n \right\}$$

in the neighborhood of  $t = \omega_-$  and

$$(3.10) \quad \phi(t) = \left\{ \frac{2(\alpha+2)}{\alpha^2} \right\}^{1/\alpha} e^{-\lambda\omega_+} (\omega_+ - t)^{-2/\alpha} \left\{ 1 + \sum_{n=1}^{\infty} c_n (\omega_+ - t)^n \right\}$$

in the neighborhood of  $t = \omega_+$  where  $c_n$  are constants.

Furthermore if  $c \neq 0, \pm\rho$ , then the solution  $\phi(t)$  of (E) cannot be obtained.

Proof. In the proof of Lemma 3.1, we get  $c = 0, \pm\rho$ . If  $c = 0$ , then we put

$$y = \frac{1}{\eta}, \quad \eta^{1/2} = \xi$$

and get from (3.5)

$$\xi^{-2/\alpha} \left\{ 1 + \sum_{m+n \geq 1} \tilde{a}_{mn} \xi^{m-((\alpha+2)/\alpha)n} \right\} = \frac{t - \omega_-}{\alpha C} \quad \text{or} \quad \frac{t - \omega_+}{\alpha C}.$$

Here if we put

$$\theta = \xi^{-2/\alpha}, \quad \tau = \frac{t - \omega_-}{\alpha C} \quad \text{or} \quad \frac{t - \omega_+}{\alpha C}$$

then

$$\theta \left\{ 1 + \sum_{m+n \geq 1} \tilde{a}_{mn} \theta^{-(\alpha/2)m} \theta^{((\alpha+2)/2)n} \right\} = \tau.$$

Hence we have

$$(3.11) \quad \tau^{-\alpha/2} = \theta^{-\alpha/2} \left\{ 1 + \sum_{m+n \geq 1} b_{mn} \theta^{-(\alpha/2)m} \theta^{((\alpha+2)/2)n} \right\}$$

$$(3.12) \quad \tau^{(\alpha+2)/2} = \theta^{(\alpha+2)/2} \left\{ 1 + \sum_{m+n \geq 1} c_{mn} \theta^{-(\alpha/2)m} \theta^{((\alpha+2)/2)/n} \right\}.$$

Applying the inverse function theorem to (3.11) and (3.12), we obtain

$$\theta^{-\alpha/2} = \tau^{-\alpha/2} \left\{ 1 + \sum_{m+n \geq 1} \hat{a}_{mn} \tau^{-(\alpha/2)m} \tau^{((\alpha+2)/2)n} \right\}.$$

Therefore we get

$$y^{1/\alpha} = \tau \left\{ 1 + \sum_{m+n \geq 1} \hat{b}_{mn} \tau^{-(\alpha/2)m} \tau^{((\alpha+2)/2)n} \right\}$$

and so (3.7) and (3.8) through (T).

If  $c = \pm \rho$ , then from (3.6) and (T) we have (3.9) and (3.10).  $\square$

#### 4. On the other solutions of (E)

Let  $z_3(y)$  be a solution of (R) which exists in  $y > 1$  and is represented as (2.13) in the neighborhood of  $y = 1$ . Moreover, suppose that we get  $z_3(y)$  as a solution of the initial value problem (R), (2.22), if  $a > \psi(t_0)$  and  $b = b_3$ . Then if  $\phi(t)$  denotes a solution of the initial value problem (E), (I) as in Section 2, we have

**Theorem II.** *If  $0 < a < \psi(t_0)$  and  $b > b_2$ , then  $\phi(t)$  is defined for  $\omega_- < t < \infty$  where  $\omega_- > -\infty$  and represented as (3.7) in the neighborhood of  $t = \omega_-$  and (2.3) where  $B \neq 0$  in the neighborhood of  $t = \infty$ .*

*If  $a = \psi(t_0)$  and  $b = -a\lambda$ , then  $\phi(t) \equiv \psi(t)$  and if  $a = \psi(t_0)$  and  $b > -a\lambda$ , then the conclusion of the case  $0 < a < \psi(t_0)$ ,  $b > b_2$  follows.*

*If  $a > \psi(t_0)$ , then there exists  $b_3$  such that*

- (i) *if  $b = b_3$ ,  $\phi(t)$  is defined for  $\omega_- < t < \infty$  where  $\omega_- > -\infty$  and represented as (3.7) in the neighborhood of  $t = \omega_-$  and (2.1) in the neighborhood of  $t = \infty$ ,*
- (ii) *if  $b > b_3$ , the conclusion of the case  $0 < a < \psi(t_0)$ ,  $b > b_2$  follows.*

For starting the proof, recall (2.14). Then

$$(4.1) \quad \frac{d}{ds}(z - f(y)) = (\alpha + 1)\alpha^2 \lambda^2 y^3 (1 - y) < 0,$$

when a solution  $(y, z)$  of (D) passes the curve  $z = f(y)$  where  $y > 1$ . Therefore

$$z_3(y) < f(y),$$

since  $f'(1) > \mu_1/\alpha$ . Thus there exists  $y_1$  ( $1 < y_1 \leq \infty$ ) such that

$$\lim_{y \rightarrow y_1} z_3(y) = -\infty.$$

However if  $y_1$  is finite, then putting  $z = 1/\zeta$  we obtain (3.2) and a contradiction  $\zeta \equiv 0$ . Consequently we conclude

$$y_1 = \infty, \quad \lim_{y \rightarrow \infty} z_3(y) = -\infty.$$

Furthermore we suppose

$$b > b_2 \quad \text{or} \quad b \geq b_3.$$

If  $z_0 > 0$ , then the solution  $z_+(y)$  of (R), (2.22) satisfies

$$(4.2) \quad 0 \leq z_+(y) < z_2(y).$$

Moreover in the  $yz$  plane,  $z_+(y)$  connects at some point  $(\tilde{y}, 0)$  with a solution  $z_-(y)$  of (R) satisfying

$$(4.3) \quad z_-(y) \leq 0 \quad \text{if} \quad \tilde{y} < y \leq 1, \quad z_-(y) < z_3(y) \leq 0 \quad \text{if} \quad y > 1.$$

On the other hand, if  $z_0 < 0$ , then the solution  $z_-(y)$  of (R), (2.22) satisfies (4.3) and connects with a solution  $z_+(y)$  of (R) satisfying (4.2) at  $(\tilde{y}, 0)$ . If  $z_0 = 0$ , then the solution of (R), (2.22) is given as  $z_+(y)$  and  $z_-(y)$  which satisfy (4.2) and (4.3) respectively and connect mutually at  $(y_0, 0)$ . So let  $z(y)$  be a many-valued function such that

$$(4.4) \quad z(y) = z_+(y) \quad \text{if} \quad z(y) \geq 0, \quad z(y) = z_-(y) \quad \text{if} \quad z(y) \leq 0.$$

Then the same discussion as was done for  $z_3(y)$  shows

$$(4.5) \quad \lim_{y \rightarrow \infty} z(y) = -\infty.$$

Here we state Lemma 4 of [15] as follows:

**Lemma 4.1.** *Let  $z_{\pm}(y)$  be solutions of (R) such that*

$$z_+(\tilde{y}) = z_-(\tilde{y}) = 0$$

*for some  $\tilde{y}$  and*

$$z_+(y) > 0, \quad z_-(y) < 0 \quad \text{for} \quad y \neq \tilde{y}$$

(i) If  $z_0 > 0$  and  $z_+(y)$  satisfies (R), (2.22) and  $y(t)$  is a solution of

$$\frac{dy}{dt} = z_+(y), \quad y(t_0) = y_0,$$

then there exists  $t_1$  such that

$$\lim_{t \rightarrow t_1+0} y(t) = \tilde{y}$$

and  $y(t)$  can be continued in the interval  $t < t_1$  uniquely by

$$\frac{dy}{dt} = z_-(y), \quad y(t_1) = \tilde{y}$$

(ii) If  $z_0 < 0$ , we get the similar conclusion.

(iii) If  $z_0 = 0$  and  $z_{\pm}(y)$  satisfy (R), (2.22), then  $y(t)$  can be defined uniquely by

$$\begin{aligned} \frac{dy}{dt} &= z_+(y) \quad \text{if } t > t_0, & \frac{dy}{dt} &= 0 \quad \text{if } t = t_0, \\ \frac{dy}{dt} &= z_-(y) \quad \text{if } t < t_0, & y(t_0) &= y_0. \end{aligned}$$

Proof. If  $(y, z) = (p(t), q(t))$  is a solution of

$$\begin{aligned} \frac{dy}{dt} &= z \\ \frac{dz}{dt} &= \frac{(\alpha - 1)z^2 + 2\alpha\lambda yz - \alpha^2\lambda^2(y^2 - y^3)}{\alpha y} \\ y(t_0) &= y_0, \quad z(t_0) = z_0, \end{aligned} \tag{4.6}$$

then it suffices to put  $y(t) = p(t)$ .

Return to our discussion. Then for  $z(y)$  defined as (4.4) we get  $y(t)$  and a solution  $\phi(t)$  of (E) through Lemma 4.1 and (T). Recall that  $(\omega_-, \omega_+)$  denotes a domain of  $\phi(t)$ .

Since  $(y(t), z(y(t)))$  is a solution of (4.6), we have

$$\lim_{t \rightarrow \omega_-} y(t) = \infty$$

from (4.2), (4.3), (4.4), (4.5) and Poincaré-Bendixon's theorem. Hence Lemma 3.1 implies  $\omega_- > -\infty$ . Because  $z_3(y) < f(y)$ , we obtain

$$\lim_{y \rightarrow \infty} y^{-3/2} z(y) = -\infty.$$

Hence if we put

$$y = \frac{1}{\eta}, \quad z = \frac{1}{\zeta}, \quad w = \eta^{-3/2} \zeta,$$

then we get

$$\lim_{\eta \rightarrow 0} w = 0$$

and a solution  $\phi(t)$  of (E) represented as (3.7) from Corollary 3.2.

Let us now consider the case  $t \rightarrow \omega_+$ . Then if  $b = b_3$ , we have  $\phi(t)$  represented as (2.1) from Lemma 2.2, since  $z_3(y)$  is represented as (2.13). In this case, we get  $\omega_+ = \infty$ . Moreover if  $b > b_2$  or  $b > b_3$ , then since  $(y(t), z(y(t)))$  is a solution of (4.6), Poincaré-Bendixon's theorem implies

$$y \rightarrow 1, \quad z_+(y) \rightarrow 0 \quad \text{as } t \rightarrow \omega_+.$$

Hence  $y, z(y)$  are given as (2.9), (2.10) and from Lemma 2.1 we get  $\phi(t)$  represented as (2.3) where  $B \neq 0$  and  $\omega_+ = \infty$ . Now the proof of Theorem II is completed.  $\square$

Next, suppose

$$b < b_1 \quad \text{if } 0 < a < \psi(t_0), \quad b < -a\lambda \quad \text{if } a = \psi(t_0), \quad b < b_3 \quad \text{if } a > \psi(t_0).$$

Then the solution  $z(y)$  of (R), (2.22) satisfies

$$(4.7) \quad z(y) > z_1(y), \quad z(y) > z_3(y).$$

In this case, we get

**Theorem III.** *If  $a > \psi(t_0)$ , then there exist  $b_4, b_5$  ( $b_5 < b_4 < b_3$ ) such that*

- (i) *if  $b_4 < b < b_3$ , then  $\phi(t)$  is defined for  $\omega_- < t < \infty$  where  $\omega_- > -\infty$  and represented as (3.7) in the neighborhood of  $t = \omega_-$  and (2.3) where  $B \neq 0$  in the neighborhood of  $t = \infty$ ,*
- (ii) *if  $b = b_4$ , then  $\phi(t)$  is defined for  $\omega_- < t < \infty$  where  $\omega_- > -\infty$  and represented as (3.9) in the neighborhood of  $t = \omega_-$  and (2.3) where  $B \neq 0$  in the neighborhood of  $t = \infty$ ,*
- (iii) *if  $b_5 < b < b_4$ , then  $\phi(t)$  is defined for  $-\infty < t < \infty$  and represented as (2.2) as  $t \rightarrow -\infty$  and (2.3) where  $B \neq 0$  in the neighborhood of  $t = \infty$ ,*
- (iv) *if  $b = b_5$ , then  $\phi(t)$  is defined for  $-\infty < t < \omega_+$  where  $\omega_+ < \infty$  and represented as (2.2) as  $t \rightarrow -\infty$  and (3.10) in the neighborhood of  $t = \omega_+$ ,*
- (v) *if  $b < b_5$ , then  $\phi(t)$  is defined for  $-\infty < t < \omega_+$  where  $\omega_+ < \infty$  and represented as (2.2) as  $t \rightarrow -\infty$  and (3.8) in the neighborhood of  $t = \omega_+$ .*

*If  $a = \psi(t_0)$ , then there exists  $b_5 (< -a\lambda)$  such that replacing  $b_4$  with  $-a\lambda$  we get (iii), (iv), (v) and if  $0 < a < \psi(t_0)$ , then replacing  $-a\lambda$  with  $b_1$  the conclusion of the case  $a = \psi(t_0)$  follows.*

Proof. Now in (R) we put

$$(4.8) \quad y^{-1/2} = \eta, \quad z^{-1} = \eta^3(-\rho + u).$$

Then from (21) of [14] we have

$$(4.9) \quad \eta \frac{du}{d\eta} = \frac{2(\alpha+2)}{\alpha^2 \lambda} \eta + \left(2 + \frac{4}{\alpha}\right) u + \dots$$

In the proof of Lemma 3.1, we put

$$(4.10) \quad y = \frac{1}{\eta}, \quad z = \frac{1}{\zeta}, \quad \xi = \eta^{1/2}, \quad w = \eta^{-3/2} \zeta, \quad \theta = w - c$$

where  $c = \pm\rho$  and obtained the differential equation similar to (4.9). Since  $2 + 4/\alpha < -2$ , Lemma 2.5 implies that there exists the unique solution  $u(\eta)$  of (4.9) such that  $u(0) = 0$ . Moreover  $u(\eta)$  is holomorphic in the neighborhood of  $\eta = 0$ . Hence we get a solution of (R) such as

$$(4.11) \quad z = y^{3/2} \left( -\rho^{-1} + \sum_{n=1}^{\infty} \tilde{a}_n y^{-n/2} \right).$$

Since  $y \rightarrow \infty$  as  $\eta \rightarrow 0$ , this is valid in the neighborhood of  $y = \infty$ . Moreover from the uniqueness of  $u(\eta)$ , (4.11) is uniquely determined. So we denote (4.11) as  $z_4(y)$ . Owing to

$$\lim_{y \rightarrow \infty} \frac{z_4(y)}{f(y)} = 0$$

and (4.1), we have

$$(4.12) \quad z_3(y) < f(y) < z_4(y).$$

Furthermore through (4.11) and (T) we get a solution  $\phi(t)$  of (E) and  $\omega_- > -\infty$  from Lemma 3.1. Using the notation of (4.10), we obtain

$$w \rightarrow c\#(c = -\rho) \quad \text{as} \quad \eta \rightarrow 0.$$

Hence from Corollary 3.2,  $\phi(t)$  is represented as (3.8) in the neighborhood of  $t = \omega_-$ .

Moreover since

$$\lim_{y \rightarrow \infty} z_4(y) = -\infty,$$

we get

$$\lim_{y \rightarrow 1+0} z_4(y) = 0$$

from Lemma 2.3 and (4.12). Hence we may suppose that  $z_4(y)$  is obtained from (R), (2.22) if  $a > \psi(t_0)$  and  $b = b_4$ . Furthermore from (4.12) we have  $b_3 > b_4$ .

If  $b_4 < b < b_3$ , then the solution  $z(y)$  of (R), (2.22) satisfies

$$(4.13) \quad z_3(y) < z(y) < z_4(y).$$

Hence we obtain

$$\lim_{y \rightarrow \infty} z(y) = -\infty.$$

Therefore defining  $y$  and  $\phi(t)$  through (T) and noting that  $(y, z(y))$  is a solution of (4.6), we get an alternative as  $t \rightarrow \omega_-$  from Lemma 3.1 as follows:

$$(4.14) \quad \omega_- > -\infty, \quad \lim_{t \rightarrow \omega_-} \phi(t) = 0$$

$$(4.15) \quad \omega_- > -\infty, \quad \lim_{t \rightarrow \omega_-} \phi(t) = \infty.$$

In case of (4.15) we get a contradiction

$$\lim_{t \rightarrow \omega_-} y = \lim_{t \rightarrow \omega_-} \lambda^{-2} e^{\alpha \lambda t} \phi(t)^\alpha = 0.$$

Next we consider the case (4.14). For this we use

$$(4.16) \quad \lim_{t \rightarrow \omega_-} y^{-3/2} z = \lim_{t \rightarrow \omega_-} \frac{\alpha \phi'(t)}{y^{1/2} \phi(t)}.$$

Since  $\phi''(t) > 0$ ,  $\phi(\omega_-) = 0$ , we obtain

$$0 \leq \phi'(\omega_-) < \infty.$$

In the case  $0 < \phi'(\omega_-) < \infty$ , we get

$$\lim_{t \rightarrow \omega_-} \frac{\phi'(t)^2}{y \phi(t)^2} = \lim_{t \rightarrow \omega_-} \frac{\phi'(t)^2}{\lambda^{-2} e^{\alpha \lambda t} \phi(t)^{\alpha+2}} = \infty.$$

Therefore from (4.16) we have

$$\lim_{t \rightarrow \omega_-} w = \lim_{t \rightarrow \omega_-} y^{3/2} z^{-1} = 0.$$

This implies that  $\phi(t)$  is represented as (3.7) from Corollary 3.2. On the other hand, in the case  $\phi'(\omega_-) = 0$  l'Hospital's theorem implies

$$\lim_{t \rightarrow \omega_-} \frac{\phi'(t)^2}{y \phi(t)^2} = \frac{2\lambda^2}{\alpha+2}$$

and from (4.16) we obtain

$$\lim_{t \rightarrow \omega_-} y^{-3/2} z = -\rho^{-1}.$$

Hence if we put  $u = \theta$  where  $\theta$  is defined as (4.10), then  $u$  is a solution of (4.9) with  $u(0) = 0$ . Since  $u$  exists uniquely, we get a contradiction  $z(y) \equiv z_4(y)$ .

Suppose  $b_4 \leq b < b_3$ . Then we have

$$\begin{aligned} z_3(y) &< z(y) \leq z_4(y), \\ y \rightarrow 1, \quad z(y) \rightarrow 0, \quad \frac{z(y)}{y-1} &\rightarrow \frac{\mu_2}{\alpha} \quad \text{as } t \rightarrow \omega_+, \end{aligned}$$

since only  $z_3(y)$  satisfies

$$\frac{z(y)}{y-1} \rightarrow \frac{\mu_1}{\alpha} \quad \text{as } y \rightarrow 1.$$

Therefore  $y, z(y)$  are represented as (2.9), (2.10) and from Lemma 2.1 we obtain a solution  $\phi(t)$  of (E) expressed as (2.3) where  $B \neq 0$ . Moreover  $\omega_+ = \infty$ . Now we conclude (i), (ii) of Theorem III.

Next, suppose  $b < b_4$ . Then we get

$$(4.17) \quad z(y) > z_4(y).$$

Here we consider the case  $t \rightarrow \omega_-$ . If  $y \rightarrow \infty, z(y) \rightarrow -\infty$  as  $t \rightarrow \omega_-$ , then in the neighborhood of  $y = \infty$

$$y^{-3/2} z(y) < 0.$$

So if we put

$$(4.18) \quad \eta = \frac{1}{y}, \quad \zeta = \frac{1}{z}, \quad w = \eta^{-3/2} \zeta, \quad \xi = \eta^{1/2},$$

then we have (3.4). Supposing that  $c$  is an accumulation point of  $w$  as  $\xi \rightarrow 0$ , we obtain

$$c \leq -\rho$$

from (4.17). However if  $c = -\rho$ , then we conclude a contradiction  $z(y) \equiv z_4(y)$ . Hence we get

$$c < -\rho.$$

From Corollary 3.2, this implies that there exists no solution of (3.4) whose accumulation points contain  $c$  as  $y \rightarrow \infty$  and that

$$y \rightarrow \infty, \quad z(y) \rightarrow -\infty$$

does not occur.

Therefore from Lemma 2.3, we have  $z(y) > 0$  or  $z(y)$  becomes a many-valued function such that

$$z(y) = z_+(y) \quad \text{if } z(y) \geq 0, \quad z(y) = z_-(y) \quad \text{if } z(y) \leq 0$$

where for some  $\tilde{y}$ ,  $z_-(y)$  is defined on  $1 \leq y \leq \tilde{y}$  and  $z_+(y)$  on  $0 \leq y \leq \tilde{y}$  so that

$$z_+(0) = z_+(\tilde{y}) = z_-(1) = z_-(\tilde{y}) = 0, \quad z_+(y) \geq 0, \quad z_-(y) \leq 0.$$

Indeed if  $z \rightarrow \gamma$  ( $\gamma \neq \pm\infty$ ) as  $y \rightarrow \infty$ , then from (3.1) we get a contradiction  $\eta = 1/y \equiv 0$ .

Since  $z_{\pm}(y)$  just defined satisfy the assumption of Lemma 4.1, we define  $y(t)$  as in Lemma 4.1. If  $(\omega_-, \omega_+)$  denotes a domain of  $y(t)$ , then we have

$$\lim_{t \rightarrow \omega_-} (y(t), z(y(t))) = (0, 0),$$

since  $(y(t), z(y(t)))$  satisfies (4.6). Hence it follows from Lemma 2.9 that (2.21) is obtained for  $(y(t), z(y(t)))$ . Therefore from Lemma 2.11 we get  $\omega_- = -\infty$  and  $\phi(t)$  is represented as (2.2) as  $t \rightarrow -\infty$ . Similarly in the case  $z(y) > 0$  we have

$$\lim_{y \rightarrow 0} z(y) = 0$$

from Lemma 2.3 and hence (2.2) as  $t \rightarrow -\infty$ .

Next we consider the case  $t \rightarrow \omega_+$ . Then there exist the following possibilities:

$$(4.19) \quad \omega_+ < \infty, \quad \lim_{t \rightarrow \omega_+} \phi(t) = 0$$

$$(4.20) \quad \omega_+ < \infty, \quad \lim_{t \rightarrow \omega_+} \phi(t) = \infty$$

$$(4.21) \quad \omega_+ = \infty, \quad \lim_{t \rightarrow \omega_+} \phi(t) = 0$$

$$(4.22) \quad \omega_+ = \infty, \quad 0 < \lim_{t \rightarrow \omega_+} \phi(t) < \infty$$

$$(4.23) \quad \omega_+ = \infty, \quad \lim_{t \rightarrow \omega_+} \phi(t) = \infty.$$

Here we define  $y$  and  $z$  through (T).

In the case (4.19) we get

$$\lim_{t \rightarrow \omega_+} y = \infty.$$

Suppose that

$$z \rightarrow \gamma (\gamma \neq \pm\infty) \quad \text{as } t \rightarrow \omega_+.$$

Then from (3.1) we have a contradiction. Therefore

$$z \rightarrow \infty \quad \text{as } t \rightarrow \omega_+$$

since if  $z < 0$ , then  $dy/dt < 0$  and  $y \rightarrow \infty$  as  $t \rightarrow \omega_+$  is impossible. Now we use (4.18) and transform (R) into (3.4). If  $c$  denotes an accumulation point of a solution  $w$  of (3.4), then it follows from Corollary 3.2 that in the neighborhood of  $t = \omega_+$  we get a solution  $\phi(t)$  of (E) represented as (3.8) for  $c = 0$  and (3.10) for  $c = \rho$ . Moreover since we get  $z > 0$  and  $c \geq 0$ , we do not obtain a solution of (E) for  $c \neq 0, \rho$ . As is shown in the proof of Lemma 3.1 and Corollary 3.2, (3.10) is obtained from the unique holomorphic solution

$$w = \rho + \sum_{n=1}^{\infty} a_n \xi^n$$

of (3.4), namely from a solution

$$z = \rho^{-1} y^{3/2} \left( 1 + \sum_{n=1}^{\infty} \tilde{a}_n y^{-n/2} \right)$$

of (R). Existence of this is unique and so we denote this as  $z_5(y)$ . Furthermore from the unique existence of  $z_5(y)$ , existence of (3.10) is also unique. So for (3.10) we put

$$\phi'(t_0) = b_5.$$

If (3.8) is got from a solution  $z(y)$  of (R), then from  $0 < \rho$ , we get

$$z_5(y) < z(y).$$

Therefore if we put

$$\phi'(t_0) = b$$

for (3.8), then from (2.23) we have

$$b_5 > b.$$

If the case (4.20), we obtain

$$\lim_{t \rightarrow \omega_+} y = 0$$

which is impossible. Moreover if the cases (4.21) and (4.22) occur, we get

$$\lim_{t \rightarrow \omega_+} y = \infty.$$

This contradicts Lemma 3.1.

Finally we suppose (4.23). Then if  $\phi'(t)$  is bounded as  $t \rightarrow \omega_+$ , we have

$$\lim_{t \rightarrow \omega_+} \frac{\phi(t)}{e^{-\lambda t}} = \lim_{t \rightarrow \omega_+} \frac{\phi'(t)}{-\lambda e^{-\lambda t}} = 0$$

and

$$\lim_{t \rightarrow \omega_+} y = \lim_{t \rightarrow \omega_+} \lambda^{-2} \left( \frac{\phi(t)}{e^{-\lambda t}} \right)^\alpha = \infty.$$

This implies a contradiction

$$\omega_+ < \infty.$$

If  $\phi'(t)$  is unbounded as  $t \rightarrow \omega_+$ , then from l'Hospital's theorem we get

$$\lim_{t \rightarrow \omega_+} \frac{\phi(t)}{e^{-\lambda t}} = \lim_{t \rightarrow \omega_+} \frac{\phi''(t)}{\lambda^2 e^{-\lambda t}} = \lim_{t \rightarrow \omega_+} \frac{1}{\lambda^2} \left( \frac{\phi(t)}{e^{-\lambda t}} \right)^{1+\alpha}.$$

On the other hand, since the orbit of the solution  $(y, z)$  of (D) cannot cross the  $y$  axis twice and  $z(=dy/dt)$  does not vanish twice,

$$\lim_{t \rightarrow \omega_+} y$$

exists and hence

$$\lim_{t \rightarrow \omega_+} \frac{\phi(t)}{e^{-\lambda t}}$$

does. Therefore this is equal to 0,  $\lambda^{2/\alpha}$ ,  $\infty$  and we get

$$\lim_{t \rightarrow \omega_+} y = \infty, 1, 0$$

respectively. However

$$\lim_{t \rightarrow \omega_+} y = \infty, 0$$

cannot occur as above. Thus we have

$$\lim_{t \rightarrow \omega_+} y = 1.$$

This occurs only if the orbit  $z = z(y)$  of (D) gets into the region where  $dz/ds > 0$  in the  $yz$  plane (cf. Lemma 2.3). Hence we obtain

$$z(y) < z_5(y), \quad b_5 < b(= \phi'(t_0))$$

and (2.3) where  $B \neq 0$  from  $z(y)$  as above.

Now the case

$$0 < a < \psi(t_0), \quad b < b_1 \text{ or } a = \psi(t_0), \quad b < -a\lambda$$

is left. However from (3.2) the solution  $z(y)$  of (R), (2.22) cannot diverge to  $\infty$  as  $y$  tends to a finite value. Moreover if  $z(y)$  converges to 0 as  $y \rightarrow 1 - 0$ , then we get the alternative of (2.11) and (2.12) and therefore a contradiction

$$z(y) \leq z_1(y).$$

Thus the present case is reduced to the case

$$a > \psi(t_0), \quad b < b_4$$

and the proof of Theorem III is completed.  $\square$

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