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<td>Author(s)</td>
<td>Fujita, Kento; Nakagawa, Dai; Nakano, Chika; Dao, Quang-Duy; Fujii, Akihiko; Ozaki, Masanori</td>
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Abstract. Small-molecule bulk-heterojunction solar cells utilizing a binary-blended donor material composed of a liquid crystalline phthalocyanine and its analogue have been studied. The short-circuit current density ($J_{sc}$) and power conversion efficiency (PCE) of the solar cells utilizing the blended donor material with an optimized blend ratio reached 11.6 mA/cm$^2$ and 4.8%, respectively, which were superior to those of organic solar cells utilizing each donor material. The improvement of $J_{sc}$ and PCE has been discussed from the viewpoints of the crystal structure, optical and carrier transport properties of the blended donor material.

1. Introduction

Recently, organic solar cells (OSC) based on solution-processed small-molecule (SM) semiconductors have attracted attention as potential alternatives to widely used conjugated polymer-based OSC [1-5]. For commercial applications, solution-processed SM semiconductors offer relatively simple synthesis and purification, monodispersity, well-defined structures without end group contaminants, relatively high charge carrier mobility, and better batch-to-batch reproducibility [3].

Some kinds of liquid crystalline (LC) materials have been demonstrated as novel SM semiconductors for organic thin-film transistors [5-7], organic light-emitting diodes [5,8,9], and solution-processed bulk-heterojunction (BHJ) OSC [9-12]. LC semiconductors exhibit the appropriate characteristics for thin-film optoelectronic devices, such as high solubility in common organic solvents and strong self-organizing nature, which lead to relatively easy fabrication of thin-films with large monodomain and high charge carrier mobility.

One of discotic LC semiconducting materials, 1,4,8,11,15,18,22,25-octaheptylphthalocyanine (C6PcH2) exhibits an ambipolar high carrier drift mobility exceeding 1 cm$^2$/Vs in the crystalline phase, which was evaluated by time-of-flight (TOF) measurement [7]. BHJ OSC utilizing C6PcH2 mixed with 1-(3-methoxy carbonyl)-propyl-1-phenyl-(6,6)C61 (PCBM) demonstrated high power conversion efficiency (PCE) exceeding 4% [10,13]. As further studies, fundamentals of an analogue molecule of C6PcH2, 1,4,8,11,15,18,22,25-octaheptylterabenzothiaziporphyrin (C6TBTAPH2), have been reported, and high PCE exceeding 4.9% has been demonstrated in the BHJ OSC utilizing C6TBTAPH2 mixed with 1-(3-methoxycarbonyl)propyl-1-phenyl-(6,6)C71 (PCBM) [14, 15].

In order to improve the photovoltaic properties of OSC with LC semiconducting materials, we focused on the miscibility. Miscibility is defined as mixing at the molecular level to produce one
homogeneous blend [16]. Introducing the concept of miscibility, improvement of the charge carrier mobility has been reported [17]. On the other hand, molecular alignment control of mixed materials has been reported [18]. Further studies have reported that the PCE of OSC was improved by utilizing the mixed donor materials, such as C6PcH2 and its homologue molecule, 1,4,8,11,15,18,22,25-octapentylphthalocyanine (C5PcH2) [16].

In our recent studies, it was clarified that C6TBTAPH2 and C6PcH2 possess the miscibility in their mixed state, that is, each molecule is mixed at the molecular level [23]. Though the PCE of BHJ OSC is expected to be improved by utilizing C6TBTAPH2 and C6PcH2, the detailed studies remain to be carried out.

In this study, the BHJ OSC utilizing the blended donor materials of C6TBTAPH2 and C6PcH2 are studied. The photovoltaic properties of the solar cells utilizing such binary donor materials are discussed by considering the absorption spectrum, crystal structure in the thin film and hole mobility in the crystal phase.

2. Experimental
C6TBTAPH2 and C6PcH2 were synthesized and purified as reported previously [7, 20]. [70]PCBM (Frontier carbon Ltd.) was used without further purification. The molecular structures of C6TBTAPH2, C6PcH2, and [70]PCBM are shown in Fig. 1.

The solar cells, the schematic device structure of which is also shown in Fig. 1, were fabricated by the following process. Indium tin oxide (ITO)-coated glass substrates were cleaned with detergent, then ultrasonicated in water, chloroform, acetone, and isopropyl alcohol, and subsequently treated with UV-induced ozone. 6-nm-thick MoO3 film was thermally evaporated onto the ITO-coated glass substrates at a rate of 0.1 Å/s under a vacuum of about 2×10⁻⁵ Pa. After transferring the substrates into a N2-filled glove box, a solution, which was prepared by dissolving the donor materials, C6TBTAPH2 and C6PcH2, and acceptor material, [70]PCBM in chloroform with 0.2% v/v 1,8-diiodooctane (DIO), was spin-coated onto the MoO3 layers. The blend ratio of donor and acceptor materials and total concentration were fixed to be 13:10 mol% and 21.5 g/L, respectively. The blend ratios of C6PcH2 in the donor materials were mainly 0, 25, 50, 75, and 100 mol%. The thickness of the active layer with the BHJ of C6TBTAPH2:C6PcH2:[70]PCBM, which was estimated using atomic force microscopy (AFM, Keyence VN-8000), was approximately 120 nm. Finally, 80-nm-thick aluminum electrode was deposited through a shadow mask by thermal evaporation under a vacuum of about 2×10⁻⁵ Pa. The active area of the device was 4 mm².

The absorption spectrum of a thin film prepared on quartz substrate was measured using a spectrophotometer (Shimadzu UV-3150) at room temperature. The energy level at the highest occupied molecular orbital (HOMO) of a material in a thin film prepared on ITO-coated glass substrate was determined by using a measurement system of photoelectron yield spectroscopy (Sumitomo Heavy Industries PYS-202-H) under a vacuum of 10⁻³ Pa at room temperature.

The current density–voltage (J–V) characteristics under an illumination from a full spectrum solar simulator (Bunkoukeiki, OTENTO-SUNIII P2) at an irradiation intensity of 100 mW/cm² (AM1.5G) were measured in vacuum at room temperature with a source measure unit (Keithley 2400). The external quantum efficiency (EQE) spectra were measured by using a photosensitivity measurement system (Bunkoukeiki, SM-250NA) with a Xe lamp light passing through a monochromator as a light source, the intensity of which was calibrated using a Si reference cell. The integrated EQE values indicated good agreements with the measured short-circuit current density ($J_{sc}$).

The charge carrier mobility under a low vacuum (10⁻¹ Pa) was determined by the TOF technique. A sandwich cell consisting of two ITO-coated glass substrates and a 12-μm-thick PET spacer was fabricated. The actual cell gap was evaluated by an interference technique of light transmittance. Subsequently, binary-blended materials of C6TBTAPH2 and C6PcH2, which were prepared by vaporizing the solvent in the solution of C6PcH2 and C6TBTAPH2, were heated to a temperature in the isotropic phase for injecting into the cell by capillary action, and cooled down to room temperature.
in vacuum. A DC bias was applied between the electrodes using batteries connected in series, and the films were irradiated using a Nd:YAG laser (wavelength: 355 nm, pulse width: 1 ns). The transient current response was observed using an oscilloscope (Teledyne LeCroy HDO4054) in vacuum at room temperature, and the charge carrier mobility $\mu$ was calculated according to the equation $\mu = \frac{d^2}{V\tau}$, where $d$, $V$, and $\tau$ are the film thickness, applied bias, and transit time, respectively.

**Figure 1.** Molecular structures of C6TBTAPH2, C6PcH2, and [70]PCBM, and schematic device structure of the C6TBTAPH2:C6PcH2:[70]PCBM BHJ OSC.

### 3. Results and Discussion

Figure 2 shows the absorption spectra of C6TBTAPH2:C6PcH2:[70]PCBM thin films with various C6PcH2 blend ratios. Two high absorption peaks at 300-450 nm and 600-800 nm were observed, which correspond to the optical transition at the B-band and Q-band of C6TBTAPH2 or C6PcH2, respectively. The absorbance of C6TBTAPH2 at the B-band is higher than that of C6PcH2, while the absorbance of C6TBTAPH2 at the Q-band is lower than that of C6PcH2 [15]. In the C6TBTAPH2:C6PcH2:[70]PCBM thin films, therefore, the absorbance at the B-band decreased, and that at the Q-band increased with increasing the blend ratio of C6PcH2. It is noted that the marked change of the optical properties, such as absorption spectrum, appeared by adjusting the blend ratio of C6TBTAPH2:C6PcH2, in spite of the slight difference of the molecular structures.

**Figure 2.** Absorption spectra of C6TBTAPH2:C6PcH2:[70]PCBM thin films with various C6PcH2 blend ratios.

**Figure 3.** $J-V$ characteristics of the C6TBTAPH2:C6PcH2:[70]PCBM BHJ OSC with various C6PcH2 blend ratios.
The $J$-$V$ characteristics of C6TBTAPH$_2$:C6PcH$_2$:PCBM BHJ OSC are shown in Fig. 3, and the photovoltaic properties depending on the C6PcH$_2$ blend ratio are summarized in Table I. The open-circuit voltage ($V_{oc}$) of the solar cell with C6TBTAPH$_2$:PCBM BHJ or C6PcH$_2$:PCBM BHJ was 0.70 or 0.73 V, respectively. The $V_{oc}$ of the solar cells with the blended donors of C6TBTAPH$_2$:C6PcH$_2$ were 0.71 V, which was close to that with C6TBTAPH$_2$:PCBM BHJ. The $V_{oc}$ of OSC is explained by an energy gap between the HOMO level of donors and the lowest occupied molecular orbital (LUMO) level of acceptors ($\Delta E_{DA}$) and/or work function of metal electrodes [21, 22]. Although the difference between $V_{oc}$ and $\Delta E_{DA}$ is still subject to debate, a linear relationship between $V_{oc}$ and $\Delta E_{DA}$ has been experimentally confirmed [22]. The HOMO levels of C6TBTAPH$_2$ and C6PcH$_2$ were estimated to be 5.25 eV and 5.32 eV, respectively. Hence, the energy gap between the HOMO level of C6TBTAPH$_2$ and the LUMO level of PCBM tends to be predominant for the $V_{oc}$ of BHJ OSC with the blended donors of C6TBTAPH$_2$ and C6PcH$_2$. The fill factors (FF) of the solar cells with various blend ratios of C6PcH$_2$ were evaluated to be 0.56-0.58.

<table>
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<tr>
<th>Blend Ratio of C6PcH$_2$ (mol%)</th>
<th>$V_{oc}$ (V)</th>
<th>$J_{sc}$ (mA/cm$^2$)</th>
<th>FF</th>
<th>PCE (%)</th>
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<tr>
<td>0</td>
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<td>0.57</td>
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<td>25</td>
<td>0.71</td>
<td>11.6</td>
<td>0.58</td>
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</tr>
<tr>
<td>50</td>
<td>0.71</td>
<td>11.3</td>
<td>0.58</td>
<td>4.6</td>
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<tr>
<td>75</td>
<td>0.71</td>
<td>10.9</td>
<td>0.56</td>
<td>4.3</td>
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<tr>
<td>100</td>
<td>0.73</td>
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The $J_{sc}$ and PCE of C6TBTAPH$_2$:C6PcH$_2$:PCBM BHJ OSC were markedly depending on the C6PcH$_2$ blend ratio as shown in Fig. 4. The $J_{sc}$ and PCE of C6TBTAPH$_2$:PCBM BHJ OSC were 11.3 mA/cm$^2$ and 4.5%, respectively, while those of C6PcH$_2$:PCBM BHJ OSC were 10.6 mA/cm$^2$ and 4.1%, respectively. In the case of the C6TBTAPH$_2$:C6PcH$_2$:PCBM BHJ OSC at 25 mol% of C6PcH$_2$ blend ratio, the $J_{sc}$ and PCE were improved to be 11.6 mA/cm$^2$ and 4.8%, respectively. It is suggested that the enhancement of the $J_{sc}$ mainly contribute to the improvement of the PCE.

Figure 5 shows the EQE spectra of C6TBTAPH$_2$:C6PcH$_2$:PCBM BHJ OSC with various C6PcH$_2$ blend ratios. The main origins of the EQE peaks should be optical excitations at the B-band and Q-band of C6TBTAPH$_2$ or C6PcH$_2$. The EQE at B-band in the C6TBTAPH$_2$:PCBM BHJ OSC is higher than that in the C6PcH$_2$:PCBM BHJ OSC, while the EQE at Q-band in the C6TBTAPH$_2$:PCBM BHJ OSC was lower than that in the C6PcH$_2$:PCBM BHJ OSC. With increasing the blend ratio of C6PcH$_2$ in the C6TBTAPH$_2$:C6PcH$_2$:PCBM BHJ OSC, the EQE at B-band decreased and that at Q-band slightly changed, which must be strongly related with the blend ratio dependence of the absorption spectra as abovementioned.

The EQE is defined as the ratio of the number of charge carriers collected at the electrodes in the solar cell to that of incident photons. In solar radiation spectrum of AM1.5, the number of photons around the wavelength of Q band is about two times more than that at B band. Comparing the EQE spectra at 0 mol% and 25 mol% of C6PcH$_2$ blend ratios, the quenching of the EQE at B-band with increasing the C6PcH$_2$ blend ratio seems to be more marked than the enhancement of the EQE at Q-band, however, the enhanced $J_{sc}$ could be confirmed by integrating the EQE spectra indeed. The enhanced $J_{sc}$ could be explained by the photon numbers used for the photon-electron conversion rather
than the EQE. Therefore, the photosensitivity at the Q-band rather than the B-band should be still predominant for the device performance at 0 mol% and 25 mol% of C6PcH2 blend ratios.

In the case of the C6PcH2 blend ratio beyond 25 mol%, the absorbance at 400 nm monotonically decreased as shown in Fig. 6. With increasing the C6PcH2 blend ratio, the absorption loss at B-band should be a main factor, which caused the decreasing of the PCE as well as \( J_{sc} \). As a result, the highest PCE appeared at 25 mol% of C6PcH2 blend ratio.

In order to discuss about the carrier transport in the C6TBTAPH2:C6PcH2 blended film, the hole mobility as a fundamental property of the blend film was investigated by TOF technique. Figure 6 also shows the C6PcH2 blend ratio dependence of hole mobility in the blended films of C6TBTAPH2 and C6PcH2. The hole mobility in the C6TBTAPH2 film was 0.23 cm²/Vs. In the blended films below 25 mol% of C6PcH2 blend ratio, the mobility could not be evaluated, because it is hard to determine the transit time due to the dispersive decay curve with weak photocurrent. The hole mobility could be evaluated to be 0.47 cm²/Vs at 25 mol% of C6PcH2 blend ratio, and gradually enhanced with increasing the blend ratio of C6PcH2, eventually 0.96 cm²/Vs at 100 mol% of C6PcH2. It is noted that the hole mobility changes depending on the blend ratio in spite of the slight difference on the molecular structures and that the blank region of the hole mobility exist below 25 mol% of C6PcH2 blend ratio.

The crystal structures in the single crystal or thin film were reported to be based on a triclinic lattice for C6TBTAPH2 and a monoclinic one for C6PcH2 [23-25]. Further studies exhibited that the crystal structures in blended single crystals of C6TBTAPH2 and C6PcH2 at the C6PcH2 blend ratio below 25 mol% were similar to the triclinic lattice of C6TBTAPH2 [23]. In the range of 25-100 mol% of C6PcH2 blend ratio, the crystal structures in the blended single crystals of C6TBTAPH2 and C6PcH2 were similar to the monoclinic lattice of C6PcH2 [23]. It is considered, therefore, that the relatively high hole mobility originated from the C6PcH2–based molecular packing structure [19].

As shown in Fig. 6, with increasing a blend ratio of C6PcH2, the hole mobility was enhanced, while the absorbance at B-band was suppressed. At 25 mol% of C6PcH2 blend ratio, the carrier generation induced by the photo-excitation at Q-band rather than B-band was predominant, and the hole mobility was higher than that of C6TBTAPH2 because of the C6PcH2–based crystal structure. Thus, the highest \( J_{sc} \) and PCE were achieved at 25 mol% of C6PcH2 blend ratio.
Figure 6. C6PcH2 blend ratio dependence of the hole mobility in the C6TBTAPH2:C6PcH2 blended films and the absorbance at 400 nm of the C6TBTAPH2:C6PcH2:[70]PCBM BHJ films.

4. Conclusions
The charge carrier mobility in the blend films of C5PcH2 and C6PcH2 were investigated. The phase diagram of the mixture of C5PcH2 and C6PcH2 was determined from the thermal properties, it was found that the mixture possess the miscibility. In the case of C5PcH2 blend ratio below 25 mol%, both hole and electron mobilities were kept high values comparable with those of C6PcH2, while both hole and electron mobilities markedly decreased by one order of magnitude in the case of C5PcH2 blend ratio above 30 mol%. From the results of XRD measurement, the molecular structure of the mixture film obviously changed depending on the blend ratio of C5PcH2, that must be strongly related with the behavior of the charge carrier transport.

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