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SOME EXAMPLES OF HYPOELLIPTIC OPERATORS OF INFINITELY DEGENERATE TYPE

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0. Introduction

The object of the present paper is to study some examples of the operators of the form

$$(1) \quad P = D_x^2 + a(x)D_y^2 + b(x)D_y,$$

on \mathbf{R}^2 where $D_x = -i\frac{\partial}{\partial x}$, $D_y = -i\frac{\partial}{\partial y}$, $a(x)$ and $b(x)$ are functions satisfying:

- $$(2) \quad \begin{aligned} (i) \quad & a(x), b(x) \in C^\infty(\mathbf{R}), \\ (ii) \quad & a(x) > 0 \text{ for } x \neq 0, \partial^\alpha a(0) = \partial^\alpha b(0) = 0 \text{ for any } \alpha. \end{aligned}$$

We consider here C^∞ -hypoellipticity of the operator P on $x=0$. In general it is hypoelliptic if $b(x)$ is small compared with $a(x)$, and conversely, not hypoelliptic if $b(x)$ is big. Such conditions for the hypoellipticity were investigated in the previous paper [5]. But the examples considered here cannot be explained by the method of [5] (we cannot regard $b(x)$ small nor big in what follows). They are analogous to the one which A. Menikoff considered in [6], i.e., the finitely degenerate case where $a(x)=x^{2h}$ and $b(x)=bx^{h-1}$. We prove the following theorems.

Theorem 1. *Let $a(x)=|x|^{-4} \exp(-2|x|^{-1})$ and $b(x)=b \cdot |x|^{-4} \exp(-|x|^{-1})$ with b being a complex constant. Then the operator P is hypoelliptic if and only if b is not odd integer.*

Theorem 2. *Let $a(x)=|x|^{-4} \exp(-2|x|^{-1})$ and $b(x)=b \cdot \operatorname{sgn} x \cdot |x|^{-4} \exp(-|x|^{-1})$ with b being a complex constant. Then the operator P is hypoelliptic.*

REMARK 1: By the similar argument of the proof of theorem 1 in T. Morioka [8], we can conclude that P is micro-hypoelliptic when P is hypoelliptic.

The hypoellipticity of P is closely connected to the branching of singularities of solutions for the weakly hyperbolic operator $Q = -D_x^2 + a(x)D_y^2 + b(x)D_y$.

G.R. Aleksandryan [1] dealt with the one for Q which corresponds the cases in Theorem 1 and Theorem 2. In Section 1, we shall prove the non-hypoellipticity part of Theorem 1, by using the observation of Aleksandryan. Section 2 is devoted to the proof of hypoellipticity parts of Theorem 1 and Theorem 2. We shall show them by constructing the parametrix of P explicitly.

The author is grateful to Professor K. Kajitani for introducing him the article [1] of Aleksandryan.

1. Proof of non-hypoellipticity

In this section we prove that P is not hypoelliptic if $a(x)$ and $b(x)$ are those in Theorem 1, and b satisfies $b=2n+1$ for some $n \in \mathbb{Z}$. Also we shall explain the reason why Theorem 2 is free from such a condition. Here we adopt the notations from Aleksandryan [1].

At first, let us set $\Lambda(x) = \exp(-|x|^{-1})$ and $\mu(x) = \Lambda'(x)$ ($= \operatorname{sgn} x \cdot |x|^{-2} \exp(-|x|^{-1})$). Then the partial Fourier transform of the equation $Pu=0$ with respect to y can be written in the following form:

$$(3) \quad -\hat{u}_{xx} + \left(\mu(x)^2 \eta^2 + b \frac{\mu(x)^2}{\Lambda(x)} \eta \right) \hat{u} = 0.$$

Furthermore making a change in such a way that $\hat{u}(x, \eta) = xw(\tau)$, $\tau = \Lambda(x)\eta$, it becomes

$$(4) \quad -w_{\tau\tau} - \frac{w_\tau}{\tau} + \left(1 + \frac{b}{\tau} \right) w = 0.$$

Set now $z=2\tau$ and $f(z) = e^{z^2/2} w\left(\frac{z}{2}\right)$. Then (4) turns into Kummer's equation

$$(5) \quad zf''(z) + (1-z)f'(z) - \alpha f(z) = 0,$$

where $\alpha = \frac{1+b}{2}$. Hence we have the following

Proposition 1. (i) Suppose $\eta > 0$. Then there exist solutions $\hat{u}_1(x, \eta)$ and $\hat{u}_2(x, \eta)$ of (3) which have the following expressions:

$$\hat{u}_1(x, \eta) = xe^{-\Lambda(x)\eta} \Psi(\alpha, 1; 2\Lambda(x)\eta) \quad \text{for } x > 0,$$

and

$$\hat{u}_2(x, \eta) = -xe^{-\Lambda(x)\eta} \Psi(\alpha, 1; 2\Lambda(x)\eta) \quad \text{for } x < 0,$$

where $\Psi(\alpha, 1; z)$ is a solution of (5) for $z > 0$ defined in A. Erdelyi et al. [2, page 255-256].

(ii) Suppose $\eta < 0$. Then there exist solutions $\hat{u}_1(x, \eta)$ and $\hat{u}_2(x, \eta)$ of (3)

which have the following expressions :

$$\hat{u}_1(x, \eta) = xe^{\Lambda(x)\eta} \Psi(1-\alpha, 1; -2\Lambda(x)\eta) \quad \text{for } x > 0,$$

and

$$\hat{u}_2(x, \eta) = -xe^{\Lambda(x)\eta} \Psi(1-\alpha, 1; -2\Lambda(x)\eta) \quad \text{for } x < 0.$$

REMARK 2: It holds that $\hat{u}_2(x, \eta) = \hat{u}_1(-x, \eta)$ for $x < 0$. Generally, it does not hold that $\hat{u}_1(x, \eta) = -\hat{u}_2(x, \eta)$ (they are linearly independent in generic case), because $\Psi(\alpha, \gamma; z)$ is many-valued holomorphic function of z and its principal branch can be at most defined in the plane cut along negative real axis (see page 257 of [2]).

REMARK 3: Since $\Psi(\alpha, \gamma; z) = O(z^{-\alpha})$ as positive number z tends to infinity (see [2, page 278]), $\hat{u}_1(x, \eta)$ is uniformly bounded for $(\log 2|\eta|)^{-1} \leq x \leq 1$ and also $\hat{u}_2(x, \eta)$ is uniformly bounded for $-1 \leq x \leq -(\log 2|\eta|)^{-1}$.

Proof. (i) We can see the result concerning $\hat{u}_1(x, \eta)$ since $z = \Lambda(x)\eta > 0$ for $x > 0$ and $\eta > 0$ (recall that $\Psi(\alpha, 1; z)$ satisfies (5) for $z > 0$). In order to obtain the result concerning $\hat{u}_2(x, \eta)$, we make the change of variable $\tilde{x} = -x$ in the equation (3) (notice that $\tilde{x} > 0$ for $x < 0$). Then (3) becomes the same equation with respect to the variable \tilde{x} since $\Lambda(\tilde{x}) = \Lambda(x)$ and $\mu(\tilde{x})^2 = \mu(x)^2$. Thus we can see that there is a solution of (3) which have the expression: $\hat{u}_2(x, \eta) = \tilde{x}e^{-\Lambda(\tilde{x})\eta} \Psi(\alpha, 1; 2\Lambda(\tilde{x})\eta)$ for $\tilde{x} > 0$. This implies the result.

(ii) To obtain the result concerning $\hat{u}_1(x, \eta)$, we set $z = -2\tau$ and $f(z) = e^{z^2/2} w\left(-\frac{z}{2}\right)$ in the equation (4) (notice that $z = -2\Lambda(x)\eta > 0$ for $x > 0, \eta < 0$). Then (4) becomes

$$(5') \quad zf''(z) + (1-z)f'(z) - (1-\alpha)f(z) = 0,$$

and this implies the result. The argument to obtain the one concerning $\hat{u}_2(x, \eta)$ is also similar. ■

Next we investigate the Wronskian of \hat{u}_1 and \hat{u}_2 , namely,

$$W(\eta) = \hat{u}_1(0, \eta)\hat{u}_2'(0, \eta) - \hat{u}_1'(0, \eta)\hat{u}_2(0, \eta).$$

We can compute the value of $W(\eta)$ which is essential to the proof of Theorem 1.

Proposition 2. (i) For $\eta > 0$, it holds that

$$(6) \quad W(\eta) = \frac{2}{\Gamma(\alpha)^2} \{\log 2\eta + \psi(\alpha) - 2\psi(1)\},$$

where $\Gamma(\alpha)$ is Euler's Gamma function and $\psi(\alpha) = \Gamma'(\alpha)/\Gamma(\alpha)$.

(ii) For $\eta < 0$, it holds that

$$(7) \quad W(\eta) = \frac{2}{\Gamma(1-\alpha)^2} \{ \log(-2\eta) + \psi(1-\alpha) - 2\psi(1) \}.$$

Proof. Here we prove the case (i). The argument for the proof of (ii) is completely parallel if α and η are respectively replaced by $1-\alpha$ and $-\eta$.

At first, let us recall that $\Psi(\alpha, n+1; z)$ ($n=0, 1, \dots$) has the following asymptotic behavior as $z \downarrow 0$ (see page 261 of [2]):

$$(8) \quad \begin{aligned} \Psi(\alpha, n+1; z) &= \frac{(-1)^{n-1}}{n! \Gamma(\alpha-n)} \{ \Phi(\alpha, n+1; z) \log z \\ &\quad + \sum_{r=0}^{\infty} \frac{(\alpha)_r}{(n+1)_r} [\psi(\alpha+r) - \psi(1+r) - \psi(1+n-r)] \frac{z^r}{r!} \} \\ &\quad + \frac{(n-1)!}{\Gamma(\alpha)} \sum_{r=0}^{n-1} \frac{(\alpha-n)_r}{(1-n)_r} \cdot \frac{z^{r-n}}{r!}, \end{aligned}$$

where $(\alpha)_r = \alpha(\alpha+1)\cdots(\alpha+r-1)$ and

$$\Phi(\alpha, \gamma; z) = \sum_{r=0}^{\infty} \frac{(\alpha)_r}{(\gamma)_r} \cdot \frac{z^r}{r!}.$$

Hence we can conclude that

$$(9) \quad \begin{aligned} \hat{u}_1(0, \eta) &= \lim_{x \downarrow 0} x e^{-\Lambda(x)\eta} \Psi(\alpha, 1; 2\Lambda(x)\eta) \\ &= \frac{1}{\Gamma(\alpha)}. \end{aligned}$$

Next let us recall the following relation (see page 258 of [2]):

$$(10) \quad \frac{d}{dz} \Psi(\alpha, \gamma; z) = -\alpha \Psi(\alpha+1, \gamma+1; z).$$

This implies that

$$(11) \quad \begin{aligned} \hat{u}_1'(x, \eta) &= e^{-\Lambda(x)\eta} \Psi(\alpha, 1; 2\Lambda(x)\eta) \\ &\quad - x \Lambda'(x) \eta e^{-\Lambda(x)\eta} \Psi(\alpha, 1; 2\Lambda(x)\eta) \\ &\quad - 2\alpha x \Lambda'(x) \eta e^{-\Lambda(x)\eta} \Psi(\alpha+1, 2; 2\Lambda(x)\eta), \end{aligned}$$

for $x > 0$. Take now the limit of the equation (11) as $x \downarrow 0$, keeping (8) in mind. Then the cancelation will occur between the terms of order $O(x^{-1})$. Thus we get

$$(12) \quad \hat{u}_1'(0, \eta) = -\frac{1}{\Gamma(\alpha)} \{ \log 2\eta + \psi(\alpha) - 2\psi(1) \},$$

Similarly, from the expression of $\hat{u}_2(x, \eta)$ for $x < 0$, we obtain

$$(13) \quad \begin{aligned} \hat{u}_2(0, \eta) &= \frac{1}{\Gamma(\alpha)}, \\ \hat{u}_2'(0, \eta) &= \frac{1}{\Gamma(\alpha)} \{\log 2\eta + \psi(\alpha) - 2\psi(1)\}. \end{aligned}$$

The equations (9), (12) and (13) immediately give our assertion.

REMARK 4. From Proposition 2, we can see that $\hat{u}_1(x, \eta)$ and $\hat{u}_2(x, \eta)$ are linearly dependent for $\eta > 0$ if $b = -1, -3, -5, \dots$ and for $\eta < 0$ if $b = 1, 3, 5, \dots$ (recall that $\alpha = \frac{1+b}{2}$ and $\frac{1}{\Gamma(-n)} = 0$ for $n = 0, 1, 2, \dots$). Also it is clear that, for sufficiently large $|\eta|$, $\hat{u}_1(x, \eta)$ and $\hat{u}_2(x, \eta)$ are linearly independent if b is not odd integer.

Next we investigate Theorem 2. In the case of Theorem 2 we consider the equation (3) with $\Lambda(x)$ and $\mu(x)$ being respectively replaced by $\tilde{\Lambda}(x) = \operatorname{sgn} x \cdot \exp(-|x|^{-1})$ and $\tilde{\mu}(x) = \tilde{\Lambda}'(x)$. The similar argument as above gives us the following

Proposition 3. (i) For $\eta > 0$, there exist solutions which have the following expressions:

$$\hat{u}_1(x, \eta) = x e^{-\tilde{\Lambda}(x)\eta} \Psi(\alpha, 1; 2\tilde{\Lambda}(x)\eta) \quad \text{for } x > 0,$$

and

$$\hat{u}_2(x, \eta) = -x e^{-\tilde{\Lambda}(x)\eta} \Psi(1-\alpha, 1; -2\tilde{\Lambda}(x)\eta) \quad \text{for } x < 0.$$

Moreover the Wronskian of them is

$$W(\eta) = \frac{1}{\Gamma(\alpha) \Gamma(1-\alpha)} \{2 \log 2\eta - 4\psi(1) + \psi(\alpha) + \psi(1-\alpha)\}.$$

(ii) For $\eta < 0$, there exist solutions which have the following expressions:

$$\hat{u}_1(x, \eta) = x e^{\tilde{\Lambda}(x)\eta} \Psi(1-\alpha, 1; -2\tilde{\Lambda}(x)\eta) \quad \text{for } x > 0,$$

and

$$\hat{u}_2(x, \eta) = -x e^{\tilde{\Lambda}(x)\eta} \Psi(\alpha, 1; 2\tilde{\Lambda}(x)\eta) \quad \text{for } x < 0.$$

Moreover the Wronskian of them is

$$W(\eta) = \frac{1}{\Gamma(\alpha) \Gamma(1-\alpha)} \{2 \log(-2\eta) - 4\psi(1) + \psi(\alpha) + \psi(1-\alpha)\}.$$

REMARK 5. As in the case of Theorem 1, we can conclude from Proposition 3 that, for sufficiently large $|\eta|$, $\hat{u}_1(x, \eta)$ and $\hat{u}_2(x, \eta)$ are linearly independent if b is not odd integer (i.e., $\alpha \in \mathbb{Z}$). Moreover, even if $\alpha \in \mathbb{Z}$,

$$W(\eta) = \frac{1}{\Gamma(\alpha)\Gamma(1-\alpha)} \{\psi(\alpha) + \psi(1-\alpha)\} \\ = \begin{cases} 1 & \text{if } |2\alpha-1| \equiv 3 \pmod{4}, \\ -1 & \text{if } |2\alpha-1| \equiv 1 \pmod{4}. \end{cases}$$

(See page 15 of [2].) Thus we see that $\hat{u}_2(x, \eta)$ and $\hat{u}_2(x, \eta)$ are linearly independent for such α . This is the reason why Theorem 2 is free from such an assumption as in Theorem 1.

Now we turn to prove the non-hypoellipticity part of Theorem 1.

Proof of non-hypoellipticity in Theorem 1. First let us observe that, if P is hypoelliptic, we get the following inequality from the argument of Banach's closed graph theorem.

For any positive number l and for any pair of open sets Ω and Ω' satisfying $\bar{\Omega}' \subset \Omega$, there exist a positive integer m and a constant C such that

$$(14) \quad \|D_y^l u\|_{L^2(\Omega')} \leq C \left\{ \sum_{m_1+m_2 \leq m} \|D_x^{m_1} D_y^{m_2} P u\|_{L^2(\Omega)} + \|u\|_{L^2(\Omega)} \right\}, \\ \forall u \in C^\infty(\bar{\Omega}).$$

We are now going to show that the inequality (14) never holds provided b is odd integer. Let us set $\Omega = (-\delta, \delta) \times (-\delta, \delta)$ and $\Omega' = (-\delta', \delta') \times (-\delta', \delta')$ with δ and δ' satisfying $0 < \delta' < \delta < 1$. Moreover set

$$(15) \quad u_\eta(x, y) = e^{iy^\eta} \hat{u}_1(x, \eta) \\ = -e^{iy^\eta} \hat{u}_2(x, \eta),$$

with $\eta > 0$ if $b = -1, -3, \dots$ and with $\eta < 0$ if $b = 1, 3, \dots$ (Observe that $\hat{u}_1(x, \eta) = -\hat{u}_2(x, \eta)$ provided b is odd integer. To see this, compare (9), (12) and (13).) Let us substitute $u_\eta(x, y)$ into (14) and compare the asymptotic behavior of the both hand sides as $|\eta| \rightarrow \infty$. Clearly, in the right hand side, it holds that $P u_\eta = 0$.

Observe now that there exists a constant C (independent of η) such that

$$|\hat{u}_1(x, \eta)| \leq C \quad \text{for } 0 \leq x \leq 1,$$

and

$$|\hat{u}_2(x, \eta)| \leq C \quad \text{for } -1 \leq x \leq 0.$$

This can be seen from the remark after the statement of Proposition 1 and the asymptotic behaviors of $\Psi(\alpha, 1; z)$ and $\Psi(1-\alpha, 1; z)$ as $z \downarrow 0$. Indeed, for example, it follows from (8) that

$$|\hat{u}_1(x, \eta)| = |x| e^{-\Lambda(x)\eta} |\Psi(\alpha, 1; 2\Lambda(x)\eta)| \\ \leq C_1 |x| (1 + |\log 2\Lambda(x)\eta|)$$

$$\begin{aligned} &\leq C_1(|x| + |x| \log 2 + |x \log \Lambda(x)| + |x \log \eta|) \\ &\leq C_2, \end{aligned}$$

for $0 < x < (\log 2\eta)^{-1}$ and $\eta \geq e/2$. Hence, if we substitute u_η into (14), the right hand side is not larger than

$$(16) \quad \|u_\eta\|_{L^2(\Omega)} \leq 4\delta^2 \cdot C.$$

On the other hand, in the left hand side of (14), it is clear that

$$\|D_j^l u_\eta\|_{L^2(\Omega')} = |\eta|^l \cdot 2\delta' \cdot \left(\int_{-\delta'}^{\delta'} |\hat{u}_1(x, \eta)|^2 dx \right)^{1/2}.$$

Moreover, from the asymptotic behavior of $\Psi(\alpha, 1; z)$ as $z \rightarrow \infty$, it follows that there exist positive constants ε and M such that

$$|\hat{u}_1(x, \eta)| \geq \varepsilon |x| \quad \text{for } M \leq 2\Lambda(x) |\eta| \leq 2M.$$

Hence we obtain that

$$\begin{aligned} (17) \quad \|D_j^l u_\eta\|_{L^2(\Omega')} &\geq |\eta|^l \cdot 2\delta' \cdot \left(\int_{M \leq 2\Lambda(x) |\eta| \leq 2M} |\hat{u}_1(x, \eta)|^2 dx \right)^{1/2} \\ &\geq |\eta|^l \cdot 2\delta' \cdot \varepsilon \cdot 3^{-1/2} \left\{ \left(\log \frac{|\eta|}{M} \right)^{-3} - \left(\log \frac{2|\eta|}{M} \right)^{-3} \right\}^{1/2}. \end{aligned}$$

Finally taking $l \geq 1$ immediately implies the contradiction among (14), (16) and (17). ■

2. Proof of hypoellipticity

In the present section, we assume that $W(\eta) \neq 0$ for $|\eta| \geq C$, and denote by $Q(x, x'; \eta)$ the Green function of (3) (in the case of Theorem 2, $\Lambda(x)$ and $\mu(x)$ being replaced respectively by $\tilde{\Lambda}(x)$ and $\tilde{\mu}(x)$), i.e.,

$$Q(x, x'; \eta) = \begin{cases} \frac{\hat{u}_2(x, \eta) \hat{u}_1(x', \eta)}{W(\eta)} & (x < x'), \\ \frac{\hat{u}_2(x', \eta) \hat{u}_1(x, \eta)}{W(\eta)} & (x' < x). \end{cases}$$

Then we have the following

Proposition 4. *For any non-negative integer m , there exists a constant C_m such that the following inequalities hold:*

$$(18) \quad \int_{-1}^1 |\partial_\eta^m Q(x, x'; \eta)| dx' \leq C_m |\eta|^{-m} \text{ for } -1 \leq x \leq 1 \text{ and } |\eta| \geq \max\{C, e\},$$

$$(19) \quad \int_{-1}^1 |\partial_\eta^m Q(x, x'; \eta)| dx \leq C_m |\eta|^{-m} \text{ for } -1 \leq x' \leq 1 \text{ and } |\eta| \geq \max\{C, e\}.$$

Proof. Here we prove the proposition in the case of Theorem 1. First we shall verify (18) when $m=0$. Observe now the following inequality:

$$(20) \quad \int_{-1}^1 |Q(x, x'; \eta)| dx' \leq \left\{ |\hat{u}_2(x, \eta)| \int_x^1 |\hat{u}_1(x', \eta)| dx' + |\hat{u}_1(x, \eta)| \int_{-1}^x |\hat{u}_2(x', \eta)| dx' \right\} \cdot |W(\eta)|^{-1}.$$

Let us set $x_\eta = (\log |\eta|)^{-1}$ (then $\Lambda(x_\eta)|\eta| = 1$). We are going to estimate the right hand side of (20). Here we assume $\eta > 0$. In the case of $\eta < 0$, the argument is completely parallel if α is replaced by $1-\alpha$.

(I) Now we are going to show that the value of $\int_{-1}^1 |Q(x, x'; \eta)| dx'$ is uniformly bounded for $x_\eta \leq x \leq 1$ and $\eta \geq \max\{C, e\}$. Concerning the first term on the right hand side of (20), we can use the expression of $\hat{u}_1(x', \eta)$ for $x' > 0$ and the asymptotic behavior $\Psi(\alpha, \gamma, z) = O(z^{-\alpha})$ as $z \rightarrow \infty$. Hence we have

$$|\hat{u}_1(x', \eta)| \leq C e^{-\Lambda(x')\eta} (\Lambda(x')\eta)^{-\alpha} \quad \text{for } x_\eta \leq x' \leq 1.$$

We cannot use the expression of $\hat{u}_2(x, \eta)$ for $x > 0$. So let us express it by linear combination of $\hat{u}_1(x, \eta)$ and

$$\hat{u}_3(x, \eta) = x e^{-\Lambda(x)\eta} \Phi(\alpha, 1; 2\Lambda(x)\eta) \quad \text{for } x > 0$$

(concerning the definition of $\Phi(\alpha, \gamma; z)$, see page 248 of [2]). From the facts that $\hat{u}_3(0, \eta) = 0$ and $\hat{u}_3'(0, \eta) = 1$, it follows

$$(21) \quad \hat{u}_2(x, \eta) = A \hat{u}_1(x, \eta) + B \hat{u}_3(x, \eta),$$

where

$$A = 1, \\ B = \frac{2}{\Gamma(\alpha)} \{\log 2\eta + \psi(\alpha) - \psi(1)\}.$$

Hence we obtain

$$(22) \quad \begin{aligned} & |\hat{u}_2(x, \eta)| \int_x^1 |\hat{u}_1(x', \eta)| dx' \leq \{ |W(\eta)|^{-1} \cdot |\hat{u}_1(x, \eta)| + C \cdot |\hat{u}_3(x, \eta)| \} \int_x^1 |\hat{u}_1(x', \eta)| dx'. \end{aligned}$$

Now recall that

$$\Phi(\alpha, \gamma; z) = \frac{\Gamma(\gamma)}{\Gamma(\alpha)} e^z z^{\alpha-\gamma} (1 + O(|z|^{-1})) \text{ as } z \rightarrow +\infty$$

(see page 278 of [2]), and $\Lambda'(x')\eta = \mu(x')\eta \geq (x_\eta)^{-2}$ for $x_\eta \leq x' \leq 1$. Hence, concerning the second term on the right of (22), we have furthermore

$$\begin{aligned} |\hat{u}_3(x, \eta)| & \int_x^1 |\hat{u}_1(x', \eta)| dx' \\ & \leq C_1 \cdot x_\eta^2 \cdot e^{\Lambda(x)\eta} (\Lambda(x)\eta)^{\alpha-1} \int_x^1 (\Lambda(x')\eta)^{-\alpha} e^{-\Lambda(x')\eta} \Lambda'(x')\eta dx' \\ & \leq C_1 \cdot x_\eta^2 \cdot e^{\Lambda(x)\eta} (\Lambda(x)\eta)^{\alpha-1} \int_{\Lambda(x)\eta}^\infty t^{-\alpha} e^{-t} dt \\ & \leq C_2 \cdot x_\eta^2 \cdot (\Lambda(x)\eta)^{-1} \leq C_3. \end{aligned}$$

Here we have used the fact that $\int_s^\infty t^{-\alpha} e^{-t} dt = O(s^{-\alpha} e^{-s})$ as $s \rightarrow +\infty$. The similar argument is applicable for estimating the first term on the right of (22). Consequently, the first term on the right of (20) is uniformly bounded for $x_\eta \leq x \leq 1$ and $\eta \geq \max\{C, e\}$.

Concerning the second term on the right of (20), let us decompose it in the following way:

$$\begin{aligned} |\hat{u}_1(x, \eta)| & \int_{-1}^x |\hat{u}_2(x', \eta)| dx' \cdot |W(\eta)|^{-1} = |\hat{u}_1(x, \eta) \cdot W(\eta)^{-1}| \\ & \times \left\{ \int_{-1}^{-x_\eta} |\hat{u}_2(x', \eta)| dx' + \int_{-x_\eta}^{x_\eta} |\hat{u}_2(x', \eta)| dx' + \int_{x_\eta}^x |\hat{u}_2(x', \eta)| dx' \right\}. \end{aligned}$$

For $x' \leq x_\eta$, the expression of $\hat{u}_2(x', \eta)$ can be applied, and also for $x_\eta \leq x'$, $\hat{u}_2(x', \eta)$ can be decomposed as (21). Thus, by using the asymptotic behaviors of $\Psi(\alpha, 1; z)$ and $\Phi(\alpha, 1; z)$ as $z \rightarrow \infty$, we see that the first and the third terms are uniformly bounded. Concerning the integral with $-x_\eta \leq x' < 0$, the expression of $\hat{u}_2(x', \eta)$ and the asymptotic behavior of $\Psi(\alpha, 1; z)$ as $z \downarrow 0$ can be applied (see page 262 of [2]). Hence it holds that

$$\begin{aligned} |\hat{u}_1(x, \eta) \cdot W(\eta)^{-1}| & \cdot \int_{-x_\eta}^0 |\hat{u}_2(x', \eta)| dx' \\ & \leq C_1 e^{-\Lambda(x)\eta} (\Lambda(x)\eta)^{-\alpha} \cdot |W(\eta)|^{-1} \int_{-x_\eta}^0 |x' \log 2\Lambda(x')\eta| dx' \\ & \leq C_2 \cdot |W(\eta)|^{-1} \cdot \left\{ \int_{-x_\eta}^0 dx' + (\log 2\eta) \int_{-x_\eta}^0 |x'| dx' \right\} \\ & \leq C_3. \end{aligned}$$

Concerning the integral with $0 \leq x' \leq x_\eta$, we can estimate in the similar way, by using the fact (21). Thus we see that the second term on the right of (20) is also uniformly bounded for $x_\eta \leq x \leq 1$ and $\eta \geq \max\{C, e\}$.

(II) For $-1 \leq x \leq -x_\eta$, the argument for the estimate is completely parallel if we interchange the roles of $\hat{u}_1(x, \eta)$ and $\hat{u}_2(x, \eta)$. Also for $-x_\eta \leq x \leq x_\eta$, the argument is similar if we rewrite (20) as

$$\begin{aligned} & \int_{-1}^1 |Q(x, x'; \eta)| dx' \\ & \leq |W(\eta)^{-1} \cdot \hat{u}_2(x, \eta)| \cdot \left\{ \int_{x_\eta}^1 |\hat{u}_1(x', \eta)| dx' + \int_x^{x_\eta} |\hat{u}_1(x', \eta)| dx' \right\} \\ & \quad + |W(\eta)^{-1} \cdot \hat{u}_1(x, \eta)| \cdot \left\{ \int_{-1}^{-x_\eta} |\hat{u}_2(x', \eta)| dx' + \int_{-x_\eta}^x |\hat{u}_2(x', \eta)| dx' \right\}, \end{aligned}$$

and estimate the each term on the right hand side. Consequently, we see that the value of $\int_{-1}^1 |Q(x, x'; \eta)| dx'$ is uniformly bounded for $-1 \leq x \leq 1$ and $\eta \geq \max\{C, e\}$.

The argument to show (18) for $m > 0$ is similar to the above if we notice the fact (10) and

$$\frac{d}{dz} \Phi(\alpha, \gamma; z) = \frac{\alpha}{\gamma} \Phi(\alpha+1, \gamma+1; z)$$

(see page 254 of [2]). Thus the proof of Proposition 4 is clear. ■

Now we are in position to verify the hypoellipticity parts of Theorems.

Proof of hypoellipticity: First let us notice that the operator P is elliptic except $x=0$. Hence we can restrict our consideration at $(0, y_0)$. Moreover, since P is non-characteristic with respect to the variable x , the smoothness of the solution w.r.t. the variable x follows from the one w.r.t. the variable y . To be more precise, let $H^{k,l}$ be the space of distributions u satisfying $(1+\xi^2)^{k/2}(1+\eta^2)^{l/2} \hat{u}(\xi, \eta) \in L^2(\mathbf{R}^2)$ (ξ and η are the dual variables of x and y respectively). Then $u \in H^{k,l}$ and $Pu \in C^\infty$ at $(0, y_0)$ implies that $u \in \bigcap_{m=1}^\infty H^{k+2m, l-2m}$ at $(0, y_0)$. Thus $u \in H^{0,\infty}$ and $Pu \in C^\infty$ at $(0, y_0)$ implies that $u \in C^\infty$ at $(0, y_0)$. So it suffices to prove that $u \in H^{0,\infty}$ at $(0, y_0)$ when $Pu \in C^\infty$ at $(0, y_0)$.

Secondly we can assume that the support of the solution is contained in a small neighborhood of $(0, y_0)$. To observe this, let us take a function $\chi(x, y) \in C_0^\infty$ satisfying $\chi(x, y) \equiv 1$ for $|x| + |y - y_0| \leq \delta/2$ and $\chi(x, y) \equiv 0$ for $|x| + |y - y_0| \geq \delta$. Then the second term on the right of

$$Pu = P\chi u + P(1-\chi)u$$

is equal to 0 in a neighborhood of $(0, y_0)$. So it suffices to show that χu is smooth at $(0, y_0)$ provided $P\chi u$ is smooth there.

Now take a function $\phi(\eta) \in C^\infty$ such that $\phi(\eta) \equiv 0$ for $|\eta| \leq \max\{C, e\}$ and $\phi(\eta) \equiv 1$ for $|\eta| \geq 2\max\{C, e\}$, and set

$$(23) \quad Qu(x, y) = \frac{1}{2\pi} \iiint e^{i(y-y')\eta} Q(x, x'; \eta) \phi(\eta) u(x', y') dx' dy' d\eta.$$

Then it follows from (18) and (19) with $m=0$ that Q is a bounded operator $H^{0,l}((-1, 1) \times \mathbf{R})$ for all $l \in \mathbf{R}$. Moreover since $Q(x, x'; \eta)$ is the Green func-

tion of (3), it holds that $PQ=I+K_1$, where K_1 is an operator with symbol $1-\phi(\eta)$, in particular, it is regularizing operator w.r.t. y , i.e., the one from $H^{0,l}((-1,1)\times\mathbf{R})$ into $H^{0,\infty}((-1,1)\times\mathbf{R})$. Now let $Q_1(x, x'; \eta)$ be the Green function of (3) with b being replaced by \bar{b} , and let R be the adjoint operator of (23) with $Q(x, x'; \eta)$ being replaced by $Q_1(x, x'; \eta)$. Then it holds that

$$RP = I + K,$$

where K is an operator from $H^{0,l}((-1,1)\times\mathbf{R})$ into $H^{0,\infty}((-1,1)\times\mathbf{R})$. Furthermore R has pseudo-local property w.r.t. $y \bmod H^{0,\infty}$. To be more precise, let $\chi_1(y)$ be a function of class C_0^∞ satisfying $\chi_1(y) \equiv 1$ for $|y-y_0| \leq \varepsilon$. Then the second term of the right of

$$Rf = R\chi_1 f + R(1-\chi_1)f$$

belongs to $H^{0,\infty}$ at $(0, y_0)$ for any $f \in H^{0,l}$. It is a consequence of (18) and (19). Indeed, $R(1-\chi_1)f$ is expressed as

$$\frac{1}{2\pi} \iint e^{i(y'-y)\eta} dy' d\eta \int Q_2(y, y'; x, x'; \eta) f(x, y') dx',$$

where

$$Q_2(y, y'; x, x'; \eta) = \frac{1-\chi_1(y')}{(y'-y)^m} D_\eta^m \{ \bar{Q}_1(x', x; \eta) \phi(\eta) \},$$

for arbitrary positive integer m , and the values of

$$|\eta|^{m+\gamma} \int |\partial_\eta^\alpha \partial_{y'}^\beta \partial_\eta^\gamma Q_2(y, y'; x, x'; \eta)| dx'$$

and

$$|\eta|^{m+\gamma} \int |\partial_\eta^\alpha \partial_{y'}^\beta \partial_\eta^\gamma Q_2(y, y'; x, x'; \eta)| dx$$

are uniformly bounded for $\eta \in \mathbf{R}$, $|y-y_0| \leq \varepsilon$ and $y' \in \mathbf{R}$. Hence $Q_2(y, y'; \cdot, \cdot; \eta)$ is an operator valued symbol of class $S^{-\infty}((y_0-\varepsilon, y_0+\varepsilon) \times \mathbf{R}_y \times \mathbf{R}_\eta; L^2(-1, 1))$.

Thus, if Pu is of class $H^{0,\infty}$ at $(0, y_0)$ and the supports of χ and χ_1 are taken properly small, then all terms on the right hand side of the equation

$$\begin{aligned} \chi u &= RP\chi u - K\chi u \\ &= R\chi_1 P\chi u + R(1-\chi_1)P\chi u - K\chi u \end{aligned}$$

are of class $H^{0,\infty}$ at $(0, y_0)$, since χu becomes of class $H^{0,l}$ for some l and the singular support of χu becomes contained in $\{(x, y) | x=0\} \cap \text{supp } \chi$. This completes the proof. ■

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