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Main-Chain Stiffness and Helical Conformation of a Poly(quinoxaline-2,3-diyl) in Solution

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(Running Title: Chain Stiffness of a Polyquinoxaline in Solution)

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ABSTRACT: Light and small-angle X-ray scattering and viscosity measurements in tetrahydrofuran at 25 °C were made for nine helical poly[5,8-dimethyl-6,7-bis(propoxymethyl)quinoxaline-2,3-diyl] samples ranging in the weight-average molar mass M_w from 8×10^3 g mol⁻¹ to 6×10^5 g mol⁻¹ to determine the particle scattering function, the radius of gyration, and the intrinsic viscosity as a function of M_w . The dimensional and hydrodynamic properties were consistently explained in terms of the Kratky-Porod wormlike chain. The helix pitch per residue h (or the contour length per residue) and the chain stiffness parameter λ^{-1} (the Kuhn segment length or twice of the persistence length) were estimated to be $h = 0.19$ nm and $\lambda^{-1} = 43 \pm 3$ nm. The former parameter corresponds to an internal rotation angle of about 120° which is substantially the same as the most stable helical structure estimated from the internal rotation potential. The latter one (λ^{-1}) indicates that the poly(quinoxaline-2,3-diyl) has rigid helical main chain of which the dihedral angle between adjacent monomer units are restricted by steric hindrance.

Introduction

Helical polymers¹⁻⁶ have attracted considerable attention because of their unique macromolecular functions leading to chiral separation,⁷⁻⁹ asymmetric catalysis,¹⁰⁻¹¹ and chiroptical materials.¹²⁻¹⁶ To gain deeper insight into their intriguing macromolecular functions, it is particularly important to determine their conformations in dilute solution. So far, physical properties of helical polymers such as polysilanes,¹⁷⁻¹⁸ polyacetylenes,¹⁹ polyisocyanates,²⁰⁻²¹ polyisocyanides,²² and polysaccharide derivatives²³⁻²⁷ have been investigated in dilute solution to reveal their intrinsic properties arising from their helical structure.

Poly(quinoxaline-2,3-diyl)s (PQXs) are a unique class of synthetic helical polymers prepared by living polymerization of 1,2-diisocyanobenzenes with transition metal initiators.²⁸⁻³¹ PQXs have been supposed to adopt rigid helical structure in solution owing to the steric repulsion of the two substituents at the 5- and 8-positions on the quinoxaline ring.^{28,32} Recently, we reported that single-handed PQXs bearing diarylphosphino groups can serve as highly effective chiral ligands

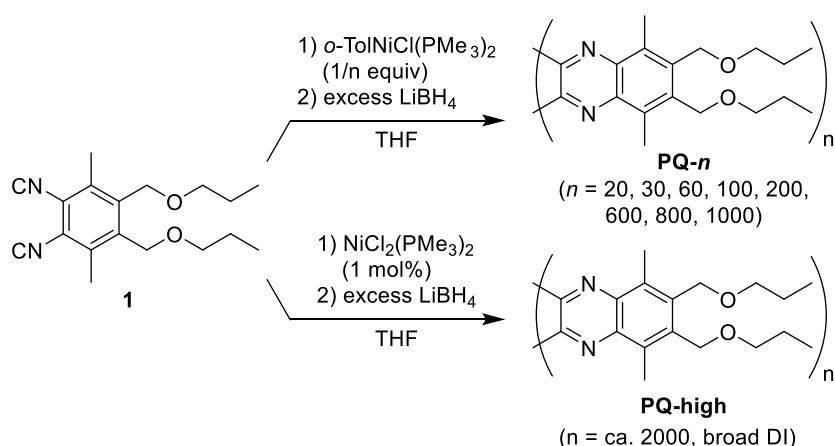
for transition metal catalysts in asymmetric reactions.³³ Furthermore, we also reported solvent-dependent helix inversion of PQXs bearing chiral side chains,³⁴⁻³⁵ which enables the highly enantioselective production of both enantiomers from a single chiral catalyst in various asymmetric reactions.³⁶⁻³⁹ This feature could also be extended to the formation of chirality-switchable cholesteric superstructures⁴⁰ and generation of handedness-switchable circularly polarized luminescence.⁴¹ It should be noted that the rigidity of the polymer main chain plays important roles in these application of PQXs. For instance, a chiral ligand should have a rigid molecular framework to exhibit high enantioselectivity. The chirality-switchable cholesteric materials are also based on the rigid main chain to stabilize the liquid crystalline state. Furthermore, a rigid, chiral, and fluorescent molecular framework is indispensable to exhibit circularly polarized luminescence with high dissymmetry factor. However, it is still difficult to predict the chain stiffness properly only from the chemical structure of semiflexible or rigid polymers because it is determined by rather small fluctuation of the internal rotation and the bond angle.

In order to develop new chiral functional materials based on PQXs, the comprehension of dimensional and hydrodynamic properties of PQXs in dilute solution is thus significantly important. These properties are characterized mainly by the contour length and the chain stiffness parameter λ^{-1} (the Kuhn segment length or twice the persistence length) in terms of the Kratky-Porod wormlike chain⁴² (or more generally, the helical wormlike chain).⁴³⁻⁴⁴ Firstly, the contour length is related to the helical structure of the polymer in solution. Although two helical conformations having different helix pitches were proposed for PQXs,²⁹ we can make it clear through the determination of the contour length of PQXs. Secondly, the chain stiffness parameter is decisively related with solution properties, that is, not only solution viscosity including concentrated solutions⁴⁵⁻⁴⁶ but also liquid crystallinity.⁴⁵

In this paper, our research interest has focused on the chain stiffness and the helix pitch (or helix rise) per residue of poly[5,8-dimethyl-6,7-bis(propoxymethyl)quinoxaline-2,3-diyl], which is the simplest class of PQXs and has been studied in detail in these two decades.⁴⁷ We studied dimensional and hydrodynamic properties of PQXs in tetrahydrofuran (THF) by a size exclusion chromatography (SEC) equipped with a multi-angle laser light scattering (MALS) photometer and a twin capillary viscometer detector and by synchrotron radiation small-angle X-ray scattering (SAXS) measurements.

Experimental Section

Sample Preparation. Eight poly[5,8-dimethyl-6,7-bis(propoxymethyl)quinoxaline-2,3-diyl] samples **PQ-*n*** of which degree of polymerization *n* ranges between 20 and 1000 were synthesized by living polymerizations of monomer **1**⁴⁷ in the presence of an organonickel complex⁴⁸ (Scheme 1). Their dispersity (*D*) defined as the ratio of weight- to number-average molar masses M_w/M_n were between 1.01 and 1.07 determined by the size exclusion chromatography described below. A high molar mass sample **PQ-high** with relatively broad *D* ($M_w/M_n = 1.33$) was prepared by using a simple nickel(II) complex, $\text{NiCl}_2(\text{PMe}_3)_2$, as reported before.⁴⁹ These polymerization proceeded quickly at room temperature and gave corresponding polymers in quantitative yields.



Scheme 1. Polymerization of 1,2-dicyanobenzene monomer **1** with nickel initiators.

Static Light Scattering (SLS) and Viscosity Measurements with Size Exclusion Chromatography (SEC). SLS and viscosity measurements for all samples in THF were made using SEC equipped with a DAWN HELEOS multi-angle laser light scattering photometer with a He-Ne laser of which wavelength λ_0 in vacuum was 658 nm, a VISCOSTAR-II online differential viscometer, and an Optilab rEX differential refractometer (all detectors supplied by Wyatt Technology, USA). These methods are widely used to characterize various polymers in solution including linear and branched polystyrenes and a polysaccharide derivative.^{26,50-52} A SEC system consisting of a LC-20AD isocratic pump, a SIL-20A autosampler supplied by SHIMADZU, Japan, and above mentioned detectors were used with three TSKgel GMH_{XL} columns (TOSOH, Japan) for **PQ-20**, **PQ-30**, ..., and **PQ-200**, with a GMH_{HR}-H column and two GMH_{XL} columns for **PQ-600**, **PQ-800**, and **PQ-1000**, or with two GMH_{XL} and a G2500H_{XL} columns for **PQ-high**. The flow rate was set to 1.0 cm³min⁻¹, the column and detector temperature were kept at room temperature (substantially the same as 25 °C), and the injection volume was 0.2 cm³. The ASTRA software (ver.5.3.3, Wyatt technology, USA) was used to collect and to analyze the data. Differential refractive index increment $\partial n/\partial c$ for **PQ-600**, **PQ-800**, **PQ-1000**, and **PQ-high** in THF was determined by using the peak area of the refractive index chromatogram and the polymer mass concentration c of injected solution assuming full recovery. The $\partial n/\partial c$ value was determined to be 0.196 cm³g⁻¹ at $\lambda_0 = 658$ nm. Angular dependence of the scattering intensity was analyzed using the Berry square-root plot⁵³ in a scattering angle range from 45.8° to 138.8° to estimate the weight-average molar mass M_w and the z -average mean-square radius of gyration $\langle S^2 \rangle_z$ of each fraction. It should be noted that the second virial term does not cause an appreciable error in M_w (and $\langle S^2 \rangle_z$) if we assume the second virial coefficient estimated from SAXS measurements described below. To verify the obtained M_w , the universal calibration method was examined for all **PQ- n** samples with polystyrene standard (Tosoh, Japan) of which M_w ranges from 1.02×10^4 to 8.42×10^6 . The resultant M_w values from the two methods agree with each other within $\pm 2\%$. The SEC chromatograms have shoulder peaks suggesting high molar mass components except for low molar mass samples as displayed in Figure 1. The shoulder peaks were omitted in the discussion of dimensional and hydrodynamic properties described below.

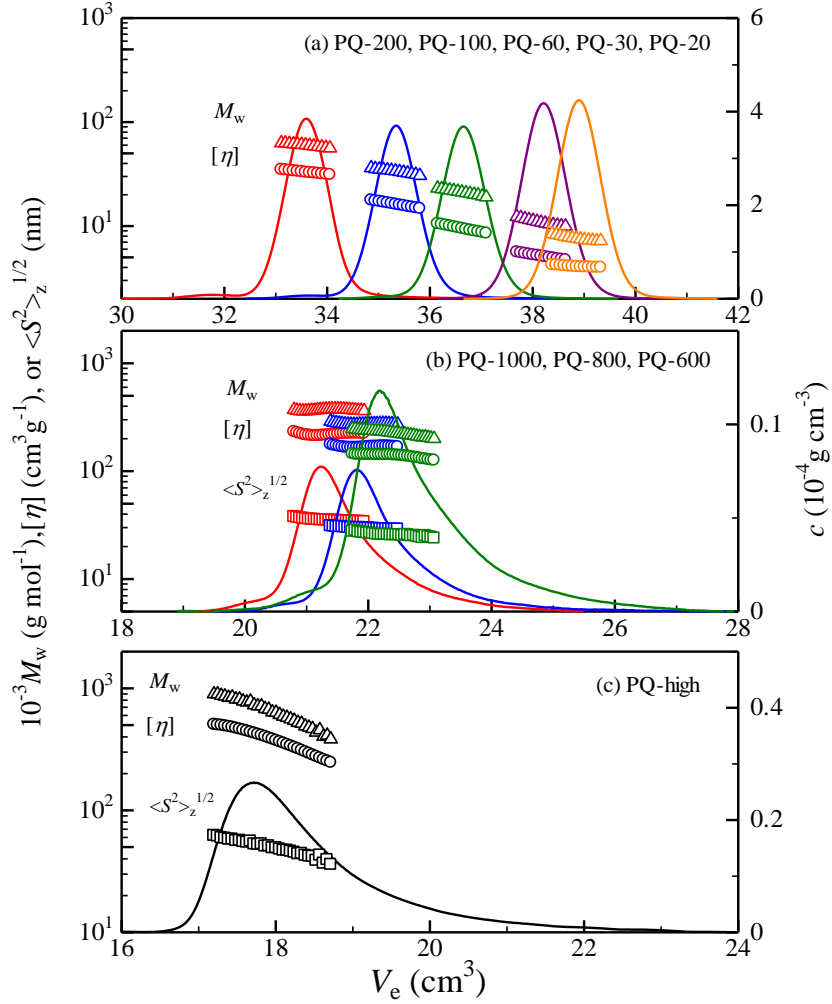


Figure 1. Retention volume V_e of the weight-average molar mass M_w (triangles), the z -average radius of gyration $\langle S^2 \rangle_z^{1/2}$ (squares), the intrinsic viscosity $[\eta]$ (circles) and polymer mass concentration c (solid curves) in THF.

Small angle X-ray Scattering (SAXS). Synchrotron radiation SAXS experiments for **PQ-20**, **PQ-30**, **PQ-60**, **PQ-100**, and **PQ-200** in THF at 25 °C were carried out at the BL-10C beamline in KEK-PF (Ibaraki, Japan). The camera length, λ_0 , and the accumulation time were set to be 980 mm, 0.1488 nm, and 300 sec, respectively. The scattered light was detected by a PILATUS3-300KW hybrid photon counting detector (DECTRIS, Switzerland). Silver behenate was used to determine the beam center and the camera length. The circular average method was utilized to determine the scattering intensity $I(q)$ as a function of the magnitude q of scattering vector. The excess scattering intensity $\Delta I(q)$ of each solution was determined from the difference between $I(q)$ for the solution and the solvent in the same capillary cell taking intensity and transparency of the incident X-ray into account. Four solutions with different polymer mass concentration c ranging in 3×10^{-3} g cm $^{-3}$ and 1.5×10^{-2} g cm $^{-3}$ was measured to extrapolate $[c/\Delta I(q)]^{1/2}$ to infinite dilution. The $P(q)$ as well as $\langle S^2 \rangle_z$ were determined by means of the extrapolation of $[c/\Delta I(q)]_{c=0}^{1/2}$ to zero q^2 . Since the $[c/\Delta I(q)]_{c=0}^{1/2}$ data for **PQ-100** and **PQ-200** were hard to extrapolate to $q^2 = 0$, the $[c/\Delta I(0)]_{c=0}^{1/2}$ values were estimated from those for **PQ-20**, **PQ-30**, and **PQ-60** considering that

$[c/\Delta I(0)]_{c=0}^{1/2}$ should be proportional to $M_w^{-1/2}$. The $\langle S^2 \rangle_z$ values were determined from the initial slope of the Berry plots as shown in Figure 2. Slight discrepancy between the experimental data and the solid curve for **PQ-100** at low- q region might be due to the small amount ($\sim 1\%$) of larger molar mass component which was observed in the SEC chromatogram. We note that values of $[c/\Delta I(0)]_{c=0}^{1/2}$ increased with increasing concentration and then the second virial coefficients were estimated to be $1.2 \times 10^{-3} \text{ mol g}^{-2} \text{ cm}^3$ and $3 \times 10^{-4} \text{ mol g}^{-2} \text{ cm}^3$ for PQ-20 and PQ-200, respectively. Preliminary experiments for **PQ- n** samples in 1,2-dichloroethane and 1,1,2-trichloroethane in which helix inversion was observed³⁴⁻³⁵ were also performed at the BL40B2 beamline in SPring-8 with the incident light of 0.1 nm wavelength. There was almost no scattering intensity owing to the very low transparency of the solvents.

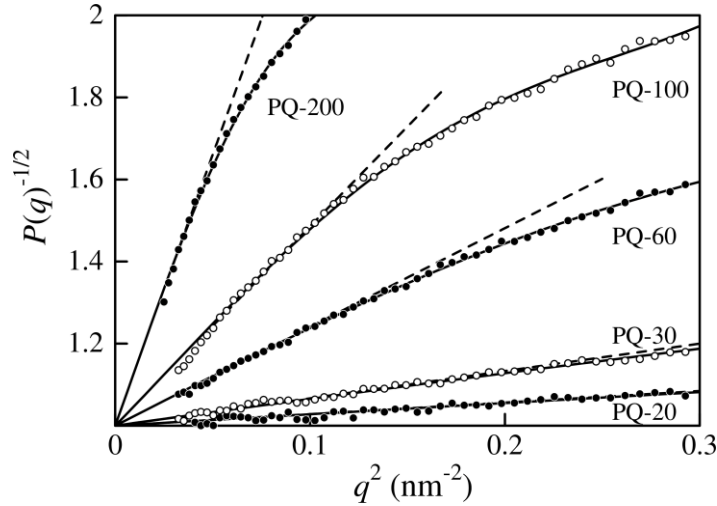


Figure 2. Berry plots for indicated **PQ- n** samples in THF at 25 °C at infinite dilution. Dashed lines indicate the initial slopes.

Conformational Calculations. The conformational energy calculations for poly(5,6,7,8-tetramethylquinoxaline-2,3-diyl) were carried out by using the method of MOPAC PM6 to estimate internal rotation potential of **PQ- n** . They were performed with MOPAC2012 software⁵⁴ with parameters of PM6 EF GNORM=0.1 GEO-OK. The total energy of quinoxaline 20-mers was calculated by varying the dihedral angle ϕ , between two adjacent quinoxaline units from 40° to 140° with an interval of 5°.

Results and Discussion

Dimensional and Hydrodynamic Properties. The resultant data for M_w , M_w/M_n , $\langle S^2 \rangle_z^{1/2}$, and the intrinsic viscosity $[\eta]$ in THF at 25 °C are summarized in Table 1. The $\langle S^2 \rangle_z^{1/2}$ and $[\eta]$ values were determined by the integration of the peak area of each SEC chromatogram except for **PQ-high**. The Flory viscosity factor Φ defined as the following equation

$$\Phi = \frac{[\eta]M_w}{(6\langle S^2 \rangle_z)^{2/3}} \quad (1)$$

are also listed in Table 1. The Φ values ($1.1 \times 10^{23} \text{ mol}^{-1}$, in average) for higher M_w samples ($M_w > 2 \times 10^5 \text{ g mol}^{-1}$) are appreciably smaller than those for typical flexible polymers ($\sim 2.6 \times 10^{23} \text{ mol}^{-1}$)^{43,55} and rather close to that for a typical semiflexible poly(*n*-hexylisocyanate) (0.8 to $1.4 \times 10^{23} \text{ mol}^{-1}$),⁵⁶⁻⁵⁷ indicating the semiflexible nature of the **PQ-*n*** in THF.

Table 1. Molecular Characteristics of Poly[5,8-dimethyl-6,7-bis(propoxymethyl)quinoxaline-2,3-diyl] Samples and Physical Properties in THF at 25 °C or Room Temperature

| Sample | $M_w (10^{-3} \text{ g mol}^{-1})^a$ | M_w / M_n | $\langle S^2 \rangle_z^{1/2} (\text{nm})$ | $[\eta] (\text{cm}^3 \text{ g}^{-1})$ | $\Phi (10^{23} \text{ mol}^{-1})$ |
|---------|--------------------------------------|-------------|---|---------------------------------------|-----------------------------------|
| PQ-20 | 7.83 | 1.01 | 1.30 ^b | 4.1 | 9.9 |
| PQ-30 | 11.4 | 1.02 | 2.0 ^b | 5.1 | 4.9 |
| PQ-60 | 21.3 | 1.01 | 3.8 ^b | 9.5 | 2.5 |
| PQ-100 | 34.0 | 1.01 | 5.4 ^b | 16.3 | 2.4 |
| PQ-200 | 59.7 | 1.02 | 8.9 ^b | 33.2 | 1.91 |
| PQ-600 | 227 | 1.07 | 27.3 ^c | 132 | 1.00 |
| PQ-800 | 286 | 1.06 | 31.6 ^c | 167 | 1.03 |
| PQ-1000 | 371 | 1.03 | 36.0 ^c | 219 | 1.18 |
| PQ-high | 574 | 1.33 | — ^d | — ^d | — ^d |

^a Average value from the two methods (see text). ^b From SAXS. ^c From Light scattering. ^d Not determined.

Figure 3 illustrates reduced Holtzer plots for **PQ-200**, **PQ-100**, **PQ-60**, **PQ-30**, and **PQ-20** determined by SAXS. Data points are mostly flat except for the low- q region indicating that the stiff main chain and weak contribution from the chain cross-section. Indeed, if we estimate the chain diameter d from the cross sectional plots of $[\ln qP(q) \text{ vs } q^2]$ on the basis of the equation $[P(q) = P_0(q) \exp(-k^2 d^2/16)]$,⁵⁸ the d value was evaluated to be about 0.4 nm, where $P_0(q)$ is the scattering function of the chain contour. According to Nakamura and Norisuye, this plot has some certain linear region not only for straight rods but also wormlike chains.⁵⁹ This small value may be reasonable because the chain thickness effect in $P(q)$ from SAXS reflects electron density profile of the chain cross section including solvent molecules.^{43,60-61} We therefore analyzed the $P(q)$ data in terms of the thin wormlike chain with the following equation

$$P(q) = \frac{2}{L^2} \int_0^L (L-t) I(q;t) dt \quad (2)$$

where L and $I(q; t)$ are the contour length and the characteristic function of the wormlike chain, respectively. It should be noted that the excluded-volume effects are negligible for the current $P(q)$ data because the Kuhn segment number is calculated to be 0.88 even for **PQ-200**, which is the highest molar mass sample to determine $P(q)$. We utilized the Nakamura-Norisuye expression for $I(q; t)$.^{59,62} The theoretical $P(q)$ can be calculated for the required M_w with the Kuhn segment length λ^{-1} and the helix pitch per residue h (or the contour length per residue) which is related to L as

$$h = \frac{LM_0}{M_w} \quad (3)$$

where M_0 denotes the molar mass of the repeat unit ($= 300.2 \text{ g mol}^{-1}$). A curve fitting procedure was examined for the data to determine h . The resultant theoretical values well explain the experimental data except for the low- q region for **PQ-100** and **PQ-200**. This is most likely due to the larger molar mass component as described in the Experimental section. These features are also seen in the double logarithmic plots in Figure S9 in the supporting information. The parameter h slightly decreased with lowering M_w and the value for **PQ-20** was 17% smaller than that for **PQ-200**. This is likely because of the end effect. In this case, the contour length can be rewritten as

$$L = \frac{M_w}{M_0} h_0 + \delta \quad \text{or} \quad h = h_0 + \frac{M_0}{M_w} \delta \quad (4)$$

where h_0 and δ were the helix pitch per residue for infinitely long chain and difference in the contour lengths nearby the chain ends. Since good linearity of the plots of h vs M_w^{-1} as displayed in Figure 3(b) supports this suggestion, the two parameters were estimated to be $h_0 = 0.194 \text{ nm}$ and $\delta = -0.80 \text{ nm}$ from the line in the figure. This end effect might be due to the low electron density of side alkyl groups nearby the chain ends taking very small chain diameter (0.4 nm) into consideration. On the other hand, the chain stiffness only from the $P(q)$ data cannot be estimated because the theoretical dashed curve for the rigid rod fairly explains the experimental data including low- q region. In any case, they are well explained by the theoretical solid red curves calculated for the wormlike chain with the above h and $\lambda^{-1} = 43 \text{ nm}$ which was obtained from $\langle S^2 \rangle_z$ and $[\eta]$ described below. We should however consider that the parameter h is significantly affected by the experimental error of M_w . The above mentioned M_w dependent h might be an artifact if the M_w values for low molar mass samples were overestimated by some reason such as slight optical anisotropy or molecular weight dependent $\partial n / \partial c$. Also, if it is due to the low electron density of both chain ends, $[\eta]$ should be free from the effect. We therefore analyzed the $\langle S^2 \rangle_z$ and $[\eta]$ data both with and without taking the chain end effects into account as described below.

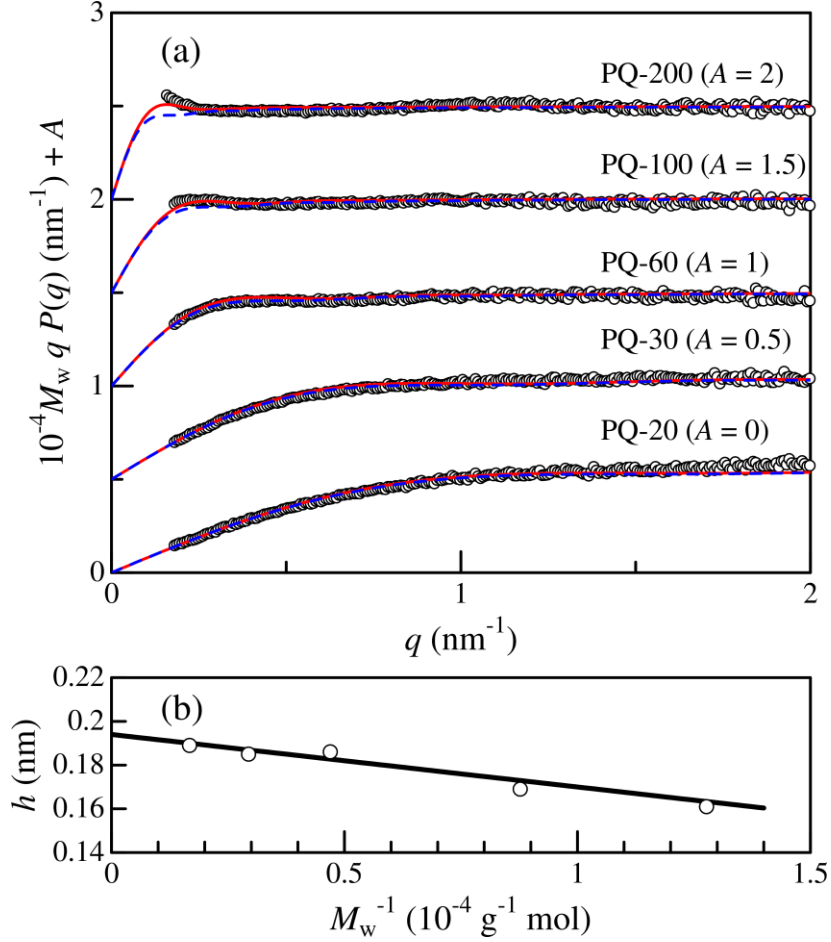


Figure 3. (a) Reduced Holtzer plots for indicated **PQ-*n*** samples in THF at 25 °C. Solid red curves and dashed blue curves denote the theoretical values for the wormlike chain and the rigid rod with the parameters in the main text. The ordinate values are shifted by A for clarity. (b) Plots of h vs M_w^{-1} for **PQ-*n*** in THF at 25 °C.

The M_w dependence of $\langle S^2 \rangle_z$ is displayed in Figure 4. While the data points for three lower M_w samples are fitted by a straight dot-dashed line for the rigid rod with $h = 0.184$ nm, those for the higher M_w samples are lower than the dot-dashed line. It is a typical feature for the wormlike chain having quite high chain stiffness. For the Kratky-Porod wormlike chain, the radius of gyration $\langle S^2 \rangle$ can be calculated by the Benoit-Doty equation as⁶³

$$\langle S^2 \rangle = \frac{L}{6\lambda} - \frac{1}{4\lambda^2} + \frac{1}{4\lambda^3 L} - \frac{1}{8\lambda^4 L^2} [1 - \exp(-2\lambda L)] \quad (5)$$

If we choose $h = 0.184$ nm and $\lambda^{-1} = 45 \pm 1$ nm, the solid curve quantitatively reproduces the experimental data. The Kuhn segment number λL for the highest M_w ($= 9.5 \times 10^5$) fraction is 13.2 and hence the intramolecular excluded-volume effect is insignificant⁶⁴⁻⁶⁵ in terms of the quasi-two-parameter theory^{43,66-67} with the Domb-Barret equation.⁶⁸ If we consider M_w dependent h with eq 4, slight smaller $\lambda^{-1} = 42$ nm was determined. The calculated dashed curve also reproduces the experimental data.

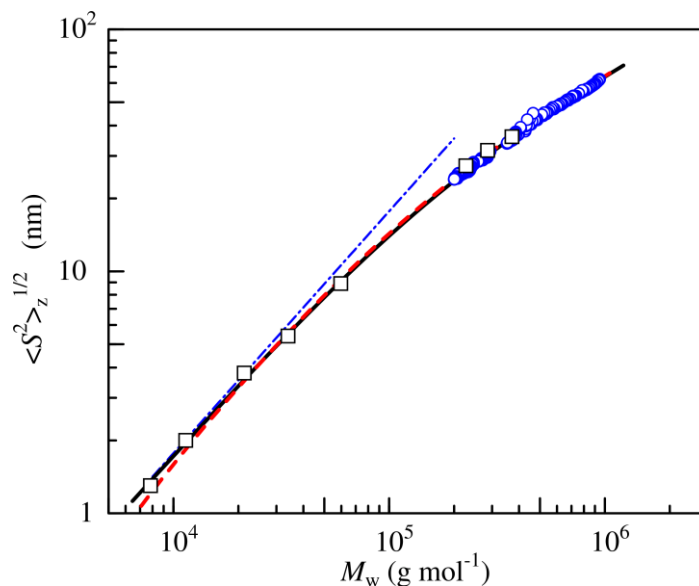


Figure 4. Double logarithmic plots of $\langle S^2 \rangle_z^{1/2}$ vs M_w for **PQ-*n*** samples in THF at room temperature. Squares, average values for each sample listed in Table 1; circles, data of each fraction from SEC. Solid and dashed curves indicate the theoretical values for the wormlike chain with fixed and M_w dependent h , respectively (see text). A dot-dashed line is for the rigid rod with $h = 0.184$ nm.

Intrinsic viscosities $[\eta]$ for the **PQ-*n*** samples in THF are plotted against M_w in Figure 5. The Yamakawa-Fujii-Yoshizaki theory^{43,69-70} allows us to calculate $[\eta]$ for the wormlike cylinder with the three parameters of h , λ^{-1} , and the chain diameter d . Considering the end effects might affect $[\eta]$ for **PQ-20** and **PQ-30** as is the case with $P(q)$, the three parameters cannot uniquely be determined from the curve fitting procedure from the seven higher M_w samples. If we thus assume h to be 0.186 nm which is the average value from $P(q)$ and $\langle S^2 \rangle_z$ for high M_w samples, the rest two parameters are estimated to be $\lambda^{-1} = 45$ nm and $d = 2.3$ nm. The resultant theoretical values in Figure 5 are fitted by the experimental data including lower M_w samples. On the one hand, if we choose eq 4 for the parameter L , theoretical values calculated with slightly different parameters, $\lambda^{-1} = 40$ nm and $d = 2.6$ nm, also reproduce the experimental data. These parameters are summarized in Table 2. Not significant difference were found both for the chain stiffness and the chain diameter from the two methods. Consequently, we may conclude that the wormlike chain parameters of poly[5,8-dimethyl-6,7-bis(propoxymethyl)quinoxaline-2,3-diyl] are as follows: $h = 0.19$ nm, $\lambda^{-1} = 43 \pm 3$ nm, and $d = 2.4 \pm 0.2$ nm.

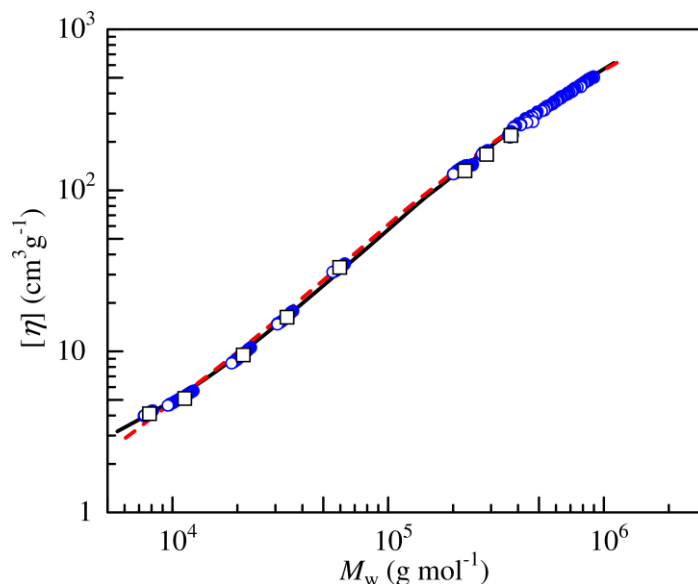


Figure 5. Double logarithmic plots of $[\eta]$ vs M_w for **PQ-*n*** samples in THF at room temperature. Squares, average values listed in Table 1; circles, data of each fraction from SEC. A solid and dashed curves indicate the theoretical values for the wormlike chain with fixed and M_w dependent h , respectively.

Table 2. Wormlike Chain Parameters of poly[5,8-dimethyl-6,7-bis(propoxymethyl)quinoxaline-2,3-diyl] in THF at 25 °C or Room Temperature

| Method | h (nm) | h_0 (nm) | δ (nm) | λ^{-1} (nm) | d (nm) ^b |
|---------------------|--------------------------------|-------------------|-----------------|---------------------|-----------------------|
| Fixed h | 0.186 ± 0.003 ^a | | | 44 ± 2 | 2.3 |
| M_w dependent h | | 0.194 ± 0.002 | -0.8 ± 0.02 | 41 ± 2 | 2.6 |

^a Except for **PQ-20** and **PQ-30**. ^b From $[\eta]$.

As mentioned in the Introduction, the helix pitch per residue h is related with the helical structure. In a previous paper,²⁹ the authors carried out the calculation of 20-mers with fixed dihedral angle ψ ($= 0^\circ$) to reduce computational complexity, and a minimum of the total energy was observed at $\phi = 135^\circ$. In this study, the total energy of quinoxaline 20-mers was calculated by varying the dihedral angle ϕ , between two adjacent quinoxaline units from 40 to 140° with an interval of 5° , along with optimizations of the dihedral angle ψ , bond lengths, and other angles by the semi-empirical molecular orbital method (MOPAC PM6). For a fixed internal rotation angle, this h value can be obtained as a function of ϕ as shown in Figure 6 in which $h = 0.19$ nm corresponds to $\phi = 120^\circ$. If we assume the evaluation error of h to be ± 0.005 nm, the resultant range of the dihedral angle between 113° and 130° . It is fairly consistent with that for minimum of the internal rotation potential. Thus, **PQ-*n*** has a rather extended helical conformation in THF as illustrated in Figure 7a. In sharp contrast, it was reported that helical conformations of poly(*o*-phenylene)s are rather similar to Figure 7b based on the single crystal X-Ray diffraction for the model oligomer and NMR measurements.⁷¹ The characteristic feature of PQXs may be attributed to the unique helical conformation in dilute solution.

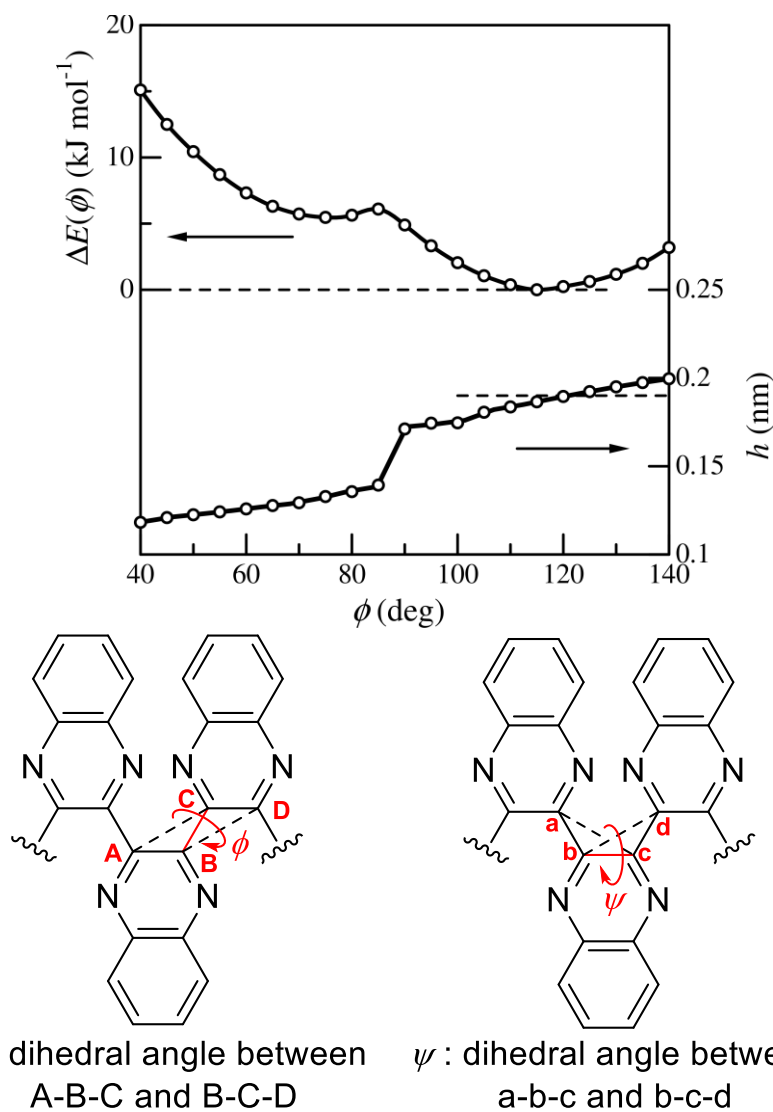


Figure 6. Dependence of the excess internal rotation potential $\Delta E(\phi)$ and h on dihedral angle ϕ for poly[5,6,7,8-tetramethylquinoxaline-2,3-diyl] (20-mer). Dihedral angle ψ , bond lengths, and other angles were optimized by the semi-empirical molecular orbital method.

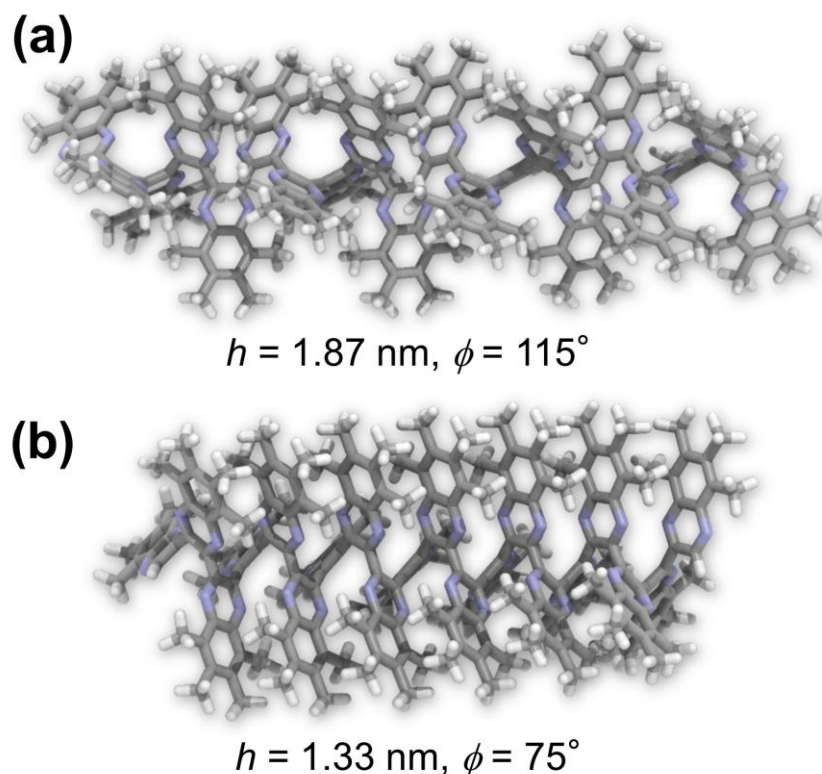


Figure 7. Plausible helical conformations of PQXs with dihedral angles (a) $\phi = 115^\circ$ and (b) $\phi = 75^\circ$. Alkoxy side chains were omitted for clarity.

The determined chain stiffness of 43 nm is close to typical semiflexible polymers, such as poly(*n*-hexylisocyanate) ($\lambda^{-1} = 42 - 84 \text{ nm}$ depending on solvent),^{56,72-73} but somewhat shorter than those for α -helical polypeptides.⁷⁴⁻⁷⁵ However, the chain stiffness of semiflexible polymers may be affected significantly by the side group and interactions with solvent molecules. For example, λ^{-1} values for polyisocyanates are reported to be from 6 nm to 84 nm,²¹ those for polysilanes are between 2 nm and 210 nm,⁷⁶⁻⁷⁷ and those for amylose carbamate derivatives from 9 nm to 75 nm.^{25,27,78} Furthermore, according to Sato et al.,^{18,79} the chain stiffness may also be affected by the helix reversal for some helical polymers. Thus, investigation of PQXs consisting of optically active side groups would allow elucidation of the origin of the chain stiffness.

Conclusions

We determined the helix pitch per residue of 0.19 nm and the Kuhn segment length of $43 \pm 3 \text{ nm}$ for poly[5,8-dimethyl-6,7-bis(propoxymethyl)quinoxaline-2,3-diyl] in THF at room temperature, indicating this polymer behaves as typical semiflexible polymer in solution. The former parameter is consistent with one of the helical conformation predicted by the semi-empirical molecular orbital method, suggesting that the PQXs has helically aligned quinoxaline rings with the dihedral angle $\phi = 113^\circ$ to 130° in dilute solution. Although helical conformations of poly(*o*-arene)s such as poly(*o*-phenylene)s,^{71,80-82} poly(*o*-naphthylene)s⁸³⁻⁸⁴ and poly(*o*-quinoline)s⁸⁵ are attracting much interest, methods for the elucidation of the conformation of poly(*o*-arene)s is still limited to single crystal X-ray diffraction for the oligomeric model compound and NMR measurement. We believe that an approach using light or X-ray scattering measurements of polymer solutions in combination with the theoretical calculations in

conformational analysis can provide powerful tools to reveal their helical conformations in dilute solutions.

Supporting Information

Additional experimental procedures, ¹H NMR spectra, and additional plots for the SAXS data. This material is available free of charge via the Internet at <http://pubs.acs.org>.

Notes

The authors declare no competing financial interest.

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