

Title	Analytic smoothing effect for solutions to Schrödinger equations with nonlinearity of integral type
Author(s)	Ozawa, T; Yamauchi, K; Yamazaki, Y
Citation	Osaka Journal of Mathematics. 2005, 42(4), p. 737-750
Version Type	VoR
URL	https://doi.org/10.18910/9186
rights	
Note	

Osaka University Knowledge Archive : OUKA

<https://ir.library.osaka-u.ac.jp/>

Osaka University

ANALYTIC SMOOTHING EFFECT FOR SOLUTIONS TO SCHRÖDINGER EQUATIONS WITH NONLINEARITY OF INTEGRAL TYPE¹

Dedicated to Professor Mitsuhiro Nakao on the occasion of his sixtieth birthday

T. OZAWA, K. YAMAUCHI and Y. YAMAZAKI²

(Received May 10, 2004)

Abstract

We study analytic smoothing effects for solutions to the Cauchy problem for the Schrödinger equation with interaction described by the integral of the intensity with respect to one direction in two space dimensions. The only assumption on the Cauchy data is the weight condition of exponential type and no regularity assumption is imposed.

1. Introduction

We study the nonlinear Schrödinger equation

$$(1.1) \quad i\partial_t u + \frac{1}{2}\Delta u = f(u),$$

where u is a complex-valued function of time and space variables denoted respectively by $t \in \mathbb{R}$ and $(x, y) \in \mathbb{R}^2$, $\partial_t = \partial/\partial t$, Δ is the Laplacian in \mathbb{R}^2 , and $f(u)$ is the nonlinear interaction given by

$$(1.2) \quad (f(u))(t, x, y) = \lambda \left(\int_{-\infty}^x |u(t, x', y)|^2 dx' \right) u(t, x, y)$$

with $\lambda \in \mathbb{C}$. The equation (1.1) with integral type nonlinearity (1.2) appears as a model of propagation of laser beams under the influence of a steady transverse wind along the x -axis [1, 3, 37] and as a special case of the Davey-Stewartson system where the velocity potential is independent of y -variable [2, 5, 6, 9, 13, 14, 18, 25, 32].

In spite of a large literature on the nonlinear Schrödinger equations (see for instance [4] and references therein), there are not many papers on the equation (1.1) with nonlinearity of integral type [1, 3, 22, 34, 37]. The existence and uniqueness of global solutions to the Cauchy problem for (1.1) is proved in the usual Sobolev spaces

¹Partially supported by Grant-in-Aid for formation of COE.

²JSPS Fellow

$H^m(\mathbb{R}^2)$ with integers $m \geq 1$ [3] and in the Lebesgue spaces $L^2(\mathbb{R}^2)$ [22, 34]. The existence of modified wave operators is proved on a dense set of small and sufficiently regular asymptotic states [22]. Smoothing properties and large time asymptotics are studied in [34] (see also [10, 12, 15, 18, 21, 23–26]). The purpose of this paper is to describe analytic smoothing properties of solutions to the Cauchy problem for (1.1) in terms of the generators of Galilei and pseudo-conformal transformations. We follow the method of Hayashi and coauthors ([10–26], especially [16, 17, 23, 24]) basically, while a systematic use of Strichartz estimates and a couple of observations on the weight condition of exponential type are new ingredients in this paper.

To state our results precisely, we introduce the following notation. The generators of Galilei transformations are denoted by $J = (J_x, J_y) = (x + it\partial_x, y + it\partial_y) = (x, y) + it\nabla$. The generator of pseudo-conformal transformations is denoted by $K = x^2 + y^2 + 2it + 2itx\partial_x + 2ity\partial_y + 2it^2\partial_t = J^2 + 2t^2L$, where $L = i\partial_t + (1/2)\Delta$. Let X be a Banach space of functions of $(t, (x, y)) \in \mathbb{R} \times \mathbb{R}^2$ or of $(x, y) \in \mathbb{R}^2$ and let A be an operator on X . Then for $a > 0$ the space $G^a(A; X)$ is defined at least formally by

$$G^a(A; X) = \left\{ f \in X; \|f; G^a(A; X)\| = \sum_{n \geq 0} \frac{a^n}{n!} \|A^n f; X\| < \infty \right\}.$$

Similarly, for operators $A = (A_1, A_2)$ of two components, we define

$$G^a(A; X) = \left\{ f \in X; \|f; G^a(A; X)\| = \sum_{\alpha \geq 0} \frac{a^{|\alpha|}}{\alpha!} \|A_1^{\alpha_1} A_2^{\alpha_2} f; X\| < \infty \right\},$$

where we have used the standard multi-index notation with $\alpha = (\alpha_1, \alpha_2)$. For simplicity, we write $G^a(A, B; X) = G^a(B; G^a(A; X))$. For $T > 0$ we define

$$\begin{aligned} X_T^a &= G^a((J_x, 0); L^\infty(0, T; L^2) \cap L^8(0, T; L_y^4 L_x^2)), \\ Y_T^a &= G^a((0, J_y); L^\infty(0, T; L^2) \cap L^8(0, T; L_y^4 L_x^2)), \\ Z_T^a &= G^a(J; L^\infty(0, T; L^2) \cap L^8(0, T; L_y^4 L_x^2)), \\ W_T^a &= G^a(J, K; L^\infty(0, T; L^2) \cap L^8(0, T; L_y^4 L_x^2)). \end{aligned}$$

We state the main result in this paper.

Theorem 1. *Let $a > 0$. Then:*

- (1) *For any $\rho > 0$ there exists $T > 0$ independent of a with the following properties:*
 - (a) *For any $\phi \in G^a((x, 0); L^2)$ with $\|\phi; G^a((x, 0); L^2)\| \leq \rho$ (1.1) has a unique solution $u \in X_T^a$.*
 - (b) *For any $\phi \in G^a((0, y); L^2)$ with $\|\phi; G^a((0, y); L^2)\| \leq \rho$ (1.1) has a unique solution $u \in Y_T^a$.*

- (c) For any $\phi \in G^a((x, y); L^2)$ with $\|\phi; G^a((x, y); L^2)\| \leq \rho$
 (1.1) has a unique solution $u \in Z_T^a$.
- (2) For any $\rho > 0$ there exists $T > 0$ depending on a and ρ such that for any $\phi \in G^a((x, y), x^2 + y^2; L^2)$ with $\|\phi; G^a((x, y), x^2 + y^2; L^2)\| \leq \rho$
 (1.1) has a unique solution $u \in W_T^a$.

REMARK 1. Theorem 1 is regarded as an infinite version of Theorem 2 of [34]. The conclusion holds when the time interval is replaced by $[-T, 0]$.

REMARK 2. Functions in the space $G^a(J; L^\infty(0, T; L^2))$ [resp. $G^a((J_x, 0); L^\infty(0, T; L^2))$, $G^a((0, J_y); L^\infty(0, T; L^2))$] are analytic in (x, y) [resp. x, y] for each $t \neq 0$ [23, 24]. Functions in the space $G^a(J, K; L^\infty(0, T; L^2))$ are analytic in $(t, (x, y))$ with $t \neq 0$ [17]. In those respects, Theorem 1 describes analytic smoothing properties of solutions. We note that no regularity assumption is imposed on the Cauchy data.

In Theorem 1, the assumptions such as $\phi \in G^a((x, 0); L^2)$ are described in terms of infinite series of weighted L^2 norms. It would be more explicit to describe the assumptions in terms of weights of exponential type.

The following proposition describes norm in the spaces $G^a((x, y); L^2)$ and $G^a((x, y), x^2 + y^2; L^2)$ in terms of weights of exponential type.

Proposition 1. For any $\varepsilon > 0$ there exists $C_\varepsilon > 0$ such that the following estimates hold:

$$(1.3) \quad \begin{aligned} \|e^{a(|x|+|y|)}\phi; L^2\| &\leq \|\phi; G^a((x, y); L^2)\| \\ &\leq C_\varepsilon \|e^{(a+\varepsilon)(|x|+|y|)}\phi; L^2\|, \end{aligned}$$

$$(1.4) \quad \begin{aligned} \|e^{a(x^2+y^2)}\phi; L^2\| &\leq \|\phi; G^a((x, y), x^2 + y^2; L^2)\| \\ &\leq C_\varepsilon \|e^{(a+\varepsilon)(x^2+y^2)}\phi; L^2\|. \end{aligned}$$

Moreover, the second inequality in (1.4) is optimal in the sense that the estimate

$$(1.5) \quad \|\phi; G^a((x, y), x^2 + y^2; L^2)\| \leq C \|e^{a(x^2+y^2)}\phi; L^2\|$$

fails to hold.

REMARK 3. The first part of Proposition 1, including (1.3) and (1.4), is a special case of Proposition 2 in Section 4 below.

In [17], Hayashi and Kato proved analyticity in space-time with $t \neq 0$ of solutions for the nonlinear Schrödinger equations

$$i\partial_t u + \frac{1}{2}\Delta u = \lambda|u|^{2p}u$$

in $\mathbb{R} \times \mathbb{R}^n$ with $\lambda \in \mathbb{C}$ and nonnegative integer p under the assumption on the Cauchy data ϕ such that $e^{i|\cdot|^2}\phi \in H^{[n/2]+1}$ for $n \geq 1$, where $[r]$ is the integer part of $r > 0$. The above assumption is relaxed as $e^{i|\cdot|^2}\phi \in H^m$, where $m = 0$ if $(n, p) = (1, 1)$ and $m = 1$ if $n = 2$ or $(n, p) = (3, 1)$ (see [17]). In view of (1.4) the assumption $\phi \in G^a((x, y), x^2 + y^2; L^2)$ is satisfied if $e^{(a+\varepsilon)(x^2+y^2)}\phi \in L^2$, which corresponds to $m = 0 < a < a + \varepsilon = 1$ in the last assumption in $n = 2$.

We refer the reader to [7, 8, 11, 19, 21, 29, 30, 31] for analyticity of solutions to other nonlinear evolution equations and to [33, 35, 36, 38] for analytic smoothing effects for linear dispersive equations.

We prove Theorem 1 in Section 3 by a contraction argument. Basic estimates for the proof of Theorem 1 are summarized in Section 2. We prove Proposition 1 in Section 4 in a general setting.

Throughout the paper we use the following notation without further comments.

$L_y^p L_x^q = L^p(\mathbb{R}_y; L^q(\mathbb{R}_x))$ with norm

$$\|u; L_y^p L_x^q\| = \|\|u; L_x^q\|; L_y^p\|.$$

$L^p = L^p(\mathbb{R}^2) = L_y^p L_x^p$. $U(t) = \exp(i(t/2)\Delta)$ denotes the free Schrödinger group acting on functions on \mathbb{R}^2 . $M(t)$ denotes the modulation operator realized as the multiplication by $\exp(i(x^2 + y^2)/(2t)) \cdot$ for $t \neq 0$. The operators J are represented as

$$J = U(t)(x, y)U(-t) = M(t)it\nabla M(-t),$$

while K satisfies the following useful identity

$$K - 2it = 2itM(t)PM(-t),$$

where $P = x\partial_x + y\partial_y + t\partial_t$.

2. Preliminaries

In this section we collect some basic estimates for the free Schrödinger group $U(t)$ and the nonlinearity $f(u)$ of integral type.

Lemma 1 (Hayashi-Ozawa [22]). *$U(t)$ satisfies the following estimates:*

- (1) For any (q, r) with $0 \leq 2/q = 1/2 - 1/r \leq 1/2$

$$\|U(\cdot)\phi; L^q(\mathbb{R}; L_y^r L_x^2)\| \leq C \|\phi; L^2\|.$$

(2) For any (q_j, r_j) with $0 \leq 2/q_j = 1/2 - 1/r_j \leq 1/2$, $j = 1, 2$, the operator G defined by

$$(Gu)(t) = \int_0^t U(t-t')u(t') dt'$$

satisfies the estimate

$$\|Gu; L^{q_1}(0, T; L_y^{r_1} L_x^2)\| \leq C \|Gu; L^{q'_2}(0, T; L_y^{r'_2} L_x^2)\|,$$

where C is independent of $T > 0$ and p' is the dual exponent to p defined by $1/p + 1/p' = 1$.

Proof. See [22, 34]. □

Lemma 2 (Hayashi-K. Kato [17]). *Let $1 \leq q, r \leq \infty$ and let $p \in \mathbb{R}$. Then for any $a, T > 0$ and $\mu \in \mathbb{C}$ with $amT < 1$, $m = \text{Max}(|\mu|/2, 1) + 1$ the following inequality holds:*

$$\|f; G^a(K + i(p-2)t + \mu t; X)\| \leq \left(1 + \frac{a|\mu|T}{1 - amT}\right) \|f; G^a(K + i(p-2)t; X)\|,$$

where $X = L^q(0, T; L_y^r L_x^2)$ or $G^a(J; L^q(0, T; L_y^r L_x^2))$.

Proof. We argue as in [17]. By the commutation relation

$$[\mu t, K + i(p-2)t] = -\frac{2i}{\mu}(\mu t)^2,$$

we have

$$\begin{aligned} & (K + i(p-2)t + \mu t)^l \\ &= \sum_{k=1}^l \binom{l}{k} \left[\prod_{j=0}^{k-1} (\mu + 2ij) \right] t^k (K + i(p-2)t)^{l-k} \\ (2.1) \quad & + (K + i(p-2)t)^l. \end{aligned}$$

We note that $\|tf; X\| \leq T \|f; X\|$. By (2.1), we obtain

$$\begin{aligned} & \|f; G^a(K + i(p-2)t + \mu t; X)\| \\ &= \sum_{l=0}^{\infty} \frac{a^l}{l!} \|(K + i(p-2)t + \mu t)^l f; X\| \end{aligned}$$

$$\begin{aligned}
&\leq \sum_{l=0}^{\infty} \sum_{k=1}^l \frac{a^l}{l!} \binom{l}{k} \left[\prod_{j=0}^{k-1} (|\mu| + 2j) \right] T^k \|(K + i(p-2)t)^{l-k} f; X\| \\
&\quad + \|f; G^a(K + i(p-2)t; X)\| \\
&\leq \sum_{l=0}^{\infty} \sum_{k=1}^l \frac{a^{l-k}}{(l-k)!} \left[\frac{1}{k!} \prod_{j=0}^{k-1} (|\mu| + 2j) \right] (aT)^k \|(K + i(p-2)t)^{l-k} f; X\| \\
&\quad + \|f; G^a(K + i(p-2)t; X)\| \\
&\leq \left[\sum_{k=1}^{\infty} \frac{(aT)^k}{k!} \prod_{j=0}^{k-1} (|\mu| + 2j) + 1 \right] \|f; G^a(K + i(p-2)t; X)\|.
\end{aligned}$$

The lemma then follows by the following inequalities:

$$\begin{aligned}
\frac{1}{k!} \prod_{j=0}^{k-1} (|\mu| + 2j) &= |\mu| \prod_{j=1}^{k-1} \left(2 + \frac{|\mu| - 2}{j + 1} \right) \\
&\leq |\mu| \prod_{j=1}^{k-1} \left(2 + \text{Max} \left(\frac{|\mu| - 2}{2}, 0 \right) \right) = |\mu| m^{k-1}. \quad \square
\end{aligned}$$

Lemma 3. Let (q_j, r_j) , $j = 0, 1, 2, 3$, satisfy $1 \leq q_j, r_j \leq \infty$, $1/q_0 = \sum_{j=1}^3 1/q_j$, $1/r_0 = \sum_{j=1}^3 1/r_j$. Then:

(1) For any $a, T > 0$

$$\begin{aligned}
&\left\| \psi_1 \int_{-\infty}^x \psi_2 \overline{\psi_3}(t, x', y) dx'; G^a((J_x, 0); L^{q_0}(I; L_y^{r_0} L_x^2)) \right\| \\
(2.2) \quad &\leq \prod_{j=1}^3 \left\| \psi_j; G^a((J_x, 0); L^{q_j}(I; L_y^{r_j} L_x^2)) \right\|,
\end{aligned}$$

$$\begin{aligned}
&\left\| \psi_1 \int_{-\infty}^x \psi_2 \overline{\psi_3}(t, x', y) dx'; G^a((0, J_y); L^{q_0}(I; L_y^{r_0} L_x^2)) \right\| \\
(2.3) \quad &\leq \prod_{j=1}^3 \left\| \psi_j; G^a((0, J_y); L^{q_j}(I; L_y^{r_j} L_x^2)) \right\|,
\end{aligned}$$

$$\begin{aligned}
&\left\| \psi_1 \int_{-\infty}^x \psi_2 \overline{\psi_3}(t, x', y) dx'; G^a(J; L^{q_0}(I; L_y^{r_0} L_x^2)) \right\| \\
(2.4) \quad &\leq \prod_{j=1}^3 \left\| \psi_j; G^a(J; L^{q_j}(I; L_y^{r_j} L_x^2)) \right\|,
\end{aligned}$$

where $I = [0, T]$.

(2) For any $a, T > 0$ with $2aT < 1$

$$(2.5) \quad \begin{aligned} & \left\| \psi_1 \int_{-\infty}^x \psi_2 \overline{\psi_3}(t, x', y) dx'; G^a(J, K - 2it; L^{q_0}(I; L_y^{r_0} L_x^2)) \right\| \\ & \leq \frac{e^a}{(1 - 2aT)^3} \prod_{j=1}^3 \left\| \psi_j; G^a(J, K; L^{q_j}(I; L_y^{r_j} L_x^2)) \right\|, \end{aligned}$$

where $I = [0, T]$.

Proof. We define $\tilde{\psi}_j = M_x^{-1} \psi_j$, $M_x = M_x(t) = \exp(ix^2/(2t))$. By the relation $J_x = M_x(it\partial_x) M_x^{-1}$, we obtain

$$(2.6) \quad \begin{aligned} & J_x^l \left(\psi_1 \int_{-\infty}^x \psi_2 \overline{\psi_3} dx' \right) = M_x(it\partial_x)^l \left(\tilde{\psi}_1 \int_{-\infty}^x \tilde{\psi}_2 \overline{\tilde{\psi}_3} dx' \right) \\ & = M_x(it)^l \left(\partial_x^l \tilde{\psi}_1 \cdot \int_{-\infty}^x \tilde{\psi}_2 \overline{\tilde{\psi}_3} dx' + \sum_{\substack{l_1+l_2=l \\ l_2 \geq 1}} \frac{l!}{l_1! l_2!} \partial_x^{l_1} \tilde{\psi}_1 \cdot \int_{-\infty}^x \partial_x^{l_2} (\tilde{\psi}_2 \overline{\tilde{\psi}_3}) dx' \right) \\ & = M_x(it)^l \sum_{l_1+l_2+l_3=l} \frac{l!}{k_1! k_2! k_3!} \partial_x^{k_1} \tilde{\psi}_1 \int_{-\infty}^x \partial_x^{k_2} \tilde{\psi}_2 \cdot \overline{\partial_x^{k_3} \tilde{\psi}_3} dx' \\ & = \sum_{l_1+l_2+l_3=l} \frac{l!(-1)^{k_3}}{k_1! k_2! k_3!} J_x^{k_1} \psi_1 \int_{-\infty}^x J_x^{k_2} \psi_2 \overline{J_x^{k_3} \psi_3} dx'. \end{aligned}$$

We estimate (2.6) by the Hölder inequalities in space-time as

$$\begin{aligned} & \left\| J_x^l \left(\psi_1 \int_{-\infty}^x \psi_2 \overline{\psi_3} dx' \right); L^{q_0}(I; L_y^{r_0} L_x^2) \right\| \\ & \leq l! \sum_{k_1+k_2+k_3=l} \prod_{j=1}^3 \frac{1}{k_j!} \left\| J_x^{k_j} \psi_j; L^{q_j}(I; L_y^{r_j} L_x^2) \right\|. \end{aligned}$$

This implies

$$\begin{aligned} & \left\| \psi_1 \int_{-\infty}^x \psi_2 \overline{\psi_3} dx'; G^a((J_x, 0); L^{q_0}(I; L_y^{r_0} L_x^2)) \right\| \\ & \leq \sum_{l=0}^{\infty} \sum_{k_1+k_2+k_3=l} \prod_{j=1}^3 \frac{1}{k_j!} \left\| J^{k_j} \psi_j; L^{q_j}(I; L_y^{r_j} L_x^2) \right\| \\ & \leq \prod_{j=1}^3 \left\| \psi_j; G^a((J_x, 0); L^{q_0}(I; L_y^{r_0} L_x^2)) \right\|. \end{aligned}$$

This proves (2.2). By a similar calculation, we obtain (2.3). Similarly, (2.4) follows from a two dimensional generalization of (2.6).

To prove (2.5), by using

$$\begin{aligned} P \int_{-\infty}^x f \, dx' &= \int_{-\infty}^x (P+1) f \, dx', \\ (P+1)^k (fg) &= \sum_{j=0}^k \binom{k}{j} P^{k-j} f \cdot (P+1)^j g, \end{aligned}$$

we compute

$$\begin{aligned} & J^\alpha (K - 2it)^l \left(\psi_1 \int_{-\infty}^x \psi_2 \overline{\psi_3} \, dx' \right) \\ &= (2it)^{|\alpha|+l} M \partial^\alpha P^l \left(\tilde{\psi}_1 \int_{-\infty}^x \tilde{\psi}_2 \overline{\tilde{\psi}_3} \, dx' \right) \\ &= (2it)^{|\alpha|+l} M \sum_{\substack{\beta' + \beta'' + \beta''' = \alpha \\ j_1 + j_2 + j_3 = l}} \frac{\alpha! l! (-1)^{|\beta'''|}}{j_1! j_2! j_3! \beta'! \beta''! \beta'''!} \\ &\quad \cdot \partial^{\beta'} P^{j_1} \tilde{\psi}_1 \int_{-\infty}^x \partial^{\beta''} P^{j_2} \tilde{\psi}_2 \cdot \overline{\partial^{\beta'''} (P+1)^{j_3} \tilde{\psi}_3} \, dx' \\ &= \sum_{\substack{\beta' + \beta'' + \beta''' = \alpha \\ j_1 + j_2 + j_3 = l}} \frac{\alpha! l! (-1)^{|\beta'''|}}{j_1! j_2! j_3! \beta'! \beta''! \beta'''!} \\ &\quad \cdot J^{\beta'} (K - 2it)^{j_1} \psi_1 \int_{-\infty}^x J^{\beta''} (K - 2it)^{j_2} \psi_2 \cdot \overline{J^{\beta'''} (K - 2it + 1)^{j_3} \psi_3} \, dx'. \end{aligned}$$

In the same way as above, we estimate

$$\begin{aligned} & \left\| \psi_1 \int_{-\infty}^x \psi_2 \overline{\psi_3} \, dx'; G^a(J, K - 2it; L^{q_0}(I; L_y^0 L_x^2)) \right\| \\ & \leq \left\| \psi_3; G^a(J, K - 2it + 1; L^{q_3}(I; L_y^{r_3} L_x^2)) \right\| \prod_{j=1}^2 \left\| \psi_j; G^a(J, K - 2it; L^{q_j}(I; L_y^{r_j} L_x^2)) \right\| \\ & \leq e^a \prod_{j=1}^3 \left\| \psi_j; G^a(J, K - 2it; L^{q_j}(I; L_y^{r_j} L_x^2)) \right\| \\ & \leq \frac{e^a}{(1 - 2aT)^3} \prod_{j=1}^3 \left\| \psi_j; G^a(J, K; L^{q_j}(I; L_y^{r_j} L_x^2)) \right\|, \end{aligned}$$

where we have used Lemma 2 and the inequality

$$\|f; G^a(A+1; X)\| \leq e^a \|f; G^a(A; X)\|,$$

which holds for any operator A with $[A, 1] = 0$.

□

3. Proof of Theorem 1

We solve the integral equation

$$(3.1) \quad u(t) = U(t)\phi - i \int_0^t U(t-t')f(u(t')) \, dt'$$

by a contraction argument on X_T^a , Y_T^a , Z_T^a , and W_T^a . Let $\phi \in G^a((x, 0); L^2)$ with $\|\phi; G^a((x, 0); L^2)\| \leq \rho$. For $R > 0$ we define

$$X_T^a(R) = \{u \in X_T^a; \|u; X_T^a\| \leq R\},$$

where

$$\|u; X_T^a\| = \text{Max}(\|u; G^a((J_x, 0); L^\infty(0, T; L^2))\|, \|u; G^a((J_x, 0); L^8(0, T; L_y^4 L_x^2))\|).$$

For $u \in X_T^a$ we define $(\Phi(u))(t)$ as the RHS of (3.1). We have

$$(3.2) \quad (J_x^l \Phi(u))(t) = U(t)x^l \phi - i \int_0^t U(t-t') (J_x^l f(u))(t') \, dt'.$$

Applying Lemma 1 to the RHS of (3.2) and using the Hölder inequality in time, we obtain

$$(3.3) \quad \begin{aligned} & \text{Max}(\|J_x^l \Phi(u); L^\infty(0, T; L^2)\|, \|J_x^l \Phi(u); L^8(0, T; L_y^4 L_x^2)\|) \\ & \leq C \|x^l \phi; L^2\| + C \|J_x^l(f(u)); L^{4/3}(0, T; L_y^1 L_x^2)\| \\ & \leq C \|x^l \phi; L^2\| + CT^{1/2} \|J_x^l(f(u)); L^4(0, T; L_y^1 L_x^2)\|. \end{aligned}$$

Multiplying both sides of (3.3) by $a^l/l!$, making a summation on l and applying Lemma 3, we obtain

$$\begin{aligned} \|\Phi(u); X_T^a\| & \leq \|\phi; G^a((x, 0); L^2)\| \\ & + CT^{1/2} \|u; G^a((J_x, 0); L^8(0, T; L_y^4 L_x^2))\|^2 \|u; G^a((J_x, 0); L^\infty(0, T; L^2))\|. \end{aligned}$$

In the same way as above, for $u, v \in X_T^a(R)$, $\Phi(u)$ satisfies

$$\begin{aligned} \|\Phi(u); X_T^a\| & \leq C\rho + CT^{1/2}R^3, \\ \|\Phi(u) - \Phi(v); X_T^a\| & \leq CT^{1/2}R^2 \|u - v; X_T^a\|. \end{aligned}$$

For $\rho > 0$ let R and T satisfy $R \geq 2C\rho$, $T \leq 1/(4C^2R^4)$. Then $\Phi(u)$ has a unique fixed point in $X_T^a(R)$. This proves Part (1). Parts (2) and (3) follow in the same way.

For Part (4) we write, with $u \in W_T^a$

$$(J^\alpha K^l \Phi(u))(t) = U(t)x^{\alpha_1}y^{\alpha_2}(x^2 + y^2)^l \phi$$

$$-i \int_0^t U(t-t') (J^\alpha(K+4it')^l f(u)) (t') dt',$$

where we have used the relation

$$K(t)U(t-t') = U(t-t') (K(t') + 4it'),$$

which follows from the commutation relation $[K, L] = -4itL$, see [17, 20]. In the same way as above, we have

$$\begin{aligned} & \|\Phi(u); W_T^a\| \\ & \leq C \|\phi; G^a((x, y), x^2 + y^2; L^2)\| \\ (3.4) \quad & + CT^{1/2} \|f(u); G^a(J, K + 4it; L^4(0, T; L_y^1 L_x^2))\|. \end{aligned}$$

By Lemmas 2 and 3, the last norm on the RHS of (3.4) is estimated as

$$\begin{aligned} & \|f(u); G^a(J, K + 4it; L^4(0, T; L_y^1 L_x^2))\| \\ & \leq \frac{1}{1 - (5/2)aT} \|f(u); G^a(J, K - 2it; L^4(0, T; L_y^1 L_x^2))\| \\ (3.5) \quad & \leq \frac{1}{1 - (5/2)aT} \frac{e^a}{(1 - 2aT)^3} \|u; G^a(J, K; L^8(0, T; L_y^4 L_x^2))\|^2 \\ & \cdot \|u; G^a(J, K; L^\infty(0, T; L^2))\|. \end{aligned}$$

By (3.4) and (3.5), we have for $u \in W_T^a(R)$

$$\|\Phi(u); W_T^a\| \leq C\rho + \frac{CT^{1/2}R^3e^a}{(1 - (5/2)aT)(1 - 2aT)^3}.$$

Similarly, for $u, v \in W_T^a(R)$

$$\|\Phi(u) - \Phi(v); W_T^a\| \leq \frac{CT^{1/2}R^2e^a}{(1 - (5/2)aT)(1 - 2aT)^3} \|u - v; W_T^a\|.$$

For $\rho > 0$ let R and T satisfy $R \geq 2C\rho$, $T \leq \text{Min}(1/(5a), 1/(20CR^2e^a)^2)$. Then $\Phi(u)$ has a unique fixed point in $W_T^a(R)$.

4. Proof of Proposition 1

First, we prove the last part of the proposition. The LHS of (1.5) is estimated as

$$\begin{aligned} & \|\phi; G^a((x, y), x^2 + y^2; L^2)\| \\ & = \sum_{j, k, l \geq 0} \frac{a^{j+k+l}}{j!k!l!} \| |x|^j |y|^k (x^2 + y^2)^l \phi; L^2 \| \end{aligned}$$

$$\begin{aligned}
&\geq \left\| \sum_{j,k,l \geq 0} \frac{a^{j+k+l}}{j!k!l!} |x|^j |y|^k (x^2 + y^2)^l \phi; L^2 \right\| \\
&= \left\| e^{a(|x|+|y|+x^2+y^2)} \phi; L^2 \right\|.
\end{aligned}$$

Therefore (1.5) fails to hold for $\phi = e^{-a(|x|+|y|+x^2+y^2)}$.

From now on we consider functions in \mathbb{R}^n . We use the standard multi-index notation.

Proposition 2. *Let n and l be positive integers and let $a > 0$. Let $w = (w_1, \dots, w_l)$ be functions of $x \in \mathbb{R}^n$. Then the following inequalities hold.*

$$\begin{aligned}
&\left\| \left(\prod_{j=1}^l e^{a|w_j|} \right) \phi; L^2(\mathbb{R}^n) \right\| \\
&\leq \sum_{\substack{\alpha \in \mathbb{Z}^l \\ \alpha \geq 0}} \frac{a^{|\alpha|}}{\alpha!} \|w^\alpha \phi; L^2(\mathbb{R}^n)\| \\
(4.1) \quad &\leq C_0^l \left\| \left(\prod_{j=1}^l (1 + a|w_j|)^{1/2} e^{a|w_j|} \right) \phi; L^2(\mathbb{R}^n) \right\|,
\end{aligned}$$

where $C_0 = (1 + 2 \log 2)^{1/2}$.

REMARK 4. To derive (1.3) from (4.1), we put $l = 2$, $w_1(x, y) = x$, $w_2(x, y) = y$. To derive (1.4) from (4.1), we put $l = 3$, $w_1(x, y) = x$, $w_2(x, y) = y$, $w_3(x, y) = x^2 + y^2$.

Proof of Proposition 2. The first inequality of (4.1) is proved by the Maclaurin expansion of e^x and the triangle inequality.

The second inequality of (4.1) is proved as follows. By the Wallis formula

$$\frac{(2k)!}{2^{2k}(k!)^2} \frac{1}{k} \sim \frac{1}{\sqrt{\pi}} \frac{1}{k^{1+1/2}},$$

we see that the series

$$\sum_{k \geq 0} \frac{(2k)!}{2^{2k}(k!)^2} \frac{1}{k_+}$$

converges to a finite value $C_0^2 (= 1 + 2 \log 2)$, where $k_+ = \max(k, 1)$. By Schwarz'

inequality, this yields the following estimate:

$$\begin{aligned}
 & \sum_{\alpha} \frac{a^{|\alpha|}}{\alpha!} \|w^{\alpha} \phi; L^2(\mathbb{R}^n)\| \\
 & \leq \left(\sum_{\alpha} \frac{(2\alpha)!}{2^{2|\alpha|} ((\alpha_1)!)^2 (\alpha_1)_+ \cdots (\alpha_l)_+} \right)^{1/2} \\
 & \quad \cdot \left(\sum_{\alpha} \frac{(\alpha_1)_+ \cdots (\alpha_l)_+ 2^{2|\alpha|}}{(2\alpha)!} \cdot a^{2|\alpha|} \|w^{\alpha} \phi; L^2(\mathbb{R}^n)\|^2 \right)^{1/2} \\
 (4.2) \quad & = C_0^l \left\langle \sum_{\alpha} \frac{(\alpha_1)_+ \cdots (\alpha_l)_+}{(2\alpha)!} (2aw)^{2\alpha} \phi, \phi \right\rangle^{1/2},
 \end{aligned}$$

where $\langle \cdot, \cdot \rangle$ is the scalar product in $L^2(\mathbb{R}^n)$. Since

$$\sum_{k=0}^{\infty} \frac{k_+}{(2k)!} \xi^{2k} = 1 + \frac{\xi}{2} \sinh \xi,$$

the RHS of (4.2) is dominated by

$$\begin{aligned}
 & C_0^l \left\langle \prod_{j=1}^l \left(1 + \frac{2aw_j}{2} \sinh 2aw_j \right) \phi, \phi \right\rangle^{1/2} \\
 & = C_0^l \left\| \prod_{j=1}^l \left(1 + \frac{2aw_j}{2} \sinh 2aw_j \right)^{1/2} \phi; L^2(\mathbb{R}^n) \right\|.
 \end{aligned}$$

Using $1 \leq (1 + \xi \sinh 2\xi) \leq (1 + |\xi|)e^{2|\xi|}$, we have the second inequality of (4.1). \square

References

- [1] J.-B. Baillon, T. Cazenave and M. Figueira: *Équation de Schrödinger avec non-linéarité intégrale*, C.R. Acad. Sci. Paris **284** (1977), 939–942.
- [2] D.J. Benney and G. L. Roskes: *Wave instabilities*, Stud. Appl. Math. **48** (1969), 377–387.
- [3] T. Cazenave: *Equations de Schrödinger non linéaires en dimension deux*, Proc. Royal Soc. Edinburgh **84** (1979), 327–346.
- [4] T. Cazenave: *Semilinear Schrödinger Equations*, Amer. Math. Soc. New York, 2003.
- [5] A. Davey and K. Stewartson: *On three-dimensional packets of surface waves*, Proc. Roy. Soc. London **338** (1974), 101–110.
- [6] V.D. Djordjević and L. G. Redekopp: *On two-dimensional packets of capillary-gravity waves*, J. Fluid Mech. **79** (1977), 703–714.
- [7] A. De Bouard: *Analytic solutions to non elliptic nonlinear Schrödinger equations*, J. Differential Equations **104** (1993), 196–213.

- [8] A. De Bouard, N. Hayashi and K. Kato: *Gevrey regularizing effect for the (generalized) Korteweg-de Vries equation and nonlinear Schrödinger equations*, Ann. Inst. H. Poincaré, Anal. Non Linéaire **12** (1995), 673–725.
- [9] J.-M. Ghidaglia and J.-C. Saut: *On the initial value problem for Davey-Stewartson systems*, Nonlinearity **3** (1990), 475–506.
- [10] N. Hayashi: *Global existence of small analytic solutions to nonlinear Schrödinger equations*, Duke Math. J. **60** (1990), 717–727.
- [11] N. Hayashi: *Analyticity of solutions of the Korteweg-de Vries equation*, SIAM J. Math. Anal. **22** (1991), 1738–1743.
- [12] N. Hayashi: *Smoothing effect of small analytic solutions to nonlinear Schrödinger equations*, Ann. Inst. H. Poincaré Phys. Théor. **57** (1992), 385–394.
- [13] N. Hayashi: *Local existence in time of small solutions to the Davey-Stewartson system*, Ann. Inst. H. Poincaré Phys. Théor. **65** (1996), 313–366.
- [14] N. Hayashi: *Local existence in time of solutions to the elliptic-hyperbolic Davey-Stewartson system without smallness condition on the data*, J. Anal. Math. **73** (1997), 133–164.
- [15] N. Hayashi: *Analytic function spaces and their applications to nonlinear evolution equations*; in Analytic Extension Formulas and their Applications (S. Saitoh, N. Hayashi, M. Yamamoto, Eds.), Kluwer Acad. Publ. 2001, 59–72.
- [16] N. Hayashi and K. Kato: *Regularity in time of solutions to nonlinear Schrödinger equations*, J. Funct. Anal. **128** (1995), 253–277.
- [17] N. Hayashi and K. Kato: *Analyticity in time and smoothing effect of solutions to nonlinear Schrödinger equations*, Commun. Math. Phys. **184** (1997), 273–300.
- [18] N. Hayashi and P.I. Naumkin: *On the Davey-Stewartson and Ishimori systems*, Math. Phys. Anal. Geom. **2** (1999), 53–81.
- [19] N. Hayashi, P.I. Naumkin and P.-N. Pipolo: *Analytic smoothing effects for some derivative nonlinear Schrödinger equations*, Tsukuba J. Math. **24** (2000), 21–34.
- [20] N. Hayashi and T. Ozawa: *Smoothing effect for some Schrödinger equations*, J. Funct. Anal. **85** (1989), 307–348.
- [21] N. Hayashi and T. Ozawa: *On the derivative nonlinear Schrödinger equation*, Phys. D **55** (1992), 14–36.
- [22] N. Hayashi and T. Ozawa: *Schrödinger equations with nonlinearity of integral type*, Discrete Contin. Dynam. Systems **1** (1995), 475–484.
- [23] N. Hayashi and S. Saitoh: *Analyticity and smoothing effect for the Schrödinger equation*, Ann. Inst. H. Poincaré Phys. Théor. **52** (1990), 163–173.
- [24] N. Hayashi and S. Saitoh: *Analyticity and global existence of small solutions to some nonlinear Schrödinger equations*, Commun. Math. Phys. **129** (1990), 27–41.
- [25] N. Hayashi and J.-C. Saut: *Global existence of small solutions to the Davey-Stewartson and Ishimori systems*, Differential Integral Equations **8** (1995), 1657–1675.
- [26] N. Hayashi and J.-C. Saut: *Global existence of small solutions to the Ishimori system without exponential decay on the data*, Differential Integral Equations **9** (1996), 1183–1195.
- [27] K. Kajitani: *Analytically smoothing effect for Schrödinger equations*, Discrete Contin. Dynam. Systems **I** (1998), 350–352.
- [28] K. Kajitani and S. Watabayashi: *Analytically smoothing effect for Schrödinger type equations with variable coefficients*; in Direct and Inverse Problems of Mathematical Physics, Kluwer Academic Publ. Dordrecht, 2000, 185–219.
- [29] K. Kato and T. Ogawa: *Analyticity and smoothing effect for the Korteweg-de Vries equation with a single point singularity*, Math. Ann. **316** (2000), 577–608.
- [30] K. Kato and K. Taniguchi: *Gevrey regularizing effect for nonlinear Schrödinger equations*, Osaka J. Math. **33** (1996), 863–880.
- [31] T. Kato and K. Masuda: *Nonlinear evolution equations and analyticity*, Ann. Inst. H. Poincaré Anal. Non Linéaire **3** (1986), 455–467.
- [32] F. Linares and G. Ponce: *On the Davey-Stewartson systems*, Ann. Inst. H. Poincaré Anal. Non Linéaire **10** (1993), 523–548.

- [33] Y. Morimoto, L. Robbiano and C. Zuily: *Remark on the analytic smoothing for the Schrödinger equation*, Indiana Univ. Math. J. **49** (2000) 1563–1579.
- [34] T. Ozawa and Y. Yamazaki: *Smoothing effect and large time behavior of solutions to Schrödinger equations with nonlinearity of integral type*, Commun. Contemp. Math. **6** (2004) 681–703.
- [35] L. Robbiano and C. Zuily: *Microlocal analytic smoothing effect for the Schrödinger equation*, Duke Math. J. **100** (1999) 93–129.
- [36] L. Robbiano and C. Zuily: *Effet régularisant microlocal analytique pour l'équation de Schrödinger: le cas des données oscillantes*, Comm. Partial Differential Equations **25** (2000) 1891–1906.
- [37] W.A. Strauss: *Mathematical aspects of classical nonlinear field equations*; in Nonlinear Problem in Theoretical Physics, Lecture Notes in Phys. **98** (1979), 123–149.
- [38] H. Takuwa: *Analytic smoothing effects for a class of dispersive equations*, Tsukuba J. Math. **28** (2004), 1–34.

T. Ozawa
Department of Mathematics
Hokkaido University
Sapporo 060-0810, Japan

K. Yamauchi
Department of Mathematics
Hokkaido University
Sapporo 060-0810, Japan

Current address:
Department of Business Administration
Faculty of Business Administration
Sapporo University
Nishioka 3-7, Sapporo 062-8520, Japan

Y. Yamazaki
Department of Mathematics
Hokkaido University
Sapporo 060-0810, Japan